# Entanglement Entropy in Nonlocal Theories

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Based on work with Joanna Karczmarek 1307.3517 and work in progress



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- Motivations:
  - Violations of area law
  - Interest in noncommutative theories
  - Scrambling

#### Theories Considered

hep-th/9907166 hep-th/9908134 hep-th/0103090

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$$(f \star g)(x,y) = \left[ e^{\frac{i}{2}\theta \left( \frac{\partial}{\partial \xi_1} \frac{\partial}{\partial \zeta_2} - \frac{\partial}{\partial \zeta_1} \frac{\partial}{\partial \xi_2} \right)} f(x + \xi_1, y + \zeta_1) g(x + \xi_2, y + \zeta_2) \right]_{\xi_1 = \zeta_1 = \xi_2 = \zeta_2 = 0}$$

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- Dipole length  $\theta p$ .
- Dipole Theory:

$$(f \tilde{\star} g)(\vec{x}) = f\left(\vec{x} - \frac{\vec{L}_g}{2}\right) g\left(\vec{x} + \frac{\vec{L}_f}{2}\right)$$

• Fixed dipole length L.

# Features of the Gravity Duals

Non trivial geometry of the compact dimensions and dilaton

$$S = \frac{1}{32\pi^6 \alpha'^4} \int d^8 \sigma e^{-2\phi} \sqrt{G_{ind}^{(8)}}$$

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⇒ Open and closed string metric

Causality is determined by closed string metric

$$\frac{ds^2}{R^2} = u^2 \left( -dt^2 + dz^2 + f(u) \left[ dx^2 + dy^2 \right] \right) + \frac{du^2}{u^2} + \dots$$

$$f^{-1}(u) = 1 + (a_{\theta}u)^4 \qquad a_{\theta} = \lambda^{\frac{1}{4}} \theta^{\frac{1}{2}}$$

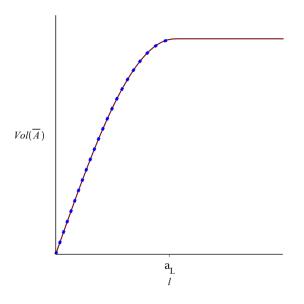
$$f^{-1}(u) = 1 + (a_L u)^2 \qquad a_L = \lambda^{\frac{1}{2}} L$$

# Entanglement Entropy of a Strip

Dipole:

$$S_A \xrightarrow[l \to \infty]{l \to 0} \frac{N^2}{2\pi} |A| \epsilon^{-3}$$

$$\xrightarrow[l \to \infty]{l \to \infty} \frac{N^2}{3\pi} |a_L| \partial A |\epsilon^{-3}|$$

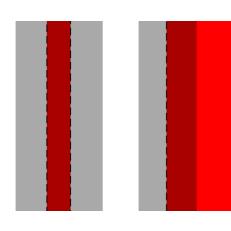


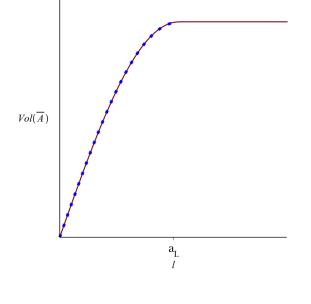
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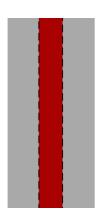


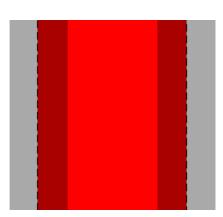
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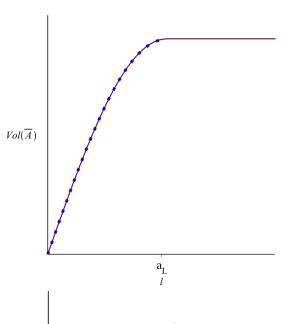
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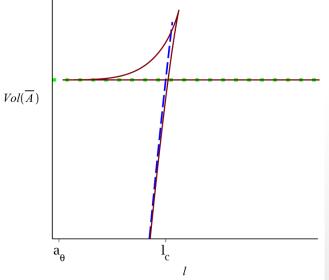




#### Noncommutative

- First order at  $l_c = \frac{1}{2}a_{\theta}^2/\epsilon$ .
- $S_A = \frac{N^2}{2\pi} \frac{|A|}{\epsilon^3}$  for all l as  $\epsilon \to 0$ .





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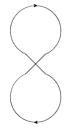
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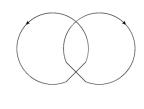
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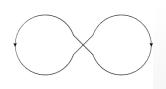
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- No! Leading divergence is same order as local theory
- Known story:
  - Only non-planar diagram gives new contribution
  - No higher order UV divergences
  - This story can be extended to gauge fields







#### The End

 Noncommutative field theory exhibits a volume law in strong coupling, large N limit.

 Intuitive picture for understanding the volume law.

# **Gravity Duals**

Noncommutative field theory:

$$\frac{ds^{2}}{R^{2}} = u^{2} \left( -dt^{2} + dz^{2} + f(u) \left[ dx^{2} + dy^{2} \right] \right) + \frac{du^{2}}{u^{2}} + d\Omega_{5}^{2}$$

$$e^{2\phi} = g_{s}^{2} f(u) \qquad B_{xy} = \frac{1}{\theta} \left( 1 - f(u) \right)$$

$$f(u) = \frac{1}{1 + (a_{\theta}u)^{4}}$$

Dipole Theory:

$$\frac{ds^{2}}{R^{2}} = u^{2} \left( -dt^{2} + dy^{2} + dz^{2} + f(u)dx^{2} \right) + \frac{du^{2}}{u^{2}} + f(u)d\psi^{2} + V_{\mathbb{C}\mathbf{P}^{2}}$$

$$e^{2\phi} = g_{s}^{2} f(u) \qquad B_{x\psi} = -\frac{1}{L} \left( 1 - f(u) \right)$$

$$f(u) = \frac{1}{1 + (\lambda Lu)^{2}}, \qquad \int V_{\mathbb{C}\mathbf{P}^{2}} = \frac{\pi^{2}}{2}.$$