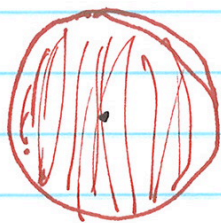


Lagrangian Mechanics

Kepler Problem

Not to Scale!



Star
 m_1



Planet
 m_2

- a) What is $U(r)$? $U(r) = -\frac{Gm_1m_2}{r} = -\frac{\gamma}{r}$
- b) Lagrange equations
- c) $U_{\text{eff}}(r) \rightarrow$ sketch Equivalent 1D problem
- d) What kind of motion allowed for fixed E ?

$$\mu \ddot{r} = -\frac{\partial U_{\text{eff}}}{\partial r}$$

Lagrangian Mechanics

Kepler Problem

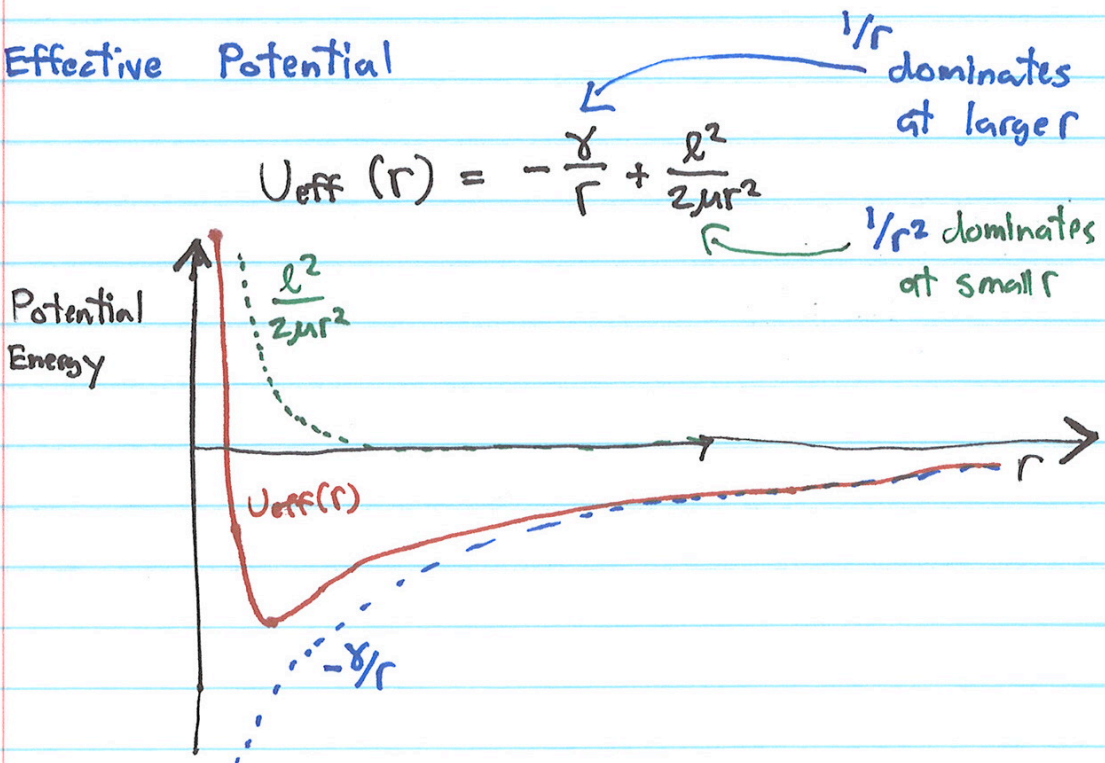
$$U(r) = -\frac{Gm_1 m_2}{r} = -\frac{\gamma}{r}$$

$$\text{So } L = T - U$$

$$L = \frac{1}{2}\mu(\dot{r}^2 + r^2\dot{\phi}^2) + \frac{\gamma}{r}$$

$$\begin{aligned} \phi: \quad l &= \mu r^2 \dot{\phi} & \frac{\partial L}{\partial r} &= \mu r \dot{\phi}^2 - \frac{\gamma}{r^2} \\ & & &= \mu r \left(\frac{l}{\mu r^2}\right)^2 - \frac{\gamma}{r^2} \\ r: \quad \mu \ddot{r} &= \frac{l^2}{\mu r^3} - \frac{\gamma}{r^2} \end{aligned}$$

