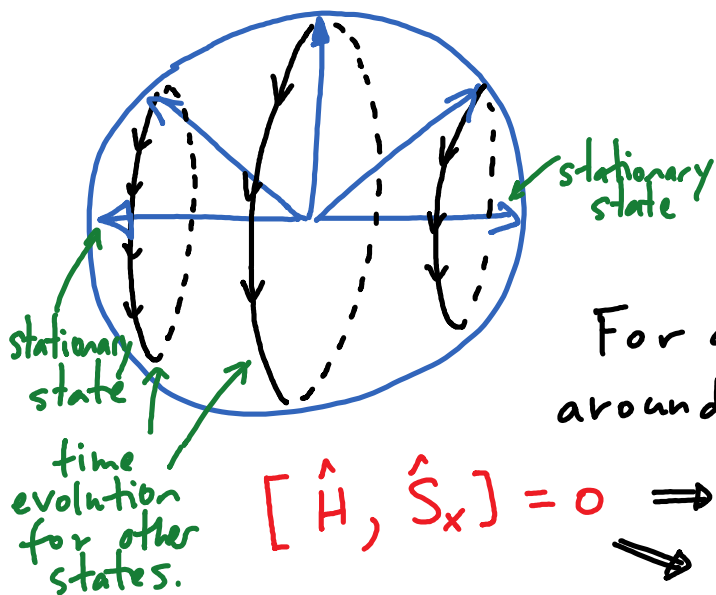


LAST TIME: e.g. time evolution with  $H = E_0 \begin{pmatrix} 3 & 1 \\ 1 & 3 \end{pmatrix}$



e.g. electron in magnetic field in  $\hat{x}$  direction + constant potential

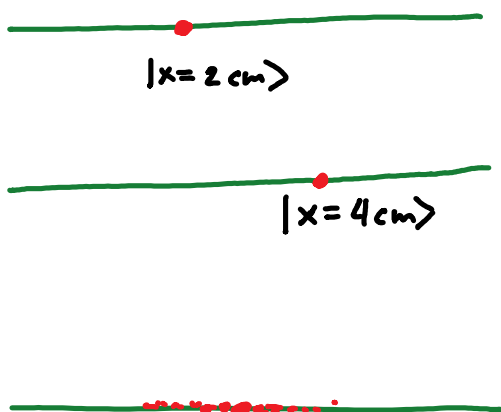
$$\hat{H} = 3E_0 \cdot \mathbb{1} + \frac{2E_0}{\hbar} \hat{S}_x$$

For any initial state,  $|\Psi(t)\rangle$  rotates around x axis.

$$[\hat{H}, \hat{S}_x] = 0 \Rightarrow \text{rotational symmetry about x axis}$$

$$\Rightarrow \hat{S}_x \text{ conserved.}$$

1D particle system:  $|x\rangle$ : eigenstate of position operator  $\hat{X}$



General state:

$$|\Psi\rangle = \int_{-\infty}^{\infty} dx \psi(x) |x\rangle$$

coeff. in superposition

$$= \text{WAVEFUNCTION } \psi(x) = \langle x | \Psi \rangle$$

Q: What is  $\langle \Psi_2 | \Psi_1 \rangle$  for states w. wavefns.  $\psi_1(x), \psi_2(x)$ ?

$$A: \langle \Psi_2 | \Psi_1 \rangle = \int_{-\infty}^{\infty} dx \psi_2(x) \langle \Psi_2 | x \rangle = \int_{-\infty}^{\infty} dx \psi_2^*(x) \psi_1(x)$$

$$\langle \Psi | \Psi \rangle = \int_{-\infty}^{\infty} dx |\psi(x)|^2$$

Hilbert space: all states where this is  $< \infty$   
(mathematicians call this  $L_2$ )

$\hat{T}(a)$ : translates by  $a$  to the right

Infinitesimal:  $\hat{T}(\epsilon) = \mathbb{1} - i \frac{\epsilon}{\hbar} \hat{P}$   $\hat{P}$  MOMENTUM OPERATOR

Wavefunction for  $\hat{T}(\epsilon)|\Psi\rangle$  is  $\psi(x-\epsilon)$

$$\langle x | (\mathbb{1} - i \frac{\epsilon}{\hbar} \hat{P}) |\Psi\rangle = \psi(x) - \epsilon \psi'(x) + \dots$$

$$O(\epsilon): \quad \langle x | \hat{P} |\Psi\rangle = \frac{\hbar}{i} \frac{d}{dx} \psi'(x)$$

Acting with  $\hat{P}$  on state  $\longleftrightarrow$  acting with  $\frac{\hbar}{i} \frac{d}{dx}$  on wave function

Q: A state has wavefunction  $\psi(x)$ . What is the expectation value for a measurement of momentum?

$$A: \langle P \rangle = \langle \Psi | \hat{P} | \Psi \rangle = \int_{-\infty}^{\infty} dx \psi^*(x) \cdot \frac{\hbar}{i} \psi'(x)$$

Wavefn for  $\hat{X}|\Psi\rangle$  is  $\langle x | \hat{X} | \Psi \rangle = x \langle x | \Psi \rangle = x \psi(x)$

Acting with  $\hat{X}$  on state  $\longleftrightarrow$  multiplying wavefn. by  $x$

$$\text{CONSEQUENCE: } [\hat{X}, \hat{P}] = i\hbar$$

By generalized uncertainty principle:

$$\Delta x \Delta p \geq \left| \langle \frac{1}{2i} [\hat{X}, \hat{P}] \rangle \right| = \frac{\hbar}{2} \quad \text{HEISENBERG UNCERTAINTY PRINCIPLE}$$

Momentum operator also has eigenstates  $|p\rangle$  with

$$\hat{P}|p\rangle = p|p\rangle \quad \langle p|q\rangle = \delta(p-q)$$

Can represent any state as

$$|\Psi\rangle = \int_{-\infty}^{\infty} dp \psi(p) |p\rangle$$

where  $\psi(p) = \langle p|\Psi\rangle$  is the MOMENTUM WAVEFUNCTION

HOMEWORK: relation between  $\psi(p)$  and  $\psi(x)$ ,  
wavefunction for  $|p\rangle$

example: time evolution in 1D system.

For classical 1D system, energy is

$$E = \frac{p^2}{2m} + V(x) \quad (*)$$

In quantum version, energy operator should be

$$\hat{H} = \frac{\hat{p}^2}{2m} + V(\hat{x})$$

guarantees that (\*) holds for expectation values

$\therefore$  Time evolution is

$$i\hbar \frac{d}{dt} |\psi\rangle = \hat{H} |\psi\rangle$$

$$\Rightarrow i\hbar \frac{d}{dt} \langle x | \psi \rangle = \langle x | \frac{\hat{p}^2}{2m} + V(\hat{x}) | \psi \rangle$$

$$= \frac{1}{2m} \langle x | \hat{p} \cdot \hat{p} | \psi \rangle + \langle x | V(x) | \psi \rangle$$

$$= \frac{1}{2m} \cdot \left( \frac{\hbar}{i} \frac{d}{dx} \langle x | \hat{p} | \psi \rangle \right) + \langle x | V(x) | \psi \rangle$$

$$= \frac{1}{2m} \left( \frac{\hbar}{i} \frac{d}{dx} \right) \left( \frac{\hbar}{i} \frac{d}{dx} \right) \langle x | \psi \rangle + V(x) \langle x | \psi \rangle$$

$$\Rightarrow i\hbar \frac{\partial \psi}{\partial t} = -\frac{\hbar^2}{2m} \frac{\partial^2 \psi}{\partial x^2} + V(x) \psi(x)$$

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