

Quasars, Pulsars, Gamma-Ray Bursts! Oh, my!



Collapsars, magnetars, X-ray pulsars, γ -ray pulsars, millisecond radio pulsars...

Giacconi et al. (1962)



- Rocket carried 3 Geiger counters to 225km - one counter failed (#1).
- Windows of counters pointed 55° from the axis of the rocket (i.e. the Zenith).
- Windows made of mica covered with lampblack.
- Rocket rotated at 2.0 rps, 350 s above 80km.

The Data

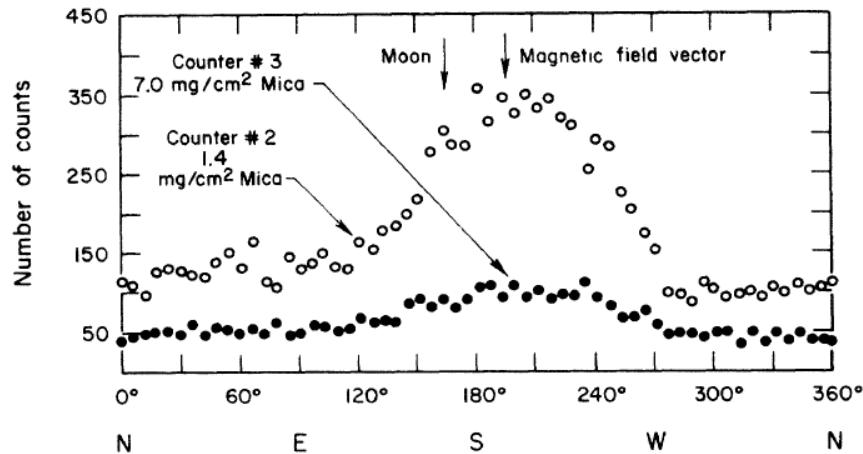
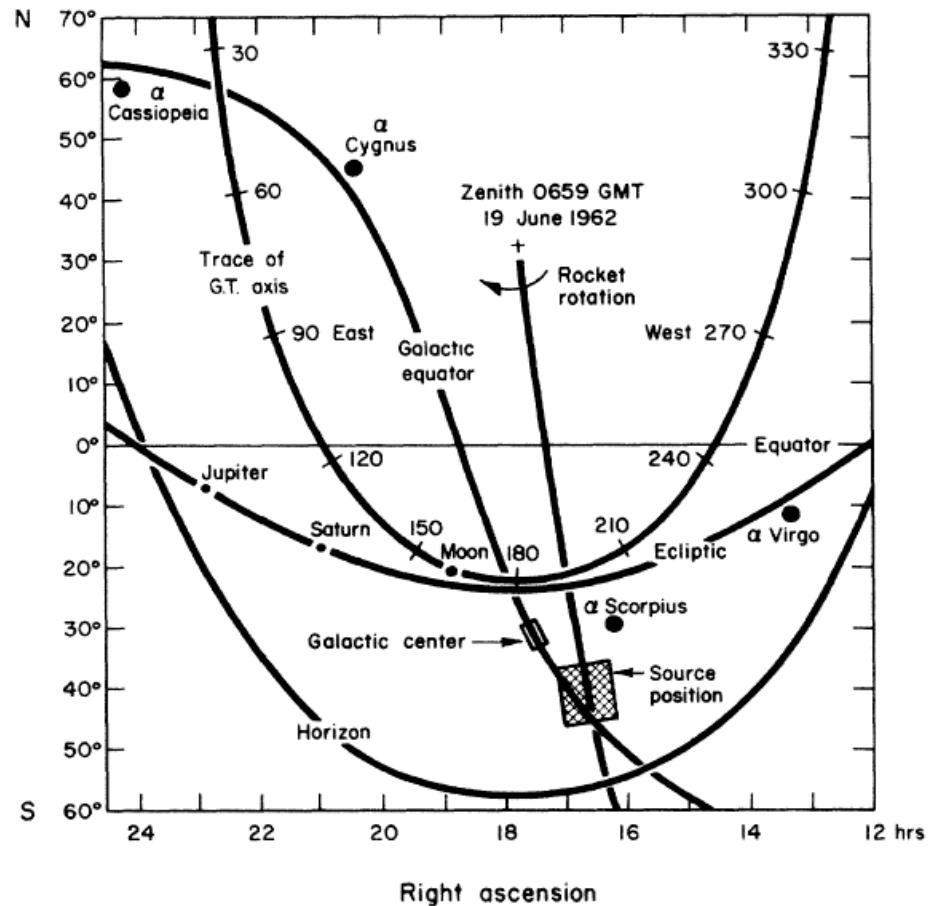


