Majorana Zero Modes in a Proximity-induced Superconducting Quantum Wire



November 23, 2018

1

Outline

- 1. Motivation
- 2. Background Theory
 - a. Cooper pairing in S-wave superconductors
 - b. Majorana Zero Modes
 - c. Non-Abelian Braiding of Majorana
- 3. Project Details
- 4. Summary

Motivation

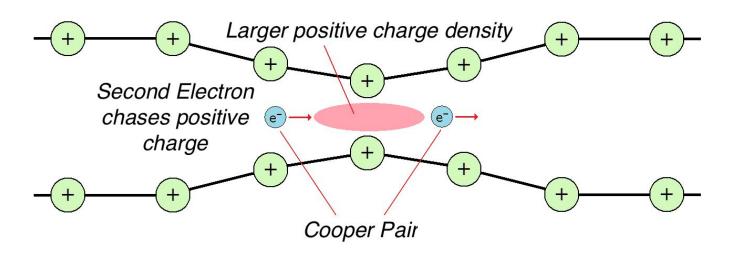
- Advantage of Quantum Computers
 - Decoherence in conventional quantum computers
- Topological Insulators as topologically-protected quantum computers
 - Physical fault-tolerance vs. correction coding
- Quasiparticle excitations in lower dimensional topological insulators
 - Majorana Zero Modes
 - Non-abelian Statistics

Motivation

- Previous successful theoretical work
 - Kitaev (2001)
 - Majorana zero modes in a p-wave superconducting wire
 - Fu and Kane (2007)
 - Majorana bound states in a chiral $p_x + ip_y 2D$ superconductor
- What's the problem?
 - P-wave superconductors theoretically exist, but are either extremely rare or non-existent
- Can we find a similar system using regular s-wave superconductors?

Theory: Cooper Pairing

- BCS Theory of Superconductivity
 - Long-range attractive force between electrons near Fermi surface

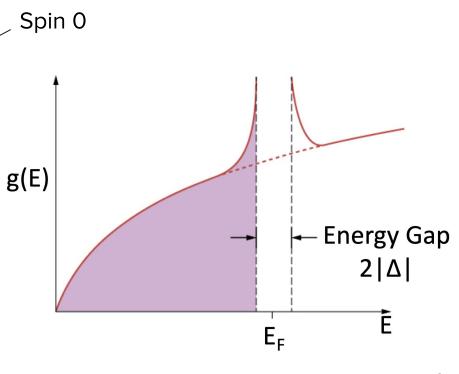


Theory: Cooper Pairing

 S-wave: Cooper pair in spin-singlet state

$$|\Psi\rangle = \frac{1}{\sqrt{2}} \left[|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle \right]$$

- Form Bose-Einstein condensate
 - Energetically favourable
 - superconducting gap △



- Excitations in superconducting systems
 - Coherent superpositions of particles and holes

$$\mathcal{H} = \int d\mathbf{r} \sum_{\sigma=\uparrow,\downarrow} (H_o(\mathbf{r})\alpha_{\sigma}^{\dagger}(\mathbf{r})\alpha_{\sigma}(\mathbf{r})) + \Delta \alpha_{\uparrow}^{\dagger}(\mathbf{r})\alpha_{\downarrow}^{\dagger}(\mathbf{r}) + \Delta^* \alpha_{\downarrow}(\mathbf{r})\alpha_{\uparrow}(\mathbf{r})$$

Single particle H (KE + PE) Creates electron at **r**with spin up
(or annihilates a
hole)

Annihilates electron at **r** with spin up (or creates a hole)

- Is there an easier way?
 - Described by Bogoliubov-de Gennes (BdG) equation

$$E_n \Phi_n = \mathcal{H}_{\mathrm{BdG}} \Phi_n$$

$$\mathcal{H} = egin{bmatrix} H_o & -i\sigma^2\Delta \ i\sigma^2\Delta^* & -H_o^* \end{bmatrix} \quad \Phi_n = egin{bmatrix} a_\downarrow \ a_\uparrow^\dagger \ a_\uparrow^\dagger \end{pmatrix}$$

