

Quantum Entanglement of Local Operators

Masahiro Nozaki

Yukawa Institute for Theoretical Physics (YITP), Kyoto University

1. Based on arXiv:1401.0539v1 [hep-th] (Phys. Rev. Lett. 112, 111602 (2014)) with Tokiro Numasawa, Tadashi Takayanagi

2. Based on arXiv:1405.5875 [hep-th]



Motivation

We study the property of

Renyi Entanglement Entropy
(R)EE.

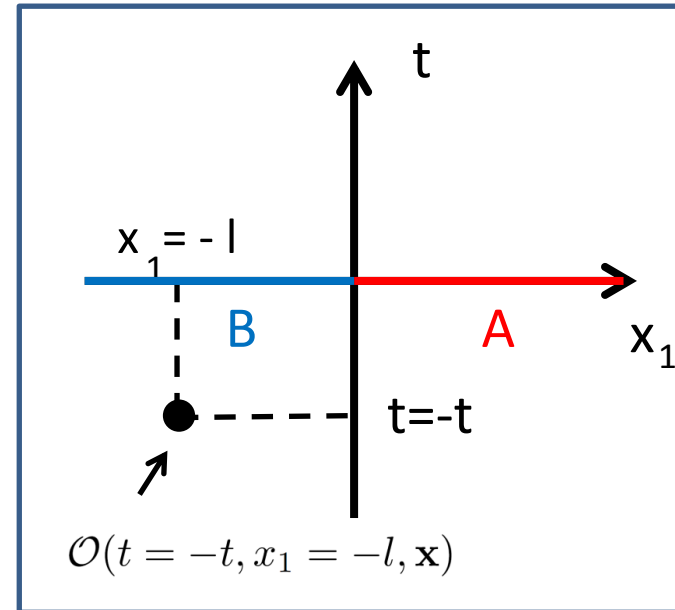
1. The size of subsystem is *infinite*.

A half of the total system:

$$x^1 \geq 0$$

2. A state is defined by acting a local operator
on the ground state:

$$|\Psi\rangle = \mathcal{N}^{-1} \mathcal{O}(t = -t, x_1 = -l, \mathbf{x}) |0\rangle .$$



Motivation

We study the property of

Renyi Entanglement Entropy
(REE)

$$\Delta S_A^{(n)}$$



At late time

*Some
Constants*

2. A state is defined by acting a local operator

on the ground state:

$$|\Psi\rangle = \mathcal{N}^{-1} \mathcal{O}(t = -t, x_1 = -l, \mathbf{x}) |0\rangle .$$

Motivation

We study the property of

Renyi Entanglement Entropy
(REE)

$$\Delta S_A^{(n)}$$



At late time

*Some
Constants*



Unique Behavior

$$|\Psi\rangle = \mathcal{N}^{-1} \mathcal{O}(t = -t, x_1 = -l, \mathbf{x}) |0\rangle .$$

The definition of $\Delta S_A^{(n)}$

$\Delta S_A^{(n)}$ is defined by the excess of REE:

