The 3D Schrodinger Eqn. and Spherical Harmonics



#### Schrodinger Eqn

$$\begin{split} \frac{-\hbar\nabla^2}{2m} \Psi + V(r)\Psi &= E\Psi & \text{Schrodinger Eqn} \\ \Psi &= R(r)Y(\theta,\phi) & \text{separation ansatz} \end{split}$$

$$\label{eq:schrodinger} \frac{-\hbar\nabla^2}{2m} \underbrace{\Psi + V(r)\Psi = E\Psi}_{\Psi = R(r)Y(\theta,\phi)} \qquad \begin{array}{ll} \mbox{Schrodinger Eqn} \\ \mbox{separation ansatz} \end{array}$$

$$\sin\theta \frac{\partial}{\partial\theta} \left( \sin\theta \frac{\partial Y}{\partial\theta} \right) + \frac{\partial^2 Y}{\partial\phi^2} = -l(l+1)\sin^2\theta Y \quad \text{angular equation}$$

$$\begin{split} \frac{-\hbar\nabla^2}{2m} \Psi + V(r)\Psi &= E\Psi & \text{Schrodinger Eqn} \\ \Psi &= R(r)Y(\theta,\phi) & \text{separation ansatz} \\ \sin\theta \frac{\partial}{\partial \theta} \left(\sin\theta \frac{\partial Y}{\partial \theta}\right) + \frac{\partial^2 Y}{\partial \phi^2} &= -l(l+1)\sin^2\theta Y & \text{angular equation} \\ & \overbrace{\varphi = \theta} & \overbrace{\varphi = \varphi} & \overbrace{\varphi = \theta} & \overbrace{\varphi = \varphi = \phi} & \overbrace{\varphi = \phi} & \overbrace{\widehat{\varphi = \varphi} & \overbrace{\widehat{\widehat{\varphi = \varphi} & \overbrace{\widehat{\widehat{\varphi = \varphi} & \overbrace{\widehat{\widehat{\varphi = \varphi} & \overbrace{\widehat{\widehat{\varphi = \varphi} & \overbrace{\widehat{\widehat{\widehat{\widehat{} [\varphi = \varphi$$

#### spherical harmonics

$$\begin{split} Y_l^m &= \epsilon \sqrt{\frac{(2l+1)(l-|m|)!}{4\pi(l+|m|)!}} \ e^{im\phi} \ P_l^m(\cos\theta) \\ &\stackrel{\text{Condon-Shortley}}{\underset{\text{phase}}{\overset{\text{}}{=}}} \ \epsilon &= (-1)^m \ m \geq 0 \\ \epsilon &= 1 \ m < 0 \end{split} \qquad \begin{array}{c} \text{orthonormal} \\ \left\langle Y_l^m | Y_{l'}^{m'} \right\rangle &= \delta_{ll'} \delta_{mm'} \end{array} \end{split}$$

The spherical harmonics are easily visualized by counting the number of zero crossings Re[Y] possess in both the latitudinal (theta) and longitudinal directions (phi).



 $\phi$  - longitudinal  $\theta$  - latitudinal

$$\begin{split} \hline Y_0^0(\theta,\phi) &= \frac{1}{2}\sqrt{\frac{1}{\pi}} \\ P_0^0(\theta,\phi) &= \frac{1}{2}\sqrt{\frac{1}{\pi}} \\ \hline Y_1^{-1}(\theta,\phi) &= \frac{1}{2}\sqrt{\frac{3}{2\pi}}\sin\theta \ e^{-i\phi} \\ Y_1^0(\theta,\phi) &= \frac{1}{2}\sqrt{\frac{3}{2\pi}}\cos\theta \\ Y_1^1(\theta,\phi) &= \frac{-1}{2}\sqrt{\frac{3}{2\pi}}\sin\theta \ e^{i\phi} \\ \end{split} \qquad l = 1, m = 0, \pm 1 \\ \hline \mathbf{Re}\left\{Y_1^{\pm 1}\right\} \\ \hline \mathbf{Re}\left\{Y_1^0\right\} \\ \hline \mathbf{Re}\left\{Y_1^0\right\} \\ \end{split}$$

$$\begin{split} Y_2^{-2}(\theta,\phi) &= \frac{1}{4}\sqrt{\frac{15}{2\pi}}\sin^2\theta \ e^{-2i\phi} \\ Y_2^{-1}(\theta,\phi) &= \frac{1}{2}\sqrt{\frac{15}{2\pi}}\sin\theta\cos\theta \ e^{-i\phi} \\ Y_2^0(\theta,\phi) &= \frac{1}{4}\sqrt{\frac{5}{\pi}}(3\cos^2\theta - 1) \\ Y_2^1(\theta,\phi) &= \frac{-1}{2}\sqrt{\frac{15}{2\pi}}\sin\theta\cos\theta \ e^{i\phi} \\ Y_2^2(\theta,\phi) &= \frac{1}{4}\sqrt{\frac{15}{2\pi}}\sin^2\theta \ e^{2i\phi} \end{split}$$

$$k = 2, m = 0, \pm 1, \pm 2$$

The spherical harmonics are easily visualized by counting the number of zero crossings they possess in both the latitudinal (theta) and longitudinal directions (phi).

They possess (I-|m|) zeros along the theta direction, and 2|m| zeros along the phi direction.

This means Y is equal to 0 along m great circles passing through the poles, and along I-m circles of equal latitude.



 $Y_l^m(\theta,\phi)$  is equal to 0 along m great circles passing through the poles, and along l-m circles of equal latitude.



 $\phi$  - longitudinal  $\theta$  - latitudinal

Figures from the Wikimedia Commons.



is equal to 0 along m great circles passing through the poles, and along I-m circles of equal latitude. The function changes sign each time it crosses one of these lines

$$\Psi_{n,l,m} = A\left(\frac{2r}{na}\right) e^{-r/na} \left[L_{n-l-1}^{2l+1}(2r/na)\right] Y_l^m(\theta,\phi)$$
  
degree of polynomial  
 $n-l-1$  radial zeros



classical picture

quantum picture

Figures from the Wikimedia Commons.

$$\Psi_{n,l,m} = A\left(\frac{2r}{na}\right) e^{-r/na} \begin{bmatrix} L_{n-l-1}^{2l+1}(2r/na) \end{bmatrix} Y_l^m(\theta,\phi)$$
  
degree of polynomial  
 $n-l-1$  radial zeros

**Ground state:** 
$$n = 1, l = 0, m = 0$$

$$\Psi_{100} = \frac{1}{\sqrt{\pi a^3}} \ e^{-r/a}$$



cross-section of the probability density black = zero density, white = highest density

$$\Psi_{n,l,m} = A\left(\frac{2r}{na}\right) e^{-r/na} \left[L_{n-l-1}^{2l+1}(2r/na)\right] Y_l^m(\theta,\phi)$$
  
degree of polynomial  
 $n-l-1$  radial zeros

#### First excited state:

$$n = 2, \ l = 0, \ m = 0$$
  
 $n = 2, \ l = 1, \ m = \{-1, 0, +1\}$   
4 states



#### one radial node

no radial nodes

The probability density in three-dimensional space is obtained by rotating the one shown here around the z-axis



# Quantum jumps: quantum state change

# Quantum jumps: excitation by absorption of light



### Quantum jumps: dexcitation by emission of light



#### Quantum jumps: dexcitation by emission



# Image of atomic Fluorescence

cloud of 50 million Lithium atoms laser cooled to ImK

2p to 2s transition produces 671 nm (red) light