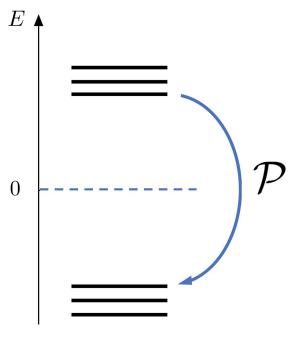
$$\Phi_n = \begin{pmatrix} a_\uparrow \\ a_\downarrow \\ a_\uparrow^\dagger \\ a_\downarrow^\dagger \\ a_\downarrow^\dagger \end{pmatrix}$$

 Particle-hole symmetry: Unitary mapping between states with positive energy E and negative energy -E

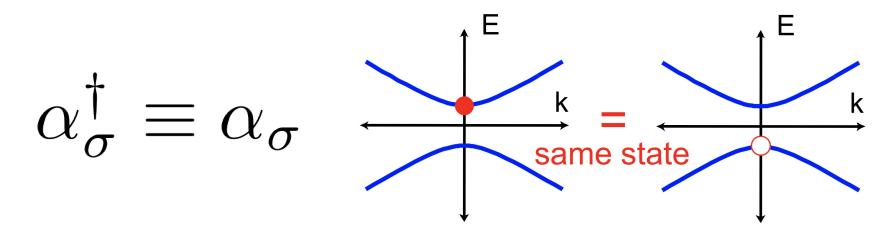
$$\Phi_{-E_n} = \mathcal{P}\Phi_{E_n}$$

$$\mathcal{PP}^{\dagger} = \mathbb{I}$$

Excitations must come in pairs!

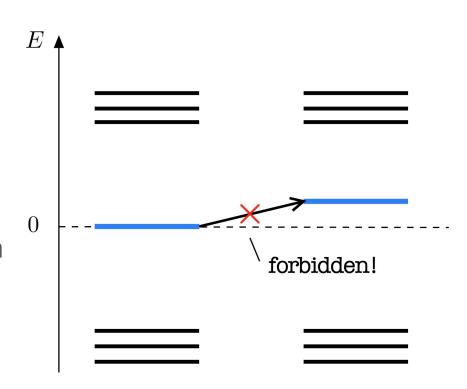


- What about E = 0?
 - Special state possible: Majorana Zero Mode



Does not support an energy pair as before

- These modes are topologically protected from decoherence
 - Local perturbations cannot influence E = 0 state
 - Change in topology of system required



Theory: Non-abelian Braiding

 In three dimensional systems, particles obey either Fermi-Dirac or Bose-Einstein Statistics

$$|\psi_1, \psi_2, ..., \psi_N\rangle = \pm |\psi_2, \psi_1, ..., \psi_N\rangle$$

• In lower dimensional systems, particles can obey different statistics

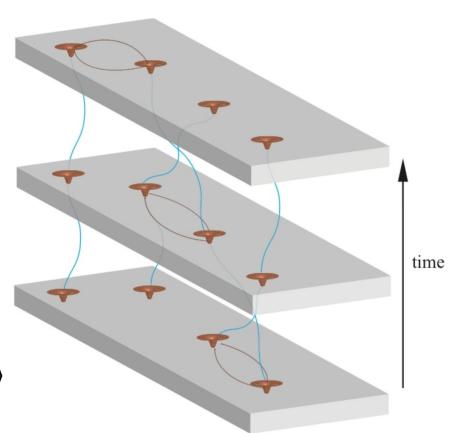
$$|\psi_1, \psi_2, ..., \psi_N\rangle = e^{i\theta} |\psi_2, \psi_1, ..., \psi_N\rangle$$