Effective Potential



Kepler Problem

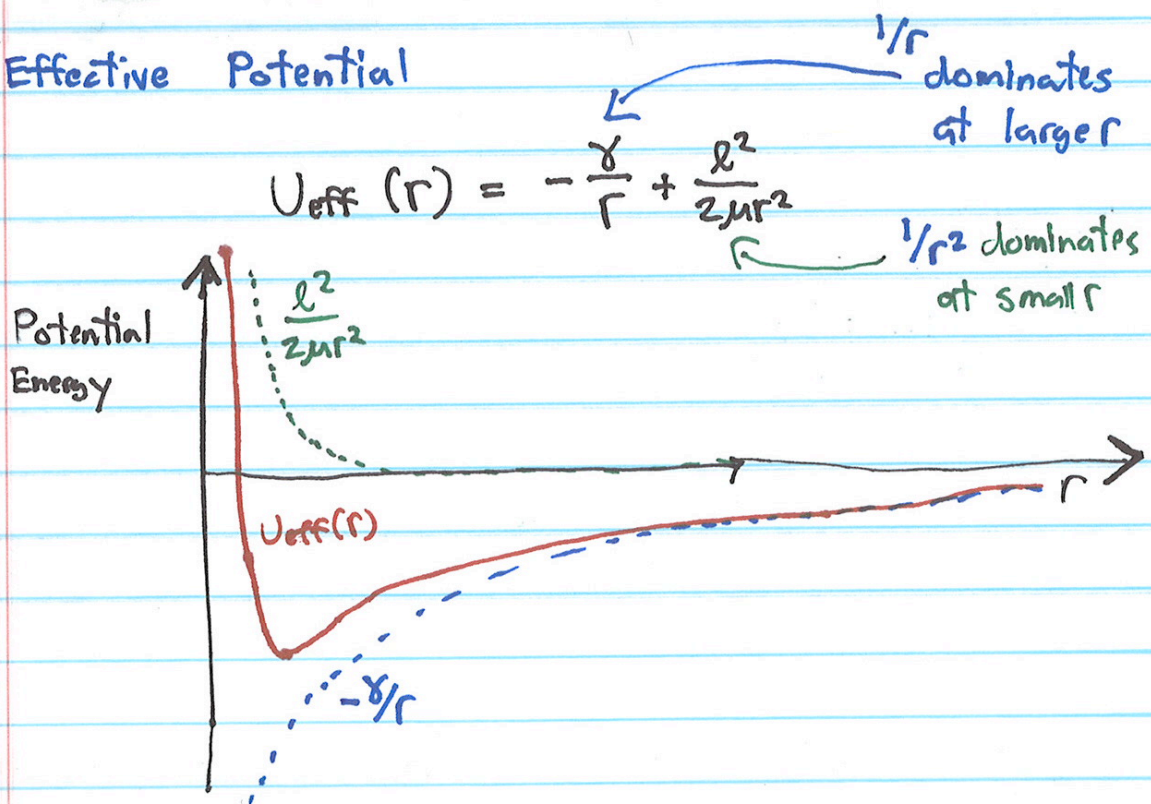
For a potential $U(r) = -\gamma/r$

The effective potential at fixed angular momentum l is:

$$U_{\text{eff}}(r) = \frac{l^2}{2\mu r^2} + U(r)$$

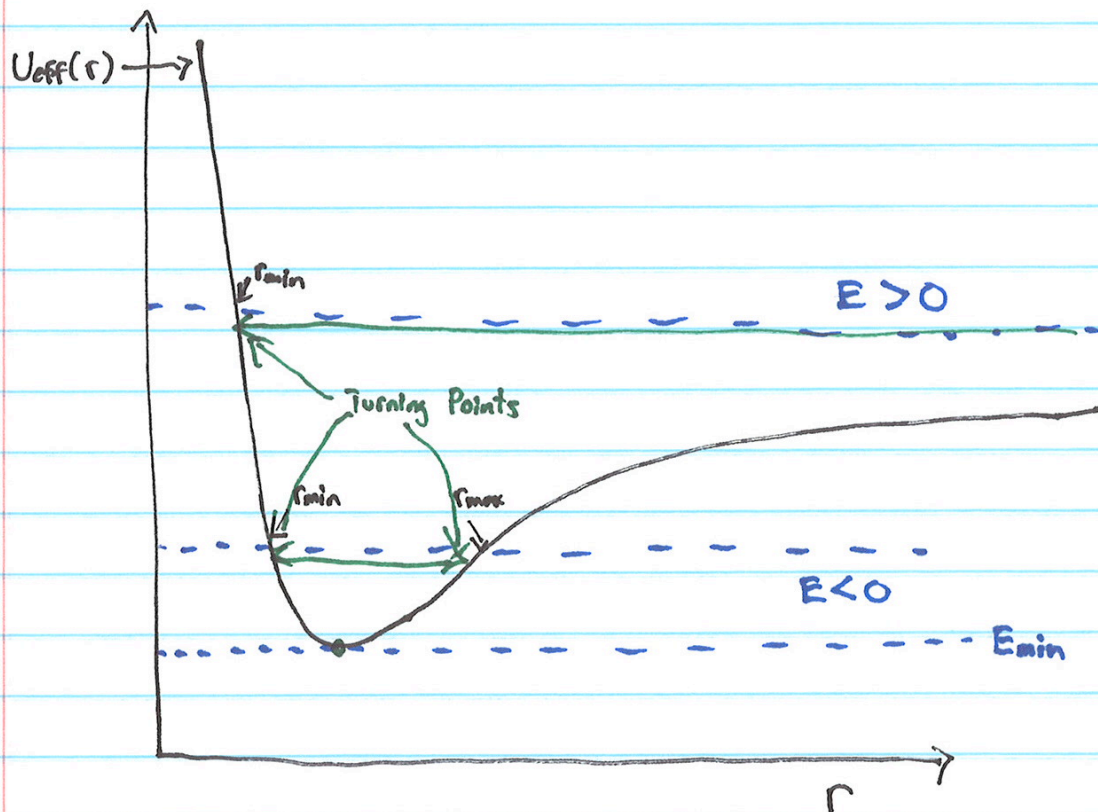
$$U_{\text{eff}}(r) = \frac{l^2}{2\mu r^2} - \frac{\gamma}{r}$$

Effective Potential



Lagrangian Mechanics

Kepler Problem



Unbound $\infty > r > r_{\min}$

Bound $r_{\max} > r > r_{\min}$

Solve $E = U_{\text{eff}}(r_{\min})$ or $E = U_{\text{eff}}(r_{\max})$
to find r_{\min}, r_{\max}

Kepler Problem

The radial equation is

$$\mu \frac{d^2 r}{dt^2} = \mu \ddot{r} = - \frac{\partial U_{\text{eff}}}{\partial r} = - \frac{\gamma}{r^2} + \frac{l^2}{2\mu r^3}$$

solve for $r(t)$

Gravity



↑
Angular Momentum
Conservation

We can use the fact that $l = r^2 \frac{d\phi}{dt} \mu$
to write $dt = \frac{r^2 \mu}{l} d\phi$

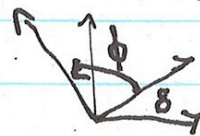
$$\Rightarrow \frac{d}{dt} = \frac{l}{r^2 \mu} \frac{d}{d\phi} = \frac{l}{\mu} u^2 \frac{d}{d\phi}$$

If we define $u = 1/r$ then the radial equation becomes

$$\mu \cdot \frac{l}{\mu} u^2 \frac{d}{d\phi} \left(\frac{l}{\mu} u^2 \frac{d}{d\phi} (1/u) \right) = -\gamma u^2 + \frac{l^2}{2\mu} u^3$$

$$\Rightarrow \frac{d^2 u}{d\phi^2} + u = \frac{\gamma \mu}{l^2} \quad \text{"shifted Harmonic Oscillator Equation"}$$

$$\Rightarrow u(\phi) = \frac{\gamma \mu}{l^2} + A \cos(\phi + \delta)$$



We can always set $\delta = 0$ by redefining $\phi = 0$

Lagrangian Mechanics

Kepler Problem: The Orbit

$$\frac{d^2 U}{d\phi^2} + U(\phi) = \frac{\gamma\mu}{l^2}$$

$$\Rightarrow \frac{d^2}{d\phi^2} \left(U - \frac{\gamma\mu}{l^2} \right) + \left(U - \frac{\gamma\mu}{l^2} \right) = 0$$

