

Consider states

$$|n \ell m\rangle \otimes |s_z\rangle$$

of a hydrogen atom, with  $\ell=1$  and  $n=2$ .

We'll denote these by  $|m\rangle \otimes |s_z\rangle$ .

a) Write each of the possible basis states, and for each state, write the eigenvalue of

$$J_z = L_z + S_z$$

state:

$$|1\rangle \otimes |+\frac{1}{2}\rangle$$

$$|1\rangle \otimes |-\frac{1}{2}\rangle$$

$$|0\rangle \otimes |+\frac{1}{2}\rangle$$

$$|0\rangle \otimes |-\frac{1}{2}\rangle$$

$$|-1\rangle \otimes |+\frac{1}{2}\rangle$$

$$|-1\rangle \otimes |-\frac{1}{2}\rangle$$

$$J_z$$

$$\frac{3}{2}$$

$$\frac{1}{2}$$

$$\frac{1}{2}$$

$$-\frac{1}{2}$$

$$-\frac{1}{2}$$

$$-\frac{3}{2}$$

b) These states aren't eigenstates of total angular momentum  $J^2$ . But some linear combinations of them are eigenstates of  $J^2$  and  $J_z$  (label these by  $|J M\rangle$ ). Based on the values of  $J_z$  you found, and the fact that states with some  $J$  always come in groups with  $M = -J, -J+1, \dots, J$ , what are the possible values of  $(J, M)$ ?

Have groups  $-\frac{3}{2}, -\frac{1}{2}, \frac{1}{2}, \frac{3}{2}$  and  $-\frac{1}{2}, \frac{1}{2}$  for  $J_z$ ,  
corresponding to  $J = \frac{3}{2}$  and  $J = \frac{1}{2}$

c) Each  $|J M\rangle$  state is a linear combination of  $|m\rangle \otimes |s_z\rangle$  states, but two of the states are just equal to a single  $|m\rangle \otimes |s_z\rangle$  state. Which are they?

$$|\frac{3}{2} \frac{3}{2}\rangle = |1\rangle \otimes |\frac{1}{2}\rangle$$

$$|\frac{3}{2} -\frac{3}{2}\rangle = |-1\rangle \otimes |-\frac{1}{2}\rangle$$

since no other states have  $m+s_z = \pm\frac{3}{2}$

d) Starting from the one with the larger  $M$ ,

act on the equation  $|J M\rangle = |m\rangle \otimes |s_z\rangle$

with  $J_- = L_- + S_-$  (use  $J_-$  on the left and  $L_- + S_-$  on the right) to determine how to write the other

$|J M\rangle$  states for this  $J$ .

$$J_- |\frac{3}{2} \frac{3}{2}\rangle = \sqrt{\frac{3}{2} \cdot \frac{5}{2} - \frac{3}{2} \cdot \frac{1}{2}} |\frac{3}{2} \frac{1}{2}\rangle = \sqrt{3} |\frac{3}{2} \frac{1}{2}\rangle$$

also  $J_- = L_- + S_-$  and  $|\frac{3}{2} \frac{3}{2}\rangle = |1\rangle \otimes |\frac{1}{2}\rangle$  so

$$\begin{aligned} \sqrt{3} |\frac{3}{2} \frac{1}{2}\rangle &= (L_- + S_-) |1\rangle \otimes |\frac{1}{2}\rangle \\ &= (L_- |1\rangle) \otimes |\frac{1}{2}\rangle + |1\rangle \otimes (S_- |\frac{1}{2}\rangle) \\ &= \sqrt{2} |0\rangle \otimes |\frac{1}{2}\rangle + |1\rangle \otimes |-\frac{1}{2}\rangle \end{aligned}$$

$$\text{Thus: } |\frac{3}{2} \frac{1}{2}\rangle = \sqrt{\frac{2}{3}} |0\rangle \otimes |\frac{1}{2}\rangle + \sqrt{\frac{1}{3}} |1\rangle \otimes |-\frac{1}{2}\rangle$$

Similarly, we get  $J_- |\frac{3}{2} \frac{1}{2}\rangle = 2 |\frac{3}{2} -\frac{1}{2}\rangle$  so

$$\begin{aligned} 2 |\frac{3}{2} -\frac{1}{2}\rangle &= \sqrt{\frac{2}{3}} (L_- |0\rangle) \otimes |\frac{1}{2}\rangle + \sqrt{\frac{2}{3}} |0\rangle \otimes (S_- |\frac{1}{2}\rangle) + \sqrt{\frac{1}{3}} (L_- |1\rangle) \otimes |-\frac{1}{2}\rangle \\ &= \frac{2}{\sqrt{3}} |1\rangle \otimes |\frac{1}{2}\rangle + \sqrt{\frac{2}{3}} |0\rangle \otimes |-\frac{1}{2}\rangle + \sqrt{\frac{2}{3}} |0\rangle \otimes |-\frac{1}{2}\rangle \end{aligned}$$

$$\Rightarrow |\frac{3}{2} -\frac{1}{2}\rangle = \frac{1}{\sqrt{3}} |-1\rangle \otimes |\frac{1}{2}\rangle + \sqrt{\frac{2}{3}} |0\rangle \otimes |-\frac{1}{2}\rangle$$

Finally, we already have  $|\frac{3}{2} -\frac{3}{2}\rangle = |-1\rangle \otimes |-\frac{1}{2}\rangle$

e) Can you find the remaining  $|J M\rangle$  states using the fact that they must be orthogonal to the ones you found?

We know that  $|\frac{1}{2} \frac{1}{2}\rangle = A|1\rangle \otimes |-\frac{1}{2}\rangle + B|0\rangle \otimes |\frac{1}{2}\rangle$   
 since the basis states on the right are the only ones with  $L_z + S_z = \frac{1}{2}$ .  
 $|\frac{1}{2} \frac{1}{2}\rangle$  should also be orthogonal to

$$|\frac{3}{2} \frac{1}{2}\rangle = \sqrt{\frac{2}{3}}|0\rangle \otimes |\frac{1}{2}\rangle + \sqrt{\frac{1}{3}}|1\rangle \otimes |-\frac{1}{2}\rangle$$

and we can assume  $|A|^2 + |B|^2 = 1$  for normalization, so (up to a phase),

$$|\frac{1}{2} \frac{1}{2}\rangle = -\sqrt{\frac{1}{3}}|0\rangle \otimes |\frac{1}{2}\rangle + \sqrt{\frac{2}{3}}|1\rangle \otimes |-\frac{1}{2}\rangle$$

Acting with  $J^- = L^- + S^-$ , we find

$$J^- |\frac{1}{2} \frac{1}{2}\rangle = -\sqrt{\frac{1}{3}} L^- |0\rangle \otimes |\frac{1}{2}\rangle - \sqrt{\frac{1}{3}} |0\rangle \otimes S^- |\frac{1}{2}\rangle + \sqrt{\frac{2}{3}} L^- |1\rangle \otimes |-\frac{1}{2}\rangle$$

$$\Rightarrow |\frac{1}{2} -\frac{1}{2}\rangle = \frac{1}{\sqrt{3}}|0\rangle \otimes |-\frac{1}{2}\rangle - \sqrt{\frac{2}{3}}|1\rangle \otimes |\frac{1}{2}\rangle$$

We can also check that this is orthogonal to  $|\frac{3}{2} -\frac{1}{2}\rangle$ .