$$\Delta S_A^{(n)} = S_A^{(n)Ex} - S_A^{(n)G},$$

where

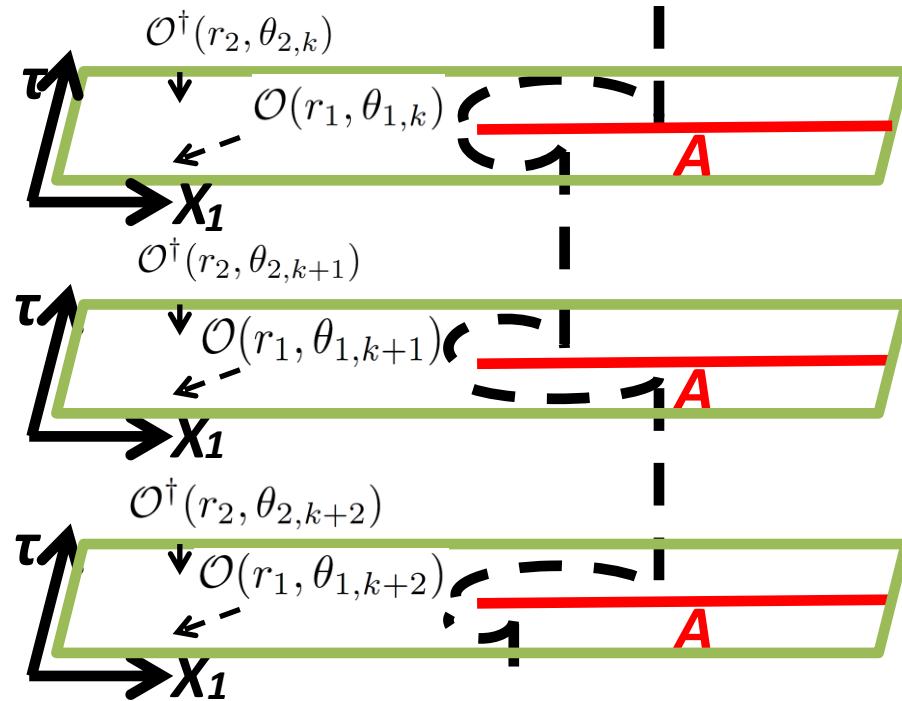
- REE for $|\Psi\rangle = \mathcal{N}^{-1} \mathcal{O}(t, x^1) |0\rangle$:

$$S_A^{(n)Ex} \sim \frac{1}{1-n} \log \left[\frac{\int D\Phi \mathcal{O}^\dagger(r_1, \theta_{1,1}) \mathcal{O}(r_2, \theta_{2,1}) \cdots \mathcal{O}^\dagger(r_1, \theta_{1,n}) \mathcal{O}(r_2, \theta_{2,n})}{\left(\int D\Phi \mathcal{O}^\dagger(r_1, \theta_{1,1}) \mathcal{O}(r_2, \theta_{2,1}) \right)^n} \right]$$

- REE for Ground State:

$$S_A^{(n)} \sim \frac{1}{1-n} \log \left[\frac{Z_n}{Z_1^n} \right]$$

The definition of $\Delta S_A^{(n)}$

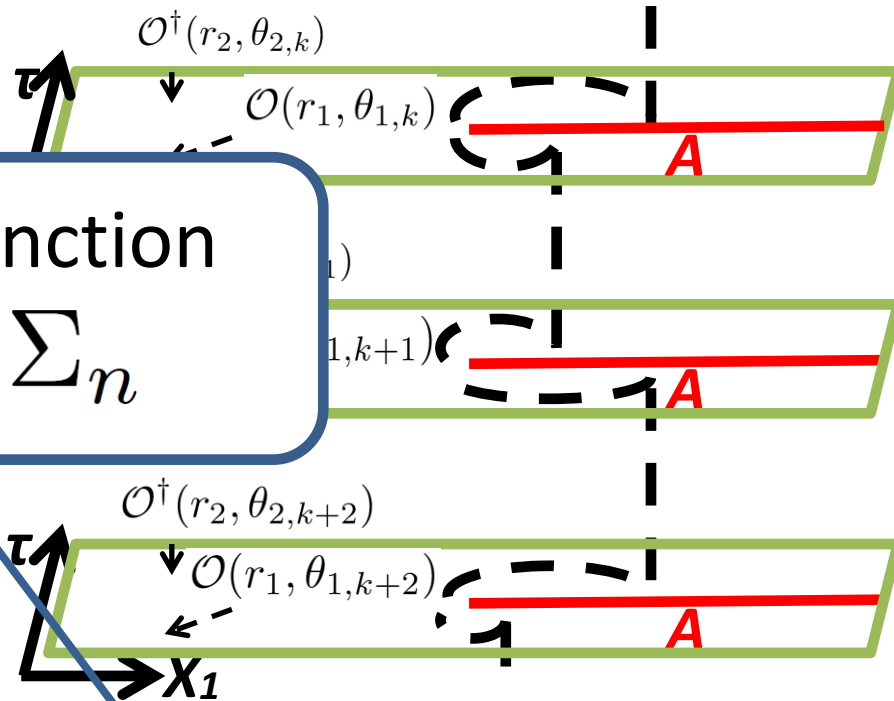


$$\Delta S_A^{(n)} = \frac{1}{1-n} \left(\log \langle \mathcal{O}^\dagger(r_2, \theta_{2,n}) \mathcal{O}(r_1, \theta_{1,n}) \cdots \mathcal{O}^\dagger(r_2, \theta_{2,1}) \mathcal{O}(r_1, \theta_{1,1}) \rangle_{\Sigma_n} - n \log \langle \mathcal{O}^\dagger(r_2, \theta_{2,1}) \mathcal{O}(r_1, \theta_{1,1}) \rangle_{\Sigma_1} \right).$$

