

| Equation of State |
|---------------------|
| Mean-Field Models |
| Nuclear Observables |
| Nuclear Collisions |
| The Symmetry Energy |
| Neutron Stars |



Neutron Rich Nuclei in Heaven and Earth

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Equation of State I: Generalities

The Bethe-Weizsäcker (BW) Mass Formula:

 $E(Z,N) = a_{\rm vol}A + a_{\rm surf}A^{2/3} + a_{\rm coul}Z^2/A^{1/3} + a_{\rm symm}(N-Z)^2/A + \dots$

- Parameters extracted from a fit to thousands of known nuclear masses
- Hidden behind its success is the saturation of the nuclear force
- BW constrains the above parameters at (or near) saturation density: $(a_{\text{vol}}, a_{\text{surf}}, a_{\text{coul}}, a_{\text{symm}}) \simeq (-16, +18, +1, +26) \text{ MeV}$



BW offers little on the density dependence of the parameters!



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|---------------------|
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| Nuclear Observables |
| Nuclear Collisions |
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Page 3 of 12

Equation of State II: Infinite Nuclear Matter

Recipe to make infinite nuclear matter:

- Turn off the long-range Coulomb force
- Let Z, N and V go to infinity with ratios remaining finite: $\rho = A/V, Y_p = Z/A, b = \delta = (N-Z)/A, \ldots$
- Only surviving terms in the thermodynamic limit: $E(Z,N)/A = a_{\rm vol} + a_{\rm symm}b^2 = \epsilon_0 + Jb^2$

Symmetric vs Asymmetric Matter:

• Expand the total energy per nucleon around b=0:

$$E(\rho; b)/A = \underbrace{E(\rho; b=0)/A}_{\text{Symmetric Matter}} + b \underbrace{\left(\frac{\partial E/A}{\partial b}\right)_{b=0}}_{0} + b^2 \underbrace{\frac{1}{2} \left(\frac{\partial^2 E/A}{\partial b^2}\right)_{b=0}}_{\text{Symmetry Energy}} + \dots$$

Pure neutron matter \approx Symmetric Matter + Symmetry Energy!



Goal: Study the density dependence of the equation of state



Equation of State Mean-Field Models Nuclear Observables Nuclear Collisions The Symmetry Energy Neutron Stars



Density Functional, Kohn-Sham, MF Theory

Improving the Bethe-Weizsäcker (BW) Mass Formula:

$$\mathcal{L}_{\rm int} = g_{\rm s} \bar{\psi} \psi \phi - g_{\rm v} \bar{\psi} \gamma^{\mu} \psi V_{\mu} - \frac{g_{\rho}}{2} \bar{\psi} \gamma^{\mu} \boldsymbol{\tau} \cdot \mathbf{b}_{\mu} \psi - e \bar{\psi} \gamma^{\mu} \tau_{p} \psi A_{\mu} - \frac{\kappa}{3!} (g_{\rm s} \phi)^{3} - \frac{\lambda}{4!} (g_{\rm s} \phi)^{4} + \Lambda_{\rm v} (g_{\rm v}^{2} V^{\mu} V_{\mu}) (g_{\rho}^{2} b^{\mu} b_{\mu}) + \frac{\zeta}{4!} g_{\rm v}^{4} (V_{\mu} V^{\mu})^{2}$$

- Parameters fitted to a large body of ground-state properties (mostly binding energies and charge radii of many nuclei)
- Ground-state observables computed at the mean-field level
- Formalism is **NOT** Hartree (Hartree-Fock) theory
- Parameters of the model encode correlations that go beyond two-body (short, long, and pairing correlations in an average way)

Resulting model unlikely to describe correctly NN physics!

Correlating Model Parameters to the Physics:

| Parameters | Constrained by |
|---------------------|--|
| $g_{ m s},g_{ m v}$ | Ground state properties of finite nuclei |
| $g_{ ho}$ | Ground state properties of heavy nuclei |
| κ,λ | Isoscalar giant monopole resonance |
| $\Lambda_{ m v}$ | Neutron radius of heavy nuclei |
| ζ | Neutron star structure |

- Existent observables insufficient to constrain all parameters
- Determination of neutron radii of neutron-rich nuclei presses!
- Simultaneous mass-radius measurement of neutron stars presses!

Crucial measurements in Heaven and Earth on the horizon!



Equation of State Mean-Field Models Nuclear Observables Nuclear Collisions The Symmetry Energy

Neutron Stars

Page 5 of 12

Accurately Calibrated Parametrizations

The Program:

- Input binding energy and charge radii of doubly magic nuclei Solve in self-consistent mean-field approximation
- Compute the linear response of the mean-field ground state Solve in self-consistent MF+RPA approximation
- Without any further adjustment, compare to EoS from nuclear collisions Up to five times nuclear-matter saturation density
- Without any further adjustment, predict neutron-star structure Only physics that neutron stars are sensitive to — is the EoS of neutron-rich matter ...

| Α | Observable | Experiment | NL3 | NL3_030 | FSUGold_000 | FSUGold* | |
|-------------------|-----------------------------|----------------|-------------|------------|-------------|-----------|--|
| | | | Stiff-Stiff | Stiff-Soft | Soft-Stiff | Soft-Soft | |
| 90 Zr | $B/A \ ({ m MeV})$ | 8.71 | 8.69 | 8.70 | 8.68 | 8.68 | |
| | $R_{ m ch}~({ m fm})$ | 4.26 | 4.26 | 4.27 | 4.25 | 4.25 | |
| | $R_n - R_p \ (\mathrm{fm})$ | | 0.11 | 0.08 | 0.09 | 0.09 | |
| | GMR (MeV) | 17.89 ± 0.20 | 18.62 | 18.75 | 17.89 | 17.98 | |
| ²⁰⁸ Pb | $B/A \ ({\rm MeV})$ | 7.87 | 7.88 | 7.89 | 7.87 | 7.89 | |
| | $R_{ m ch}~({ m fm})$ | 5.50 | 5.51 | 5.52 | 5.51 | 5.52 | |
| | $R_n - R_p \ (\mathrm{fm})$ | | 0.28 | 0.20 | 0.29 | 0.21 | |
| | GMR (MeV) | 14.17 ± 0.28 | 14.32 | 14.74 | 13.73 | 14.04 | |
| | GDR (MeV) | 13.30 ± 0.10 | 12.70 | 13.07 | 12.79 | 13.07 | |

* Disclaimer: Gold is referred to the color — not the metal!



Equation of State Mean-Field Models Nuclear Observables Nuclear Collisions The Symmetry Energy Neutron Stars



Experimental Extraction of $R_n - R_p$

JLAB Experiment 00-003 (03-011) [Michaels, Souder, Urciuoli]:

- Parity Violating Asymmetry in elastic e-Pb scattering
- Electroweak (as opposed to hadronic) probe of neutron density
- Weak-vector boson Z^0 couples strongly to neutrons
- A clean and accurate measurement of the neutron radius 1% or 0.05 fm measurement of the neutron radius of 208 Pb



| Particle | EM coupling | Weak-Vector coupling | | | | | |
|---|-------------|--|--|--|--|--|--|
| up-quark | +2/3 | $+1 - 4\sin^2\theta_{\rm w}(+2/3) \simeq +1/3$ | | | | | |
| down-quark | -1/3 | $-1 - 4\sin^2\theta_{\rm w}(-1/3) \simeq -2/3$ | | | | | |
| proton | +1 | $+1 - 4\sin^2\theta_{\rm w} \simeq 0$ | | | | | |
| neutron | 0 | -1 | | | | | |
| $g_{\rm v}^f = 2T_z^f - 4\sin^2\theta_{\rm w}Q^f, \sin^2\theta_{\rm w} \approx 0.231 \simeq 1/4$ | | | | | | | |



Equation of State Mean-Field Models Nuclear Observables Nuclear Collisions The Symmetry Energy Neutron Stars



High Densities in Earth

[Danielewicz, Lacey, and Lynch – Science 298, 1592 (2002)]

Nuclear Collisions: Constraints and Predictions

- Sole earthly tool available to compress nuclear matter
- Compressions up to several (five) times nuclear saturation density
- Imprint of the EoS left in the flow and fragment distribution



FSUGold provides a reliable extrapolation to high density ...

●First ●Prev ●Next ●Last ●Go Back ●Full Screen ●Close ●Quit



Equation of State Mean-Field Models Nuclear Observables Nuclear Collisions The Symmetry Energy Neutron Stars



Neutron Skin and Neutron-Star Radii

Question: Is there a correlation between the neutron skin of ^{208}Pb and the radius of a "canonical" $1.4M_{\odot}$ neutron star? Answer: Probably yes! Same pressure that pushes neutrons out in ^{208}Pb pushes neutrons out in a neutron star.

- Isolated radio-quiet neutron stars already discovered
- Find good candidates for mass-radius measurement

Interesting Correlation:

The neutron skin of ²⁰⁸Pb depends on the EOS below N.M. saturation density, while the radius of the neutron star is also sensitive to the high-density EOS. "The thinner the skin of ²⁰⁸Pb, the smaller the radius of the star"



Large neutron skin together with a small neutron-star radius, could provide strong signature in favor of a phase transition ...



| Equation of State |
|---------------------|
| Moon Field Models |
| Wean-Field Wodels |
| Nuclear Observables |
| |
| Nuclear Collisions |
| |
| The Symmetry Energy |
| |
| Neutron Stars |



Maximum (Limiting) Neutron-Star Mass

- Maximum mass determined by high-density behavior of EOS (ζ)
- Radius of low-mass stars determined by symmetry energy (Λ_v)

At present both parameters are poorly constrained! However, situation could improve very rapidly ...

- Find a single "heavy" neutron star (PSRJ0751?)
- Measure the neutron radius of ²⁰⁸Pb (JLAB?)





Equation of State Mean-Field Models Nuclear Observables Nuclear Collisions The Symmetry Energy Neutron Stars



Neutron Star Composition

The composition of non-exotic stars is controlled by the density dependence of the symmetry energy!

- Symmetry energy imposes a penalty for violating (N = Z) balance
- The stiffer the symmetry energy, the higher the price
- The stiffer the symmetry energy, the higher the proton fraction





Equation of State Mean-Field Models Nuclear Observables Nuclear Collisions The Symmetry Energy Neutron Stars



Electron Fraction and Neutron-Star Cooling

Enhanced (direct URCA) cooling of non-exotic stars?

- Core-collapse Supernovae generates proto-neutron star $(T_{\rm core} \simeq 10^{12} {\rm K})$
- Direct URCA process cools down the star until $(T_{\rm core} \simeq 10^9 {\rm K})$
- Depending on the EoS direct URCA may continue Is Y_p large enough to conserve momentum?
- Best case for DUrca (soft-stiff): Soft EoS for symmetric matter \rightarrow large ρ_c Stiff symmetry energy \rightarrow large Y_p

Direct URCA process:

a)
$$n \to p + e^- + \bar{\nu}$$

b)
$$p + e^- \rightarrow n + \nu_i$$

may continue cooling the neutron star. FSUGold predicts that the pulsar in 3C58 does NOT need to be an exotic (quark) star.





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| Nuclear Collisions |
| The Symmetry Energy |
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High Densities in Heaven

Neutron Stars: Constraints and Predictions

- Sole heavenly tool available to compress nuclear matter
- Compression up to several (ten?) times nuclear saturation density
- Imprint of the EoS left in limiting mass, radius, cooling history, ...

| М | odel | $l \mid k_{\mathrm{F}}^0 \; (\mathrm{fm}^{-1}) \mid \epsilon$ | | ϵ_0 | $\epsilon_0 \; (\text{MeV})$ | | K (MeV) | J (MeV) | | L (MeV) | |
|-------------|-----------------|---|------|--------------|------------------------------|----|---------|---------|-----|---------|--|
| NL3 1.30 | | -16.2 | | | 271 37. | | .4 118 | | .5 | | |
| SUGold 1.30 | | -16.3 | | 230 | 32 | .6 | 60.5 | | | | |
| | Obse | rvable | NL | 3 | NL3_03 | 0 | FSUGold | 1_000 | FSU | JGold | |
| | ρ_c (f | $m^{-3})$ | 0.05 | 52 | 0.085 | | 0.05 | 1 | 0. | 076 | |
| | R (| km) | 15.0 |)5 | 14.18 | | 13.80 |) | 12 | 2.66 | |
| | $M_{\rm max}$ | $_{\rm c}(M_\odot)$ | 2.7 | 8 | 2.75 | | 1.80 | | 1 | .72 | |
| | $ ho_{ m Urca}$ | $({\rm fm}^{-3})$ | 0.2 | 1 | 0.51 | | 0.22 | | 0 | .47 | |
| | $M_{ m Urca}$ | $_{ m a}(M_{\odot})$ | 0.8 | 4 | 2.64 | | 0.74 | : | 1 | .30 | |
| | ΔM | $I_{\rm Urca}$ | 0.3 | 8 | 0.00 | | 0.59 | | 0 | .06 | |

Some Questions and Answers (FSUGold):

• Is the pulsar in 3C58 an exotic star?

Not necessarily if $M_* > 1.3 M_{\odot}$

• Is the limiting mass of a neutron star $M_{\text{max}} \simeq 1.72 \ M_{\odot}$? Report suggests $M(PSRJ0751 + 1807) = 2.1^{+0.4}_{-0.5}$

Fascinating times ahead!