$$\Rightarrow U(\phi) = \frac{\gamma\mu}{l^2} + A \cos(\phi + \delta)$$

Can choose $\delta = 0$ by redefining $\phi = 0$

$$\therefore U(\phi) = \frac{\gamma\mu}{l^2} (1 + \epsilon \cos \phi)$$

$$r(\phi) = \frac{c}{1 + \epsilon \cos \phi}$$

$$c = \frac{l^2}{\gamma\mu}$$

$$\epsilon < 1$$

ellipse

$$\epsilon = 1$$

parabola

$$\epsilon > 1$$

hyperbola

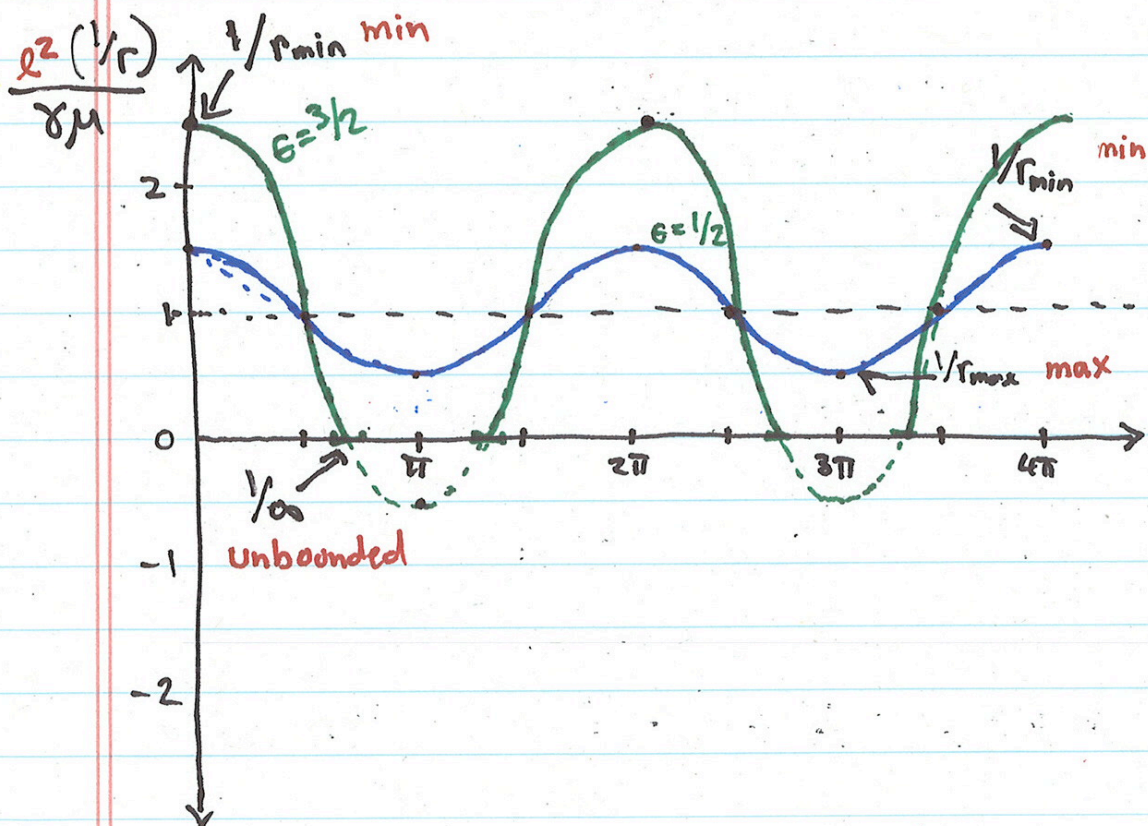
Kepler Problem

$$\text{So } U(\phi) = \frac{\gamma\mu}{l^2} + A \cos \phi$$

$$\text{or } U(\phi) = \frac{\gamma\mu}{l^2} (1 + E \cos \phi)$$

eccentricity

$$\therefore r(\phi) = \frac{(l^2/\gamma\mu)}{1 + E \cos \phi} = \frac{C}{1 + E \cos \phi}$$