$$\theta_{1,i} = \theta_1 + 2\pi(i-1)$$

$$\theta_{2,i} = \theta_2 + 2\pi(i-1)$$

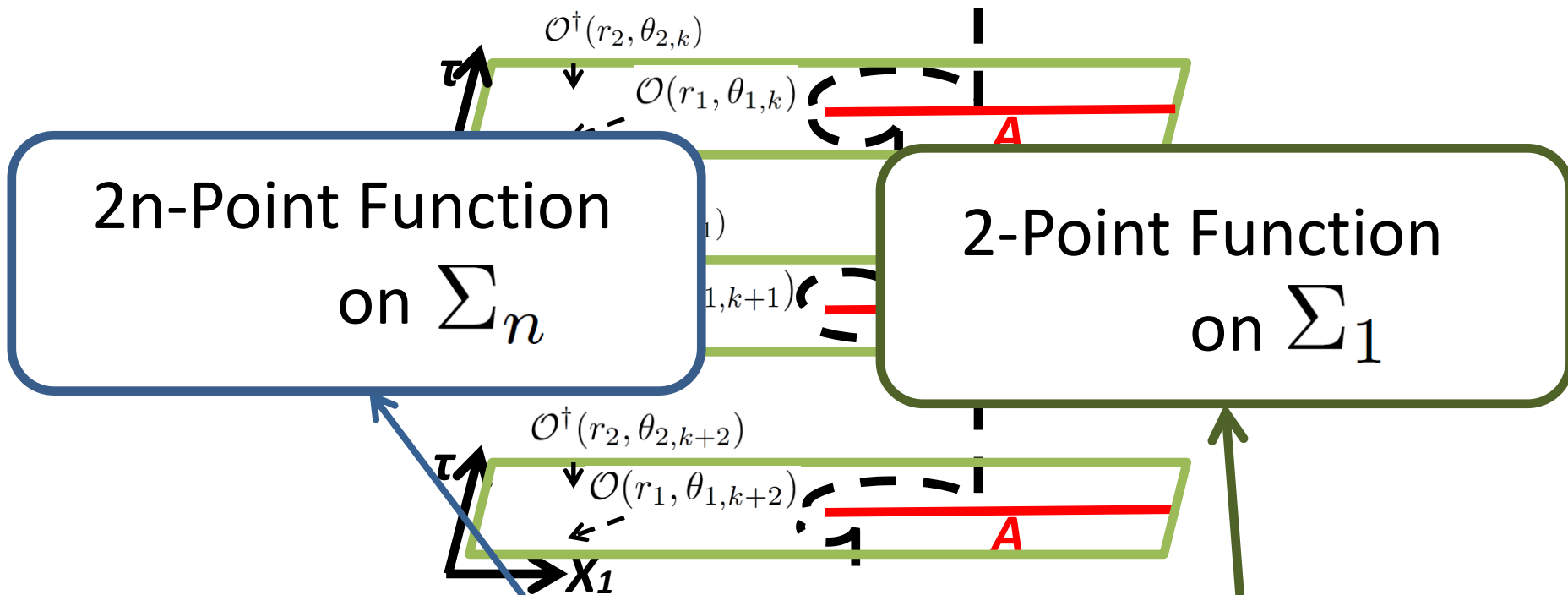
The definition of $\Delta S_A^{(n)}$



$$\Delta S_A^{(n)} = \frac{1}{1-n} \left(\log \left\langle \mathcal{O}^\dagger(r_2, \theta_{2,n}) \mathcal{O}(r_1, \theta_{1,n}) \cdots \mathcal{O}^\dagger(r_2, \theta_{2,1}) \mathcal{O}(r_1, \theta_{1,1}) \right\rangle_{\Sigma_n} - n \log \left\langle \mathcal{O}^\dagger(r_2, \theta_{2,1}) \mathcal{O}(r_1, \theta_{1,1}) \right\rangle_{\Sigma_1} \right).$$