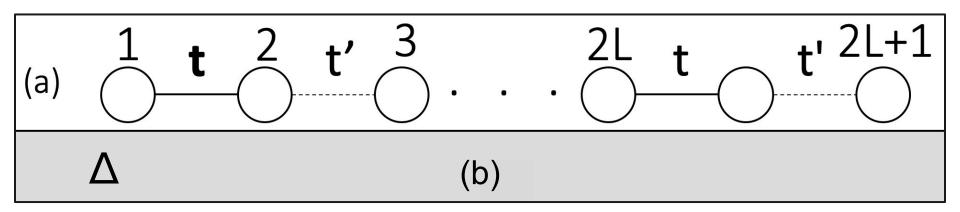
Arbitrary phase change

Theory: Non-abelian Braiding

- If these systems have degeneracy, particle exchanges can lead to state exchanges
 - Non-commuting unitary operations! Order of operation matters

$$U_1U_2|\Psi\rangle \neq U_2U_1|\Psi\rangle$$





- (a) 1D lattice with an odd number of sites (Tight-binding Model)
- (b) 3D s-wave superconductor

$$\mathcal{H} = \sum_{\sigma=\uparrow,\downarrow} \sum_{n=1}^{L} \left(t \alpha_{2n+1,\sigma}^{\dagger} \alpha_{2n,\sigma} + t^* \alpha_{2n,\sigma}^{\dagger} \alpha_{2n+1,\sigma} \right)$$

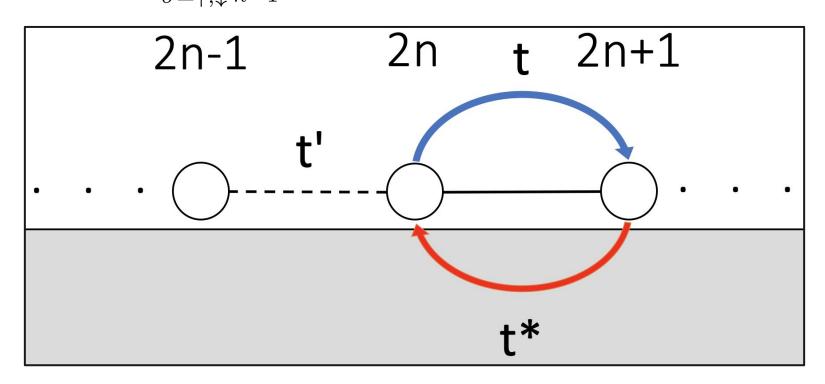
$$+ \sum_{\sigma=\uparrow,\downarrow} \sum_{n=1}^{L} \left(t' \alpha_{2n,\sigma}^{\dagger} \alpha_{2n-1,\sigma} + t'^* \alpha_{2n-1,\sigma}^{\dagger} \alpha_{2n,\sigma} \right)$$

$$+ \sum_{n=1}^{2L+1} \left(\Delta \alpha_{n,\uparrow} \alpha_{n,\downarrow} + \Delta^* \alpha_{n,\downarrow}^{\dagger} \alpha_{n,\uparrow}^{\dagger} \right)$$

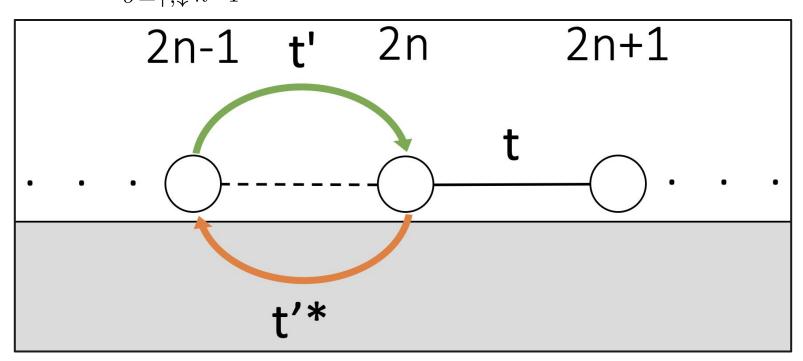
$$\{\alpha_{n,\sigma}, \alpha_{m,\rho}^{\dagger}\} = \delta_{nm}\delta_{\sigma\rho} \qquad \alpha_{0,\sigma} = \alpha_{2L+2,\sigma} = 0$$

$$\{\alpha_{n,\sigma}, \alpha_{m,\rho}\} = 0 = \{\alpha_{n,\sigma}^{\dagger}, \alpha_{m,\rho}^{\dagger}\}$$

