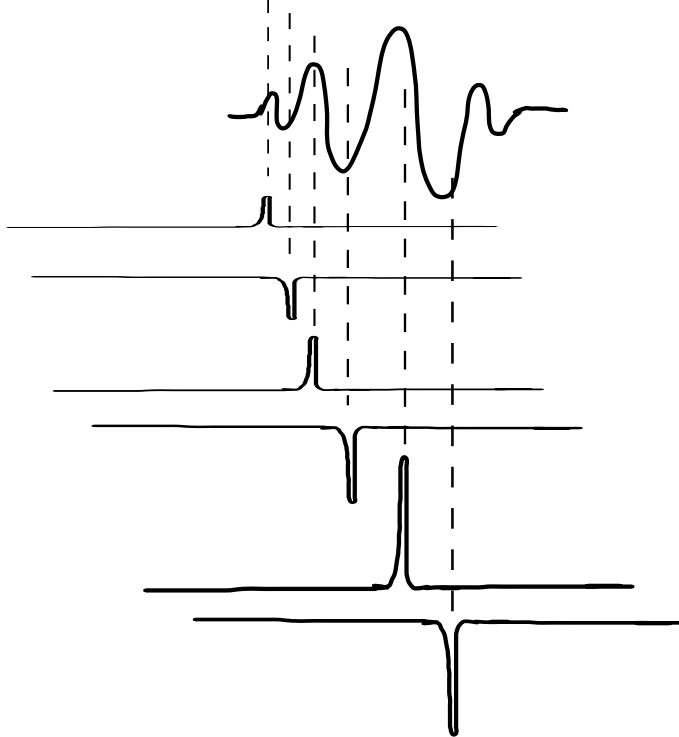
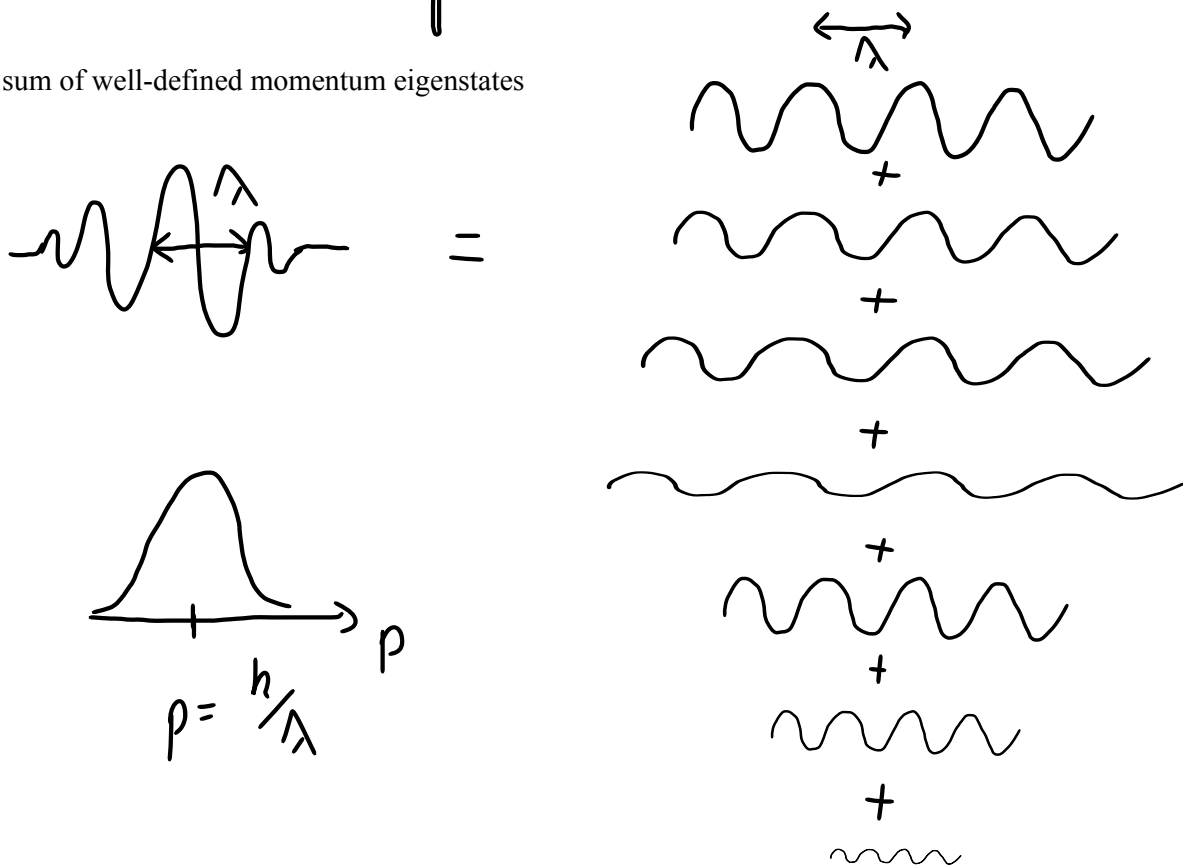


Last time we talked about the properties of wavepackets: these are wavefunctions for electrons which have a fairly well defined (but not perfectly defined) position and momentum. We can think of a wavepacket in two ways:

As a sum of well-defined position eigenstates...



or, as a sum of well-defined momentum eigenstates



The narrower wavepacket you want, the wider the required momentum amplitude function.
The narrower the momentum amplitude function, the wider the wavepacket.

$$\psi(x) = e^{i \frac{2\pi}{h} p x}$$

Perfectly well defined momentum (momentum eigenstate) = planewave, spread over all positions equally $|\psi(x)|^2 = 1$

Perfectly well defined position (position eigenstate) = superposition of all momentum eigenstates with equal weight.

$$|\hat{\psi}(p)|^2 = 1$$

	Polarizer at 0°	Polarizer at 45°
eigenstates	$ 0^\circ\rangle, 90^\circ\rangle$	$ 45^\circ\rangle, -45^\circ\rangle$
general quantum state	$a_0 0^\circ\rangle + a_{90} 90^\circ\rangle$	$a_{45} 45^\circ\rangle + a_{-45} -45^\circ\rangle$
	$a_0 = \frac{a_{45} + a_{-45}}{\sqrt{2}}$ $a_{90} = \frac{a_{45} - a_{-45}}{\sqrt{2}}$	$a_{45} = \frac{a_0 + a_{90}}{\sqrt{2}}$ $a_{-45} = \frac{a_0 - a_{90}}{\sqrt{2}}$

The better we can predict the outcome for one polarizer, the less we know about the outcome for the other polarizer.

	Position	Momentum
eigenstates	$ x\rangle$, for any x	$ p\rangle$, for any p
general quantum state	$\int \Psi(x) x\rangle dx$	$\int \tilde{\Psi}(p) p\rangle dp$
	$\Psi(x) =$ $\int \tilde{\Psi}(p) e^{i\frac{2\pi}{h}px} dp$	$\tilde{\Psi}(p) =$ $\frac{1}{2\pi} \int \Psi(x) e^{-i\frac{2\pi}{h}px} dx$

The better we can predict the outcome of a position measurement, the worse we can predict the outcome of a momentum measurement, and vice versa.