$\theta_{1,i} = \theta_1 + 2\pi(i-1)$
 $\theta_{2,i} = \theta_2 + 2\pi(i-1)$

The definition of $\Delta S_A^{(n)}$



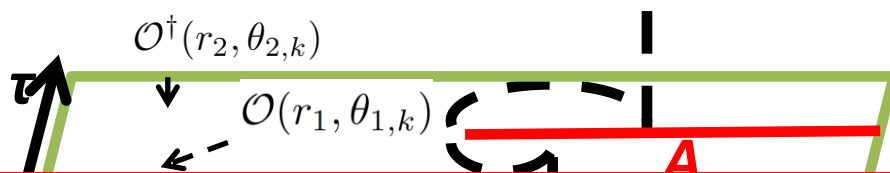
$$\Delta S_A^{(n)} = \frac{1}{1-n} \left(\log \langle \mathcal{O}^\dagger(r_2, \theta_{2,n}) \mathcal{O}(r_1, \theta_{1,n}) \cdots \mathcal{O}^\dagger(r_2, \theta_{2,1}) \mathcal{O}(r_1, \theta_{1,1}) \rangle_{\Sigma_n} - n \log \langle \mathcal{O}^\dagger(r_2, \theta_2) \mathcal{O}(r_1, \theta_1) \rangle_{\Sigma_1} \right).$$

Additional definitions for the angles are provided:

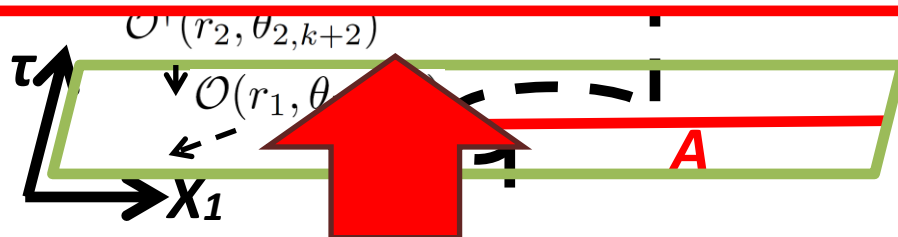
$$\theta_{1,i} = \theta_1 + 2\pi(i-1)$$

$$\theta_{2,i} = \theta_2 + 2\pi(i-1)$$

Replica Method



***This formula holds for any local operators
in general QFT in any dimensions.***



$$\Delta S_A^{(n)} =$$

$$\frac{1}{1-n} \left(\log \langle \mathcal{O}^\dagger(r_2, \theta_{2,n}) \mathcal{O}(r_1, \theta_{1,n}) \cdots \mathcal{O}^\dagger(r_2, \theta_{2,1}) \mathcal{O}(r_1, \theta_{1,1}) \rangle_{\Sigma_n} \right. \\ \left. - n \log \langle \mathcal{O}^\dagger(r_2, \theta_{2,1}) \mathcal{O}(r_1, \theta_{1,1}) \rangle_{\Sigma_1} \right).$$

$$\theta_{1,i} = \theta_1 + 2\pi(i-1)$$

$$\theta_{2,i} = \theta_2 + 2\pi(i-1)$$

Example

We consider *free massless scalar* field theory in *d+1 dim.*

Especially, we focus on that in *4 dim.*

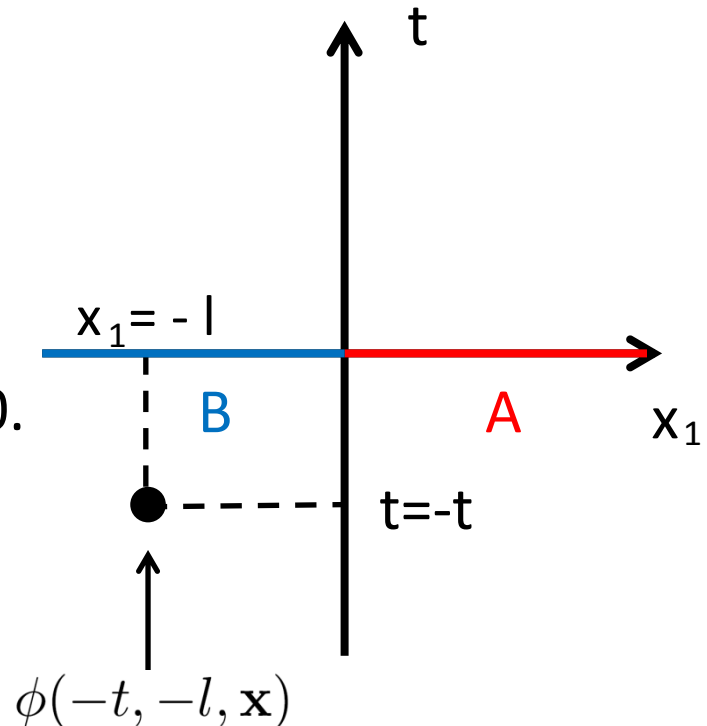
We act a local operator $\phi(-t, -l, \mathbf{x})$ on the ground state:

