

# LAST TIME: Time dependent perturbation theory

$$H = H_0 + H'(t)$$

time  
indep

small time-dep.  
perturbation

$H_0$  eigenstates.

$$|\psi(t)\rangle = \sum_n c_n(t) e^{-\frac{iE_n t}{\hbar}} |\psi_n\rangle$$

$$\text{Found } c_m(t) = c_m(0) - \frac{i}{\hbar} \sum_n \left[ \int_0^t dt' e^{i\omega_{mn}t'} H'_{mn}(t') \right] c_n(0) + \dots$$

$$\omega_{mn} = \frac{E_m - E_n}{\hbar}$$

$$H'_{mn} = \langle m | H'(t) | n \rangle$$

TRANSITION PROBABILITIES:

$$\text{if } |\psi(0)\rangle = |\psi_a\rangle$$

Probability of finding  $|\psi_b\rangle$  at time  $t$ :

$$P_b(t) = |c_b(t)|^2 = \frac{1}{\hbar^2} \left| \int_0^t dt_1 e^{i\omega_{ba}t_1} H'_{ba}(t_1) \right|^2$$

Total probability of making a transition:

$$P_{\text{tot}} = \sum_{b \neq a} |c_b(t)|^2$$

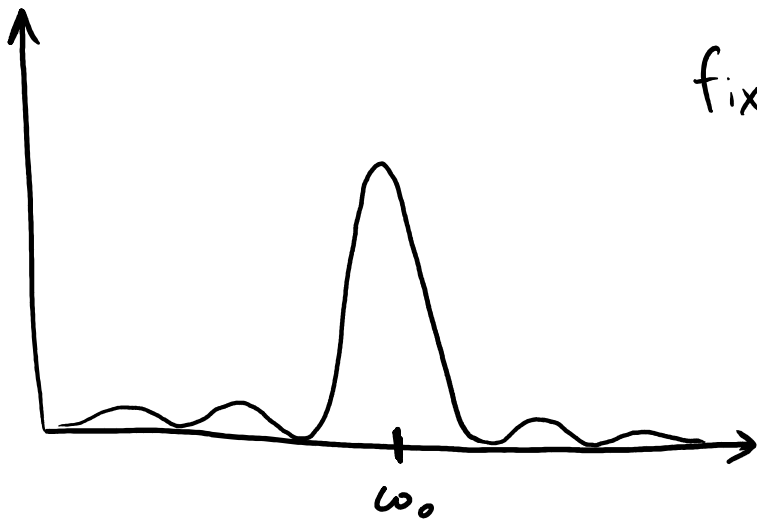
e.g. sinusoidal perturbations

$$H'(t) = V \cos(\omega t)$$

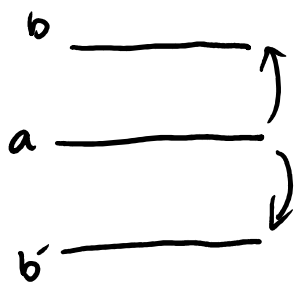
$$P_b(t) = \frac{1}{\hbar^2} \left| \int_0^t dt' e^{i\omega_{ba}t'} \cos(\omega t') V_{ba} \right|^2$$
$$= \frac{|V_{ba}|^2}{\hbar^2} \left| \frac{e^{i(\omega_{ba}+\omega)t} - 1}{2(\omega_{ba}+\omega)} + \frac{e^{i(\omega_{ba}-\omega)t} - 1}{2(\omega_{ba}-\omega)} \right|^2$$

Most significant near  $\omega \approx \omega_{ba} = \frac{|E_b - E_a|}{\hbar} \equiv \omega_0$

Then 
$$P_b(t) \approx \frac{|V_{ba}|^2}{\hbar^2} \frac{\sin^2\left(\frac{(\omega - \omega_0)t}{2}\right)}{(\omega - \omega_0)^2}$$