FIG. 1. Number of counts versus azimuth angle. The numbers represent counts accumulated in 350 seconds in each 6° angular interval.



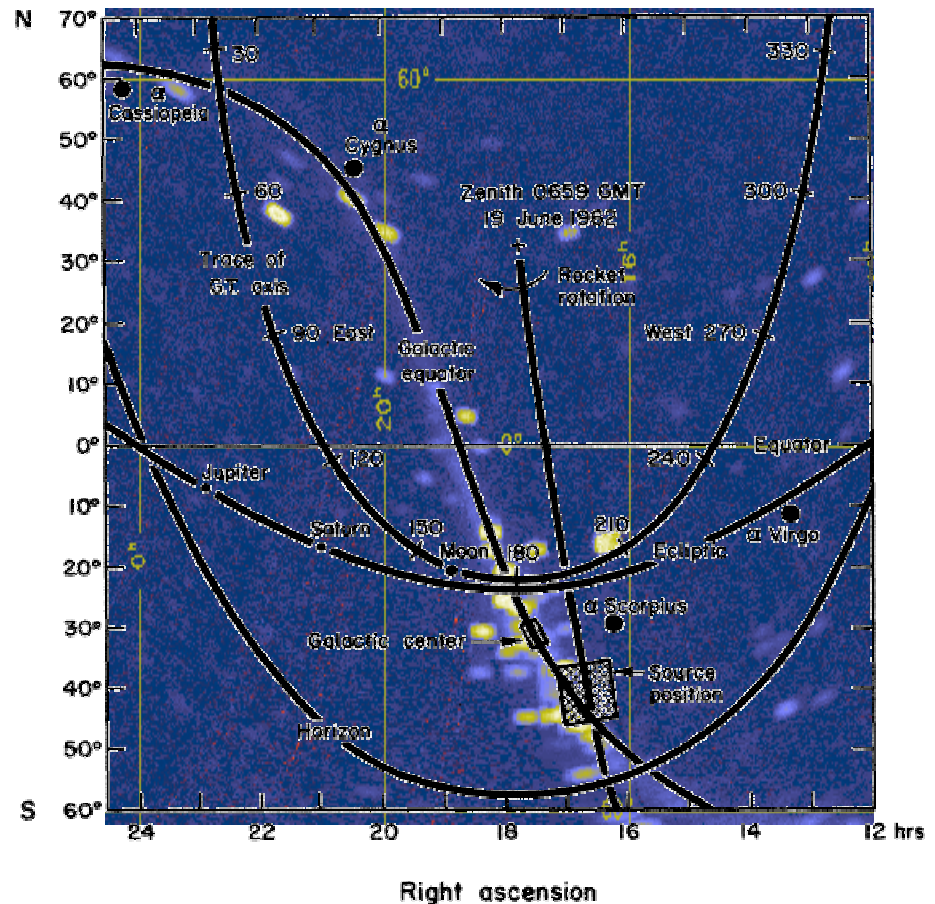
The Results



- They argued that the source was located about 10° above the horizon by comparing the absorption through the air and the mica assuming monochromatic source.