$$|\Psi\rangle = \mathcal{N}^{-1} \phi(-t, -l, \mathbf{x}) |0\rangle .$$

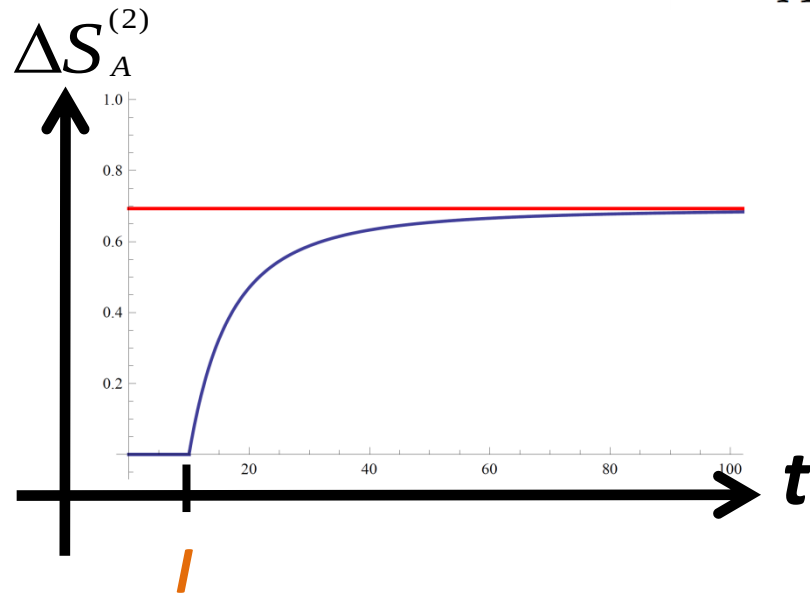
Subsystem *A*: $x^1 \geq 0$

We measure the second (Renyi) entanglement entropies $\Delta S_A^{(2)}$ at $t=0$.