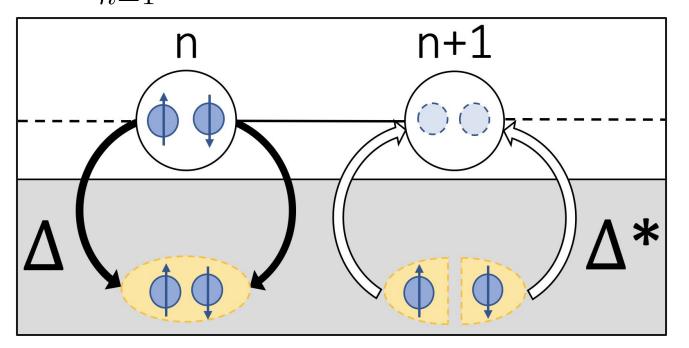
$$\mathcal{H} = \sum_{\sigma=\uparrow,\downarrow} \sum_{n=1}^{L} \left(\underline{t} \alpha_{2n+1,\sigma}^{\dagger} \alpha_{2n,\sigma} + \underline{t}^{*} \alpha_{2n,\sigma}^{\dagger} \alpha_{2n+1,\sigma} \right)$$



$$+\sum_{\sigma=\uparrow,\downarrow}\sum_{n=1}^{L}\left(\underline{t'\alpha_{2n,\sigma}^{\dagger}\alpha_{2n-1,\sigma}}+\underline{t'^{\ast}\alpha_{2n-1,\sigma}^{\dagger}\alpha_{2n,\sigma}}\right)$$



$$+\sum_{n=1}^{2L+1} \left(\Delta \alpha_{n,\uparrow} \alpha_{n,\downarrow} + \Delta^* \alpha_{n,\downarrow}^{\dagger} \alpha_{n,\uparrow}^{\dagger} \right)$$

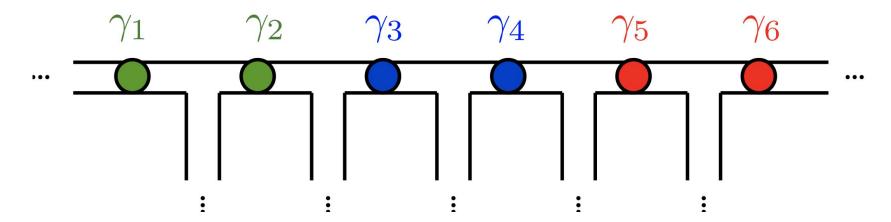


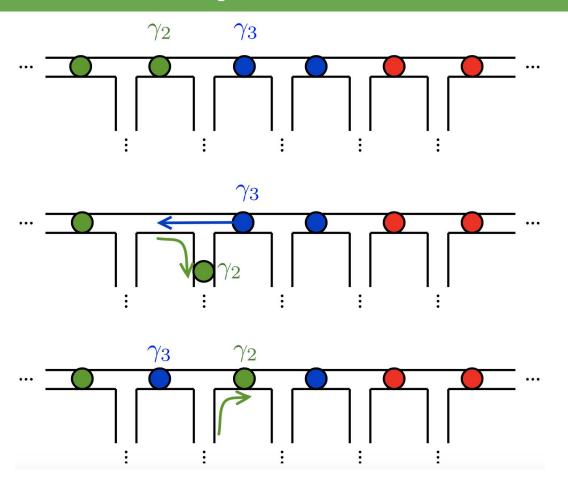
- Three important features of the system
 - 1. Odd-dimensional Hilbert space (single particle states)
 - 2. Single particle positive and negative energy states
 - 3. Particle-hole symmetry

Unpaired zero energy state

'Topological' attribute of the system

- What will we do with this system?
 - Solve for energy eigenstates via a BdG equation
 - Analyze Majorana character of E = 0 state
- Braiding operations between Majorana of neighbouring wires





Summary

- Topological Insulators may form the basis for topologically-protected quantum computing
 - Braiding operations of Majorana Zero Modes
- Physically-realizable system in which Majorana modes are possible
 - Analyze energy eigenstates of system + majorana character
 - Develop computation algorithms