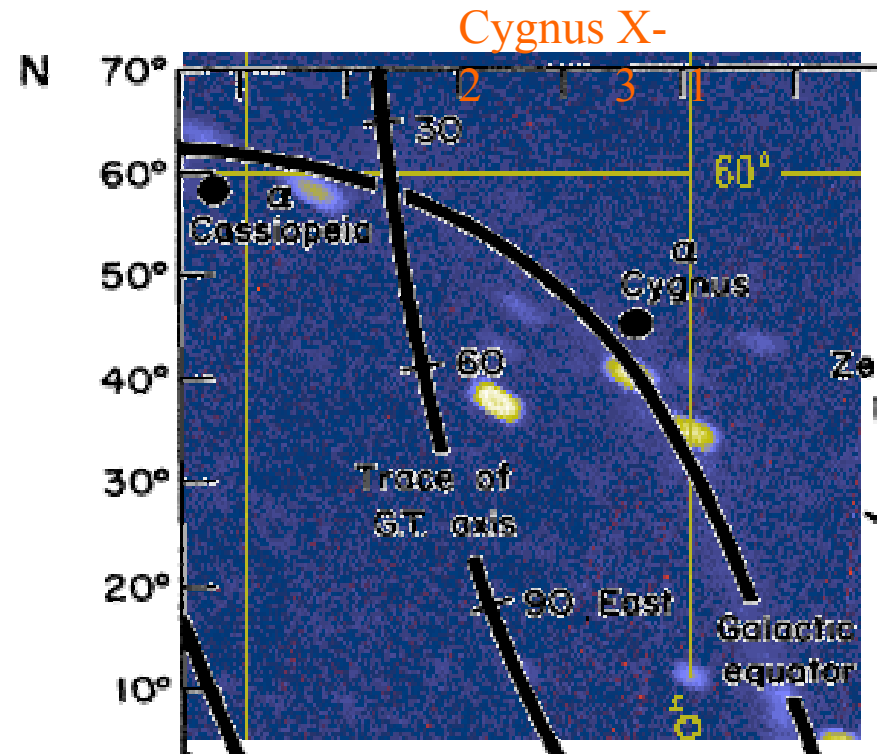
What Did the Aerobee See?

- Fig. 2 superimposed on the HEAO all-sky survey.
- The object known as Sco X-1 is the brightest in the X-ray sky at an azimuth angle of 210° along the G.T. axis.



What about the scruff at 60°?

- The detection at 60° albeit not as dramatic as the 210° result is also associated with an LMXB, Cygnus X-2.
- What is the flux of Cygnus X-2 compared with Sco X-1?



The Answer:



- Sco X-1: 14000 μJy (LMXB)
- Cyg X-1: 235-1320 μJy (HMXB)
- Cyg X-2: 450 μJy (LMXB)
- Cyg X-1: 90-430 μJy (HMXB)
- BTW:
 - What were they looking for from the moon?
 - What became of American Science and Engineering?