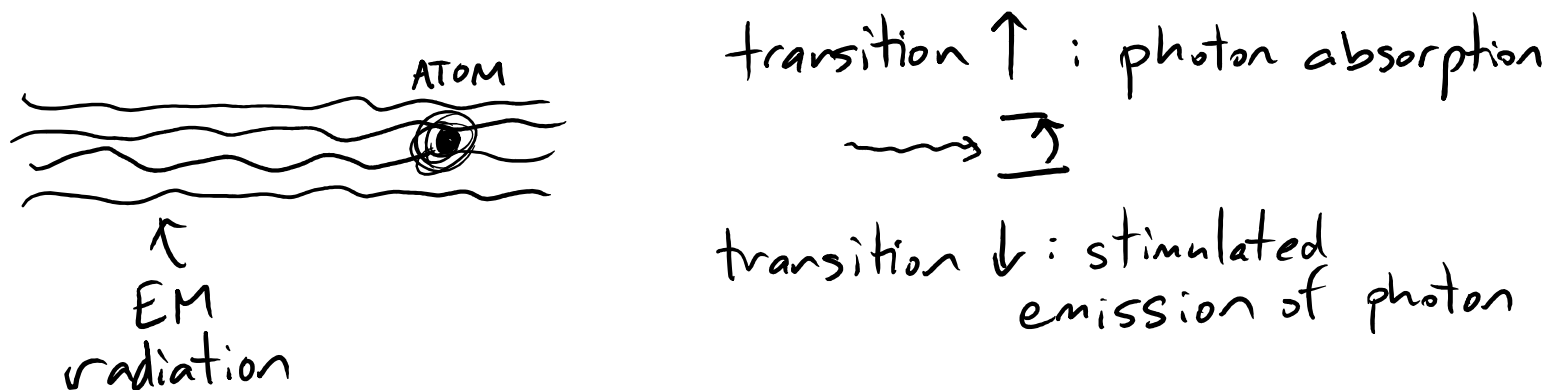
fixed  $t$ : transition probability  
highly peaked when  
driving frequency matches  
 $\frac{|E_a - E_b|}{\hbar}$



can stimulate transitions up (system absorbs energy from external perturbation) or down (system emits energy).

$P_b(t)$  looks sinusoidal with  $t$ , but this is an artifact of assuming a single  $\omega$ .

Application: atomic/molecular transitions.



Also: can have spontaneous emission



Need Hamiltonian for charged particle in EM field.

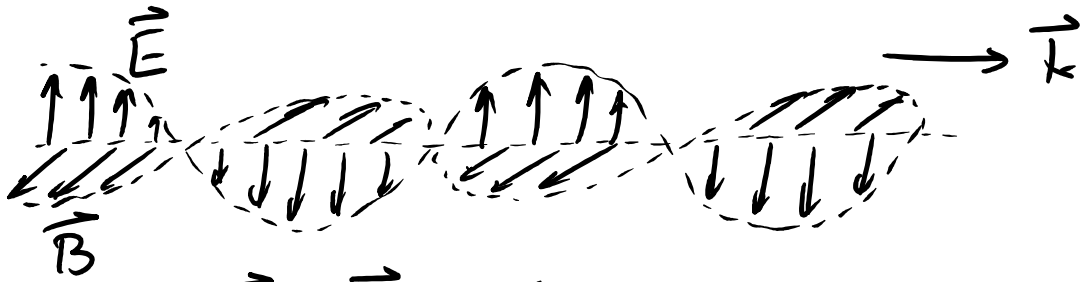
Recall: can write  $\vec{B} = \nabla \times \vec{A}$

$$\vec{E} = -\frac{\partial \vec{A}}{\partial t} - \nabla \phi$$

"Gauge transformation"  $\phi \rightarrow \phi - \partial_t \lambda$   $\vec{A} \rightarrow \vec{A} - \nabla \lambda$  doesn't change  $\vec{E}, \vec{B}$ . Useful to pick  $\nabla \cdot \vec{A} = 0$  "Coulomb Gauge". Then:

$$H = \frac{p^2}{2m} - \frac{q}{m} \vec{p} \cdot \vec{A}(\vec{x}, t) + \frac{q^2}{2m} \vec{A}^2(\vec{x}, t) + q \phi(\vec{x}, t)$$

Monochromatic EM wave:



$$\vec{E} = \vec{E}_0 \cos(\vec{k} \cdot \vec{x} - \omega t) \quad |\vec{B}_0| = \frac{1}{c} |\vec{E}_0|$$

$$\vec{B} = \vec{B}_0 \cos(\vec{k} \cdot \vec{x} - \omega t)$$

Can choose  $\phi = 0$   $\vec{A} = -\frac{1}{\omega} \vec{E}_0 \sin(\vec{k} \cdot \vec{x} - \omega t)$

for light:

$$\lambda \sim 500 \text{ nm} \gg \text{atom size}$$

$$\therefore \vec{k} \cdot \vec{x} \text{ small} \quad (k) = \frac{2\pi}{\lambda}$$

$$\phi = 0 \quad \vec{A} \approx \frac{1}{\omega} \vec{E}_0 \sin(\omega t)$$

Simplify via gauge transformation with  $\lambda = \frac{\vec{E}_0 \cdot \vec{x}}{\omega} \sin(\omega t)$

$$\phi_{\text{new}} = -\vec{E}_0 \cdot \vec{x} \cos(\omega t) \quad \vec{A}_{\text{new}} \approx 0$$

+ small terms we ignored.

Then:  $H_{EM} \approx e \vec{E}_0 \cdot \vec{x} \cos(\omega t)$

$$= -\vec{E}_0 \cdot \vec{p} \cos(\omega t)$$

DIPOLE APPROXIMATION



electric dipole moment operator  
for atom:  $\vec{p} = -e \vec{x}$