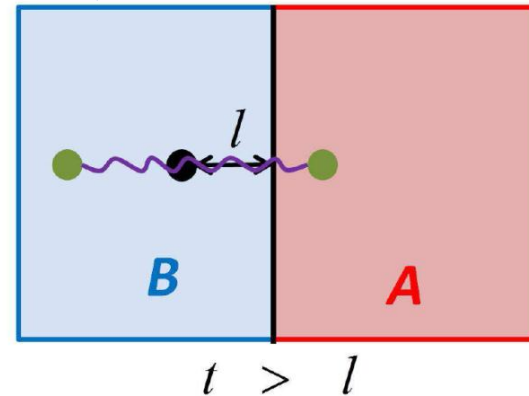
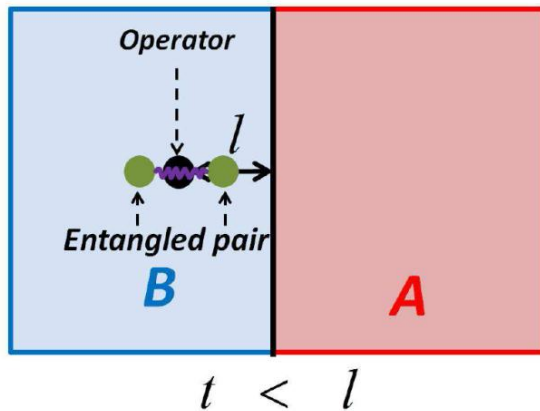
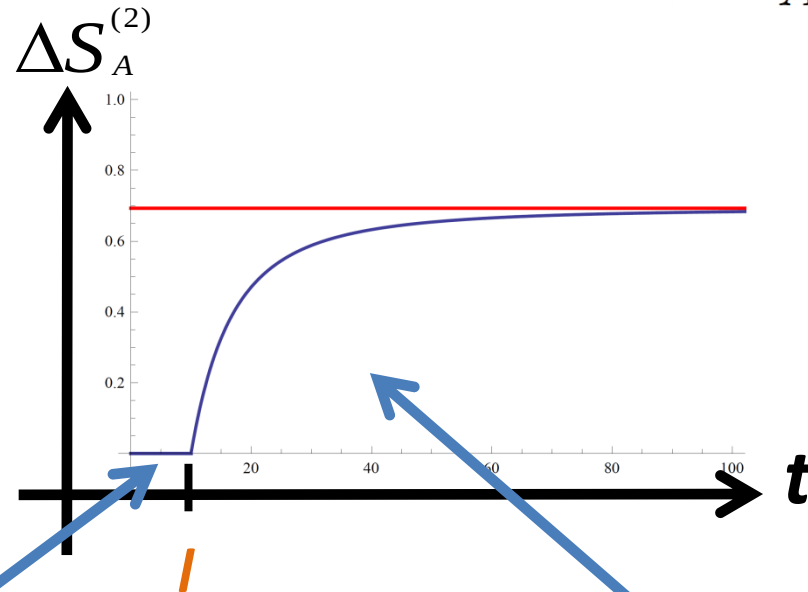
➡ *Time evolution!!*



Time Evolution of $\Delta S_A^{(2)}$

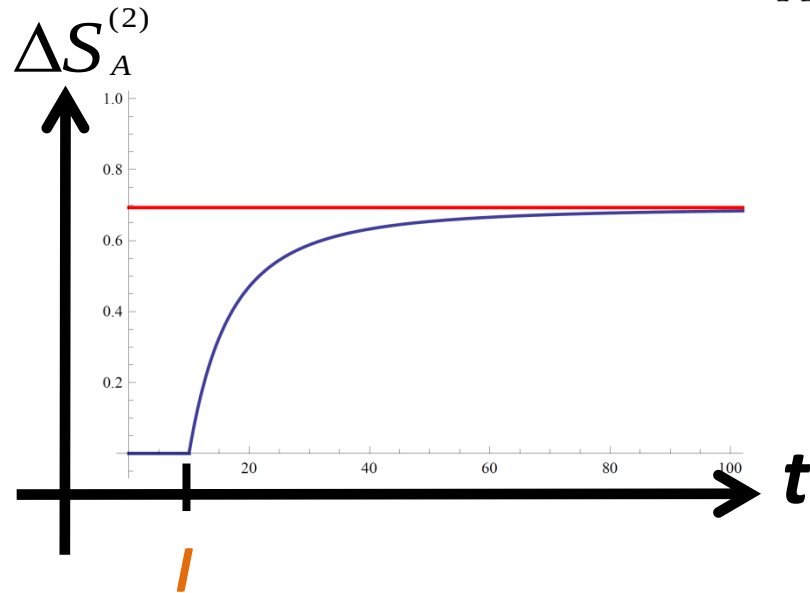


Time Evolution of $\Delta S_A^{(2)}$



We can interpret this behavior in terms of *the relativistic propagation of an entangled pair*.