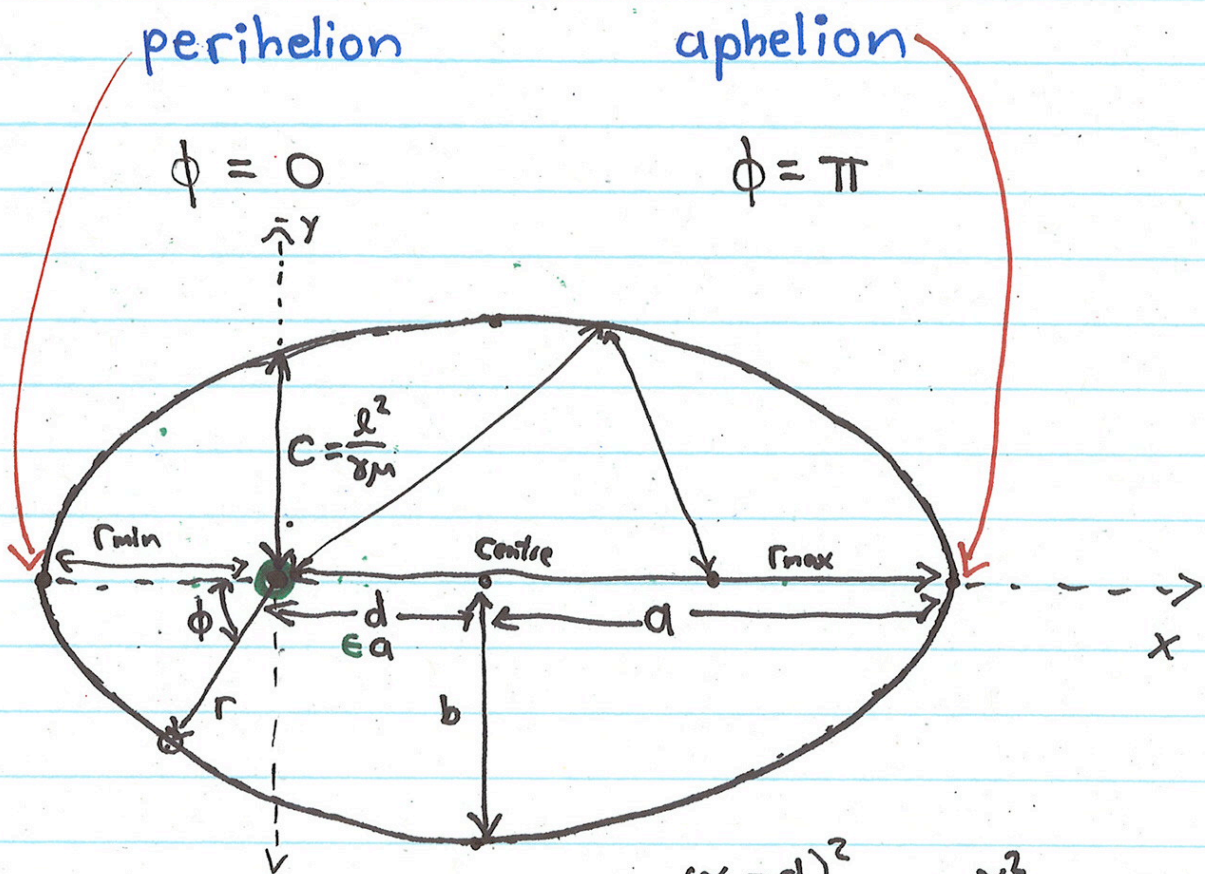
Kepler Problem

Bounded Orbits $0 \leq \epsilon < 1$

In this case the particle/planet oscillates between

Note: Circle
if $\epsilon = 0$

$$r_{\min} = \frac{c}{1+\epsilon} \quad \text{and} \quad r_{\max} = \frac{c}{1-\epsilon}$$



$$r(\phi) = \frac{c}{1 + \epsilon \cos \phi}$$

\Leftrightarrow

$$\frac{(x-d)^2}{a^2} + \frac{y^2}{b^2} = 1$$

ellipse with FOCUS at origin