LMXBs and HMXBs



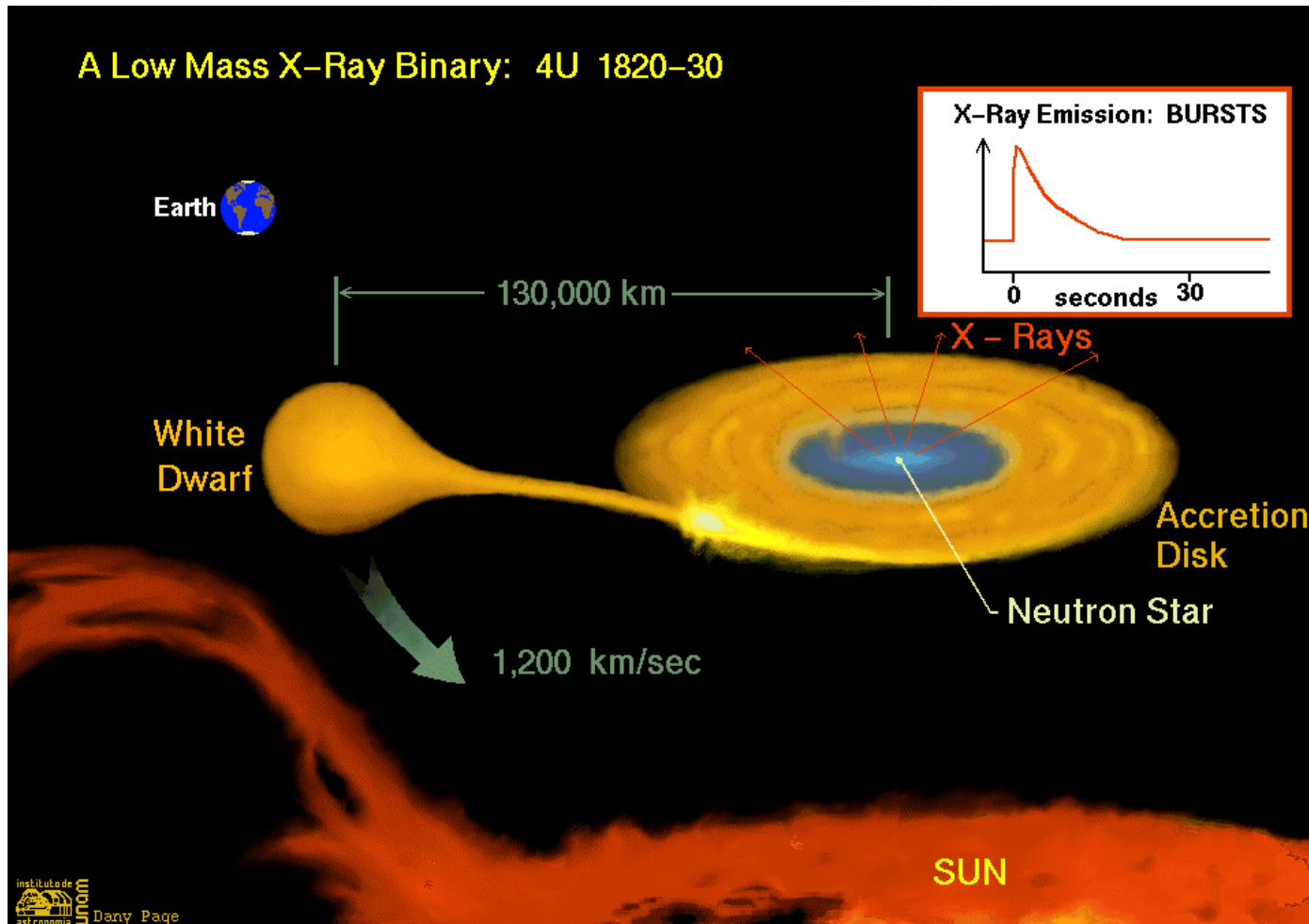
■ Low-mass X-ray Binaries

- Low mass main sequence star or white dwarf in orbit with a neutron star or black hole
- Roche lobe overflow driven by gravitational radiation

■ High-mass X-ray Binaries

- High mass main sequence star in orbit with a neutron star or black hole
- Wind accretor or Roche lobe overflow