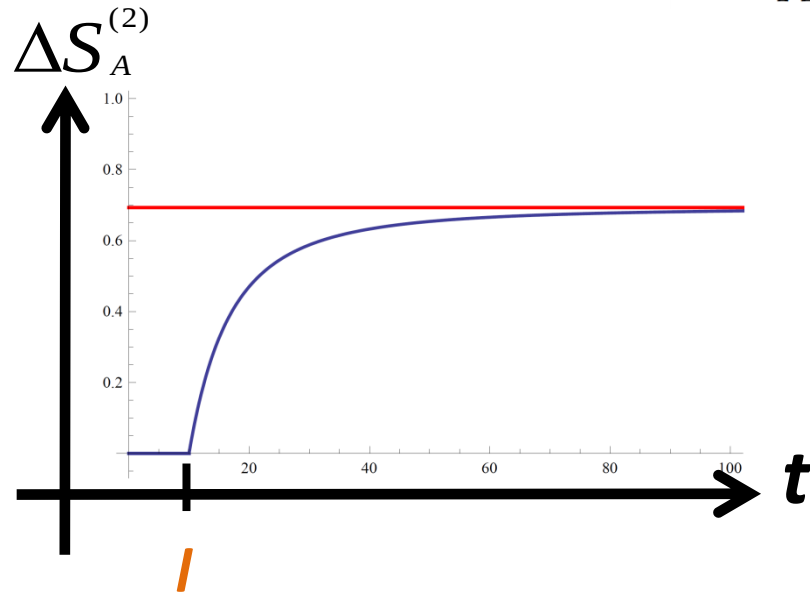
Time Evolution of $\Delta S_A^{(2)}$



$\Delta S_A^{(2)}$ for $|\Psi\rangle = \mathcal{N}^{-1} \phi(-t, -l, \mathbf{x}) |0\rangle$ approaches **constants!!**
 (log2)

We call them **the (Renyi) entanglement entropies of operators.**

Time Evolution of $\Delta S_A^{(2)}$



$\Delta S_A^{(2)}$ for $|\Psi\rangle = \mathcal{N}^{-1} \phi(-t, -l, \mathbf{x}) |0\rangle$ approaches **constants!!**
 (log2)

We call them **the (Renyi) entanglement entropies of operators.**

Generalize Results

We defined *the (Renyi) entanglement entropies of operators* by the late time values of $\Delta S_A^{(n)}$.

The (Renyi) entanglement entropies of specific operators $(: (\partial^m \phi)^k :)$ which are composed of single species operator are given by

$$\Delta S_A^{(n)f} = \frac{1}{1-n} \log \left(\frac{1}{2^{nk}} \sum_{j=0}^k \binom{k}{j} C_j^n \right).$$
$$\Delta S_A = k \cdot \log 2 - \frac{1}{2^k} \sum_{j=0}^k \binom{k}{j} C_j \log \binom{k}{j} C_j.$$

for *any dimension*.

 They characterize the local operators from the viewpoint of quantum entanglement!!

Sum rule

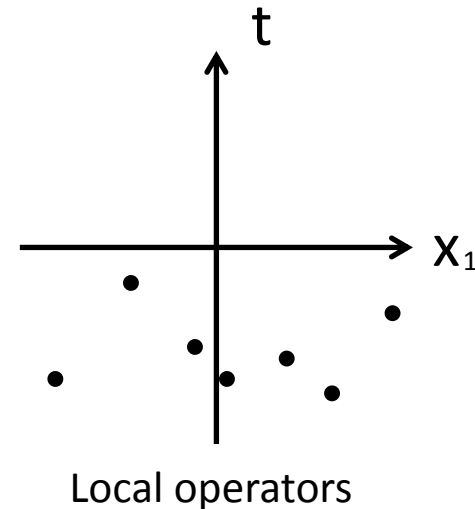
We act various local operators on the ground state.

A locally excited state is given by $|\psi\rangle = \mathcal{N}^{-1} \mathcal{T} \prod_i^k \mathcal{O}^i(t^i, x^{1,i}) |0\rangle$.

Time ordering Op.

The late time of $\Delta S_A^{(n)}$ is given by

$$\Delta S_A^{(n)} = \sum_{i=1}^k \Delta S_A^{(n)i}$$



They are given by the sum of the REE for the state defined by acting each operators $\mathcal{O}^i(t^1, x^{1,i})$ on the ground state.

Summary

- We defined the (Renyi) entanglement entropies of local operators.
 - They characterize local operators from the viewpoint of quantum entanglement.
- These entropies of the operators (constructed of single-species operator) is given by the those of binomial distribution.
 - The results we obtain in terms of entangled pair agree with the results we obtain by replica method.
- They obey the sum rule.

Future Problems

- The formula for the operators constructed of multi-species operators: $:(\partial_r^m \phi)^k \phi^l:$ (generally depend on the spacetime dimension).
- The (Renyi) entanglement entropies of operators in the interacting field theory . (also massive and finite temp.)
- The (Renyi) entanglement entropies for the excited state defined by acting non-local operators.