

## PHYSICS 402 WORKSHEET

We've seen that there is a general relation between symmetries and conserved quantities in quantum mechanics. In this worksheet, we'll use this to understand properties of the momentum operator for a particle in one dimension.

- ① Let  $\hat{T}(a)$  be the unitary operator that tells us how the quantum state transforms when we translate the system to the right by  $a$ . For small  $a$ , let

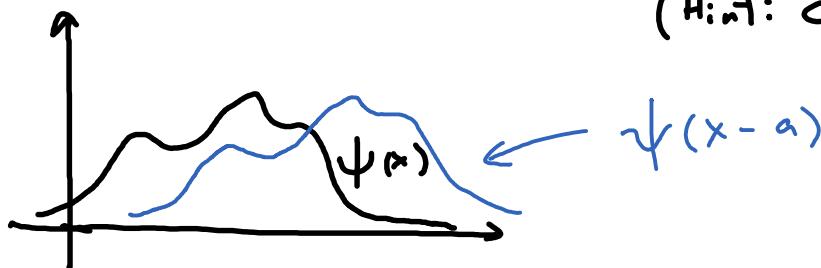
$$\hat{T} = \mathbb{1} - i \frac{a}{\hbar} \hat{P} + \dots \quad (*)$$

We know that  $\hat{P}$  corresponds to a conserved quantity for systems with translation symmetry. We call this the **MOMENTUM** operator.

- a) If the position wavefunction of a state  $|\Psi\rangle$  is

$$\langle x | \Psi \rangle = \psi(x)$$

what is the position wavefunction for  $\hat{T}(a)|\Psi\rangle$ ?  
 (Hint: can you draw it?)



- b) Using this and (\*), what is the wavefunction for the state  $\hat{P}|\Psi\rangle$ ? Have  $\psi(x-a) = \langle x | \hat{T}(a) | \Psi \rangle$

$$\begin{aligned} \text{Taylor expand around } a=0: \Rightarrow \psi(x-a) &= \langle x | (\mathbb{1} - i \frac{a}{\hbar} \hat{P}) | \Psi \rangle \\ &= \psi(x) - i \frac{a}{\hbar} \langle x | \hat{P} | \Psi \rangle + \dots \\ \Rightarrow \langle x | \hat{P} | \Psi \rangle &= \frac{i}{\hbar} \frac{d}{dx} \psi(x) \end{aligned}$$

c) Using your result, complete the following sentence:

"Acting with  $\hat{P}$  on a state corresponds to acting with  $\frac{\hbar}{i} \frac{d}{dx}$  on the wavefunction."

d) What is the wavefunction for  $\hat{X}|\Psi\rangle$ , where  $\hat{X}$  is the position operator whose eigenstates are  $|x\rangle$ ?

Wavefn. is  $\langle x|\hat{X}|\Psi\rangle = x\langle x|\Psi\rangle = x\psi(x)$ . Here we have used that since  $\hat{X}$  is Hermitian, we can act "to the left" on  $\langle x|$  to give  $\langle x|\hat{X} = x\langle x|$ . In more detail, this can be derived as:

$$\langle x|\hat{X}|\Psi\rangle = \langle \Psi|\hat{X}^\dagger|x\rangle^* = \langle \Psi|\hat{X}|x\rangle^* = \langle \Psi|x|x\rangle^* = x\langle x|\Psi\rangle$$

e) Suppose we have a particle in a potential  $V(x)$  so that its classical energy is

$$E = \frac{\hat{P}^2}{2m} + V(x)$$

In the quantum description of this system, we might expect that the energy operator should be:

$$\hat{H} = \frac{1}{2m}\hat{P}^2 + V(\hat{x})$$

In this case, the Schrödinger equation reads

$$i\hbar \frac{d}{dt}|\Psi\rangle = \left( \frac{1}{2m}\hat{P}^2 + V(\hat{x}) \right) |\Psi\rangle$$

By taking the inner product of this with  $\langle x|$  and using your results above, derive the equation for how the particle's wavefunction evolves with time.

$$\text{We get } i\hbar \frac{d}{dt} \langle x|\Psi\rangle = \frac{1}{2m} \langle x|\hat{P}^2|\Psi\rangle + \langle x|V(\hat{x})|\Psi\rangle$$

$$\text{Using } \langle x|\hat{P}|\Psi\rangle = \frac{\pm i}{\hbar} \frac{d}{dx} \langle x|\hat{P}|\Psi\rangle = \left( \frac{\pm i}{\hbar} \frac{d}{dx} \right) \left( \frac{\pm i}{\hbar} \frac{d}{dx} \right) \langle x|\Psi\rangle = -\frac{\hbar^2}{m} \frac{d^2\psi}{dx^2}$$

$$\text{and } \langle x|V(\hat{x})|\Psi\rangle = V(x) \langle x|\Psi\rangle = V(x)\psi(x)$$

$$\text{we get: } \boxed{i\hbar \frac{d}{dt} \psi(x,t) = -\frac{\hbar^2}{2m} \frac{d^2\psi(x,t)}{dx^2} + V(x)\psi(x,t)}$$

f) What is the wavefunction for  $\hat{X}\hat{P}|\Psi\rangle$ ?

It is  $\langle x | \hat{X}\hat{P} | \Psi \rangle = x \langle x | \hat{P} | \Psi \rangle = x \cdot \frac{\hbar}{i} \frac{d\psi}{dx}$

g) Using d) and c), what is the wavefunction for  $\hat{P}\hat{X}|\Psi\rangle$ ?

It is  $\frac{\hbar}{i} \frac{d}{dx} (x\psi(x)) = \frac{\hbar}{i} \psi(x) + \frac{\hbar}{i} x \frac{d\psi}{dx}$

h) Using f) and g), what is the wavefunction for  $[\hat{X}, \hat{P}]|\Psi\rangle$ ? What can you conclude about the operator  $[\hat{X}, \hat{P}]$

We have:  $\langle x | [\hat{X}, \hat{P}] | \Psi \rangle = \langle x | \hat{X}\hat{P} | \Psi \rangle - \langle x | \hat{P}\hat{X} | \Psi \rangle$   
 $= x \frac{\hbar}{i} \frac{d\psi}{dx} - \frac{\hbar}{i} x \frac{d\psi}{dx} - \frac{\hbar}{i} \psi(x)$   
 $= i\hbar \psi(x)$

We conclude  $[\hat{X}, \hat{P}] = i\hbar$

i) What does the generalized uncertainty principle

$$\Delta A \Delta B \geq \left| \langle \psi | \frac{1}{2i} [\hat{A}, \hat{B}] | \psi \rangle \right|$$
 tell us in this case?

We get  $\Delta X \Delta P \geq \left| \langle \psi | \frac{1}{2i} \cdot i\hbar | \psi \rangle \right| = \frac{\hbar}{2}$

j) Can you show that for any state,

$$m \frac{d}{dt} \langle \hat{x} \rangle = \langle \hat{p} \rangle ?$$

We have  $m \frac{d}{dt} \langle \psi | \hat{x} | \psi \rangle$

$$= m \left( \frac{d \langle \psi |}{dt} \right) \hat{x} | \psi \rangle + m \langle \psi | \hat{x} \left( \frac{d}{dt} | \psi \rangle \right)$$

Using  $\frac{d | \psi \rangle}{dt} = -\frac{i}{\hbar} \hat{H} | \psi \rangle$ , we get  $\frac{d \langle \psi |}{dt} = \frac{i}{\hbar} \langle \psi | \hat{H}$

So  $m \frac{d}{dt} \langle \psi | \hat{x} | \psi \rangle$

$$= \frac{mi}{\hbar} \langle \psi | (-\hat{x} - \hat{x} \hat{H}) | \psi \rangle$$

$$= \frac{mi}{\hbar} \langle \psi | \left( \frac{\hat{p}^2}{2m} + V(\hat{x}) \right) \hat{x} - \hat{x} \left( \frac{\hat{p}^2}{2m} + V(\hat{x}) \right) | \psi \rangle$$

Now  $[\hat{x}, V(\hat{x})] = 0$  and  $\hat{P}^2 \hat{x} - \hat{x} \hat{P}^2$

$$= \hat{P}^2 \hat{x} - \hat{P} \hat{x} \hat{P} + \hat{P} \hat{x} \hat{P} - \hat{x} \hat{P}^2$$

$$= -\hat{P} [\hat{x}, \hat{P}] - [\hat{x}, \hat{P}] \hat{P}$$

$$= -2i\hbar \hat{P}$$

Thus  $m \frac{d}{dt} \langle \psi | \hat{x} | \psi \rangle = \langle \psi | \hat{p} | \psi \rangle$