LMXBs



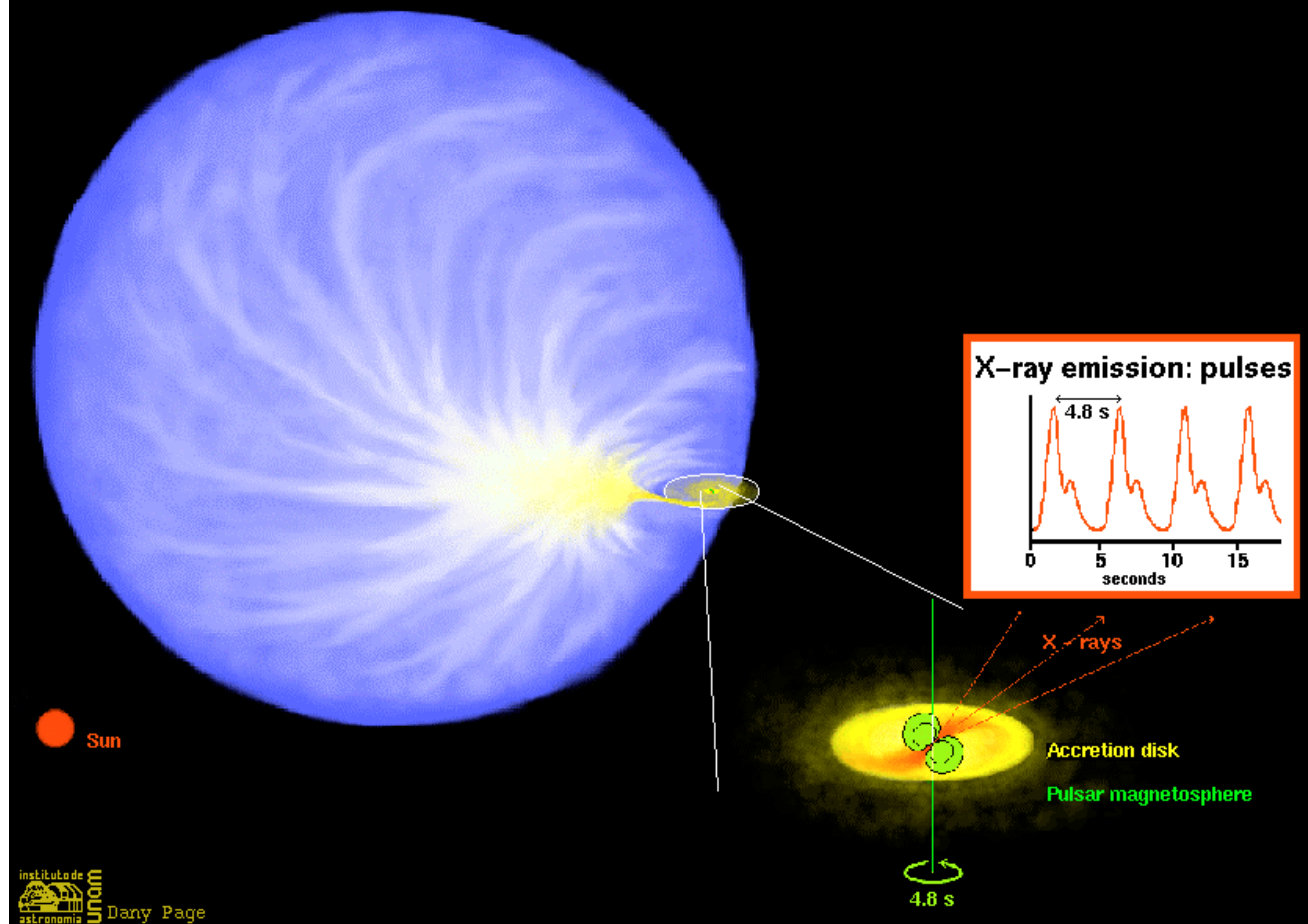
LMXB Properties



- Long-lasting system - GR drives the evolution
- NS/BH can accrete much of the donor's mass
 - The donor might be disrupted completely.
 - The neutron star gains lots of angular momentum.
 - Magnetic field of the neutron star gets "buried."
- The faint donor is tidally distorted.
 - Orbit is circularized and tidally locked.
 - Accretion rate does not depend on orbital phase.
 - It is hard to see the changing aspect of the donor in the glare of the accretion disk.
- Result: MSP with WD or alone.

HMXBs

CENTAURUS X-3: A HIGH MASS X-RAY BINARY



HMXB Properties



- Short-lived system - stellar evolution
- NS/BH accretes little of the donor's mass
 - The donor may evolve normally.
 - Magnetic field of the neutron star preserved.
- The bright donor may only be minimally distorted.
 - Orbit may be elliptical.
 - Accretion rate may depend on orbital phase.
- Result: BH/NS with BH/NS/WD or alone.

Binary Stellar Evolution

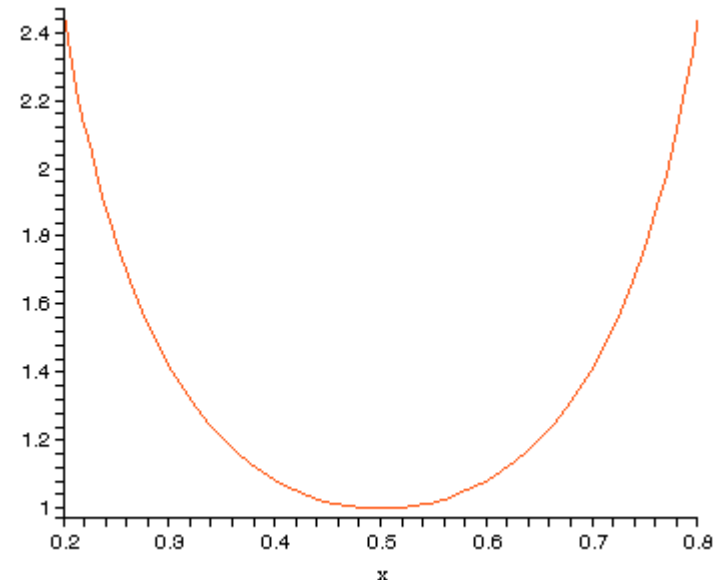
- When mass moves from one star to the other, the orbit changes. Need to remember Kepler's Laws!

- Second Law: $\Omega a^2 M_1 M_2 / (M_1 + M_2) = L$

- Third Law: $\Omega^2 a^3 = G (M_1 + M_2)$

- 2nd + 3rd:

$$a = \frac{L^2 M}{G M_1^2 M_2^2}$$



LMXBs and HMXBs



- In an LMXB, the donor is **less** massive than the neutron star, so the orbit will widen ending the mass transfer, unless L is not conserved or the donor expands as it loses mass.
- In an HMXB, the donor is **more** massive than the neutron star, so the orbit will shrink increasing the mass transfer.

Roche Lobe

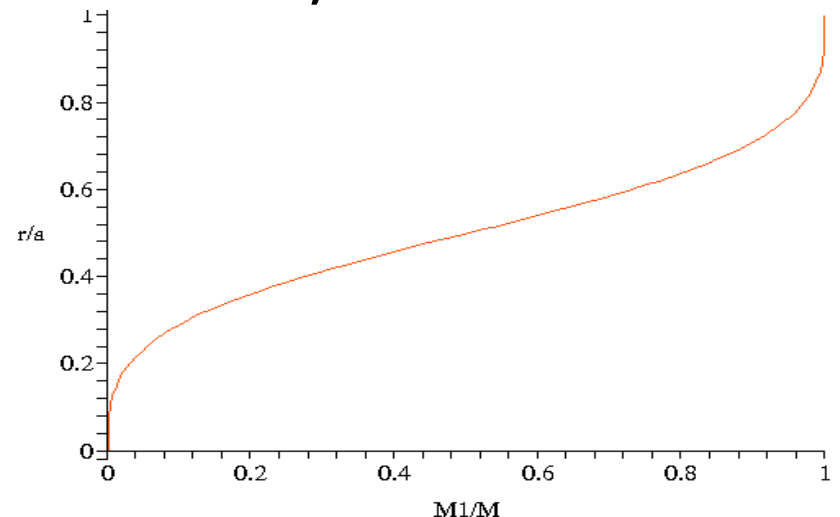
- The force pulling a particle toward the centre of the star equals the force pulling it away.

$$\frac{F}{m} = -\frac{GM_1}{r^2} + \frac{GM_2}{(a-r)^2} - \Omega^2 \left(a\frac{M_2}{M} - r \right)$$

- With some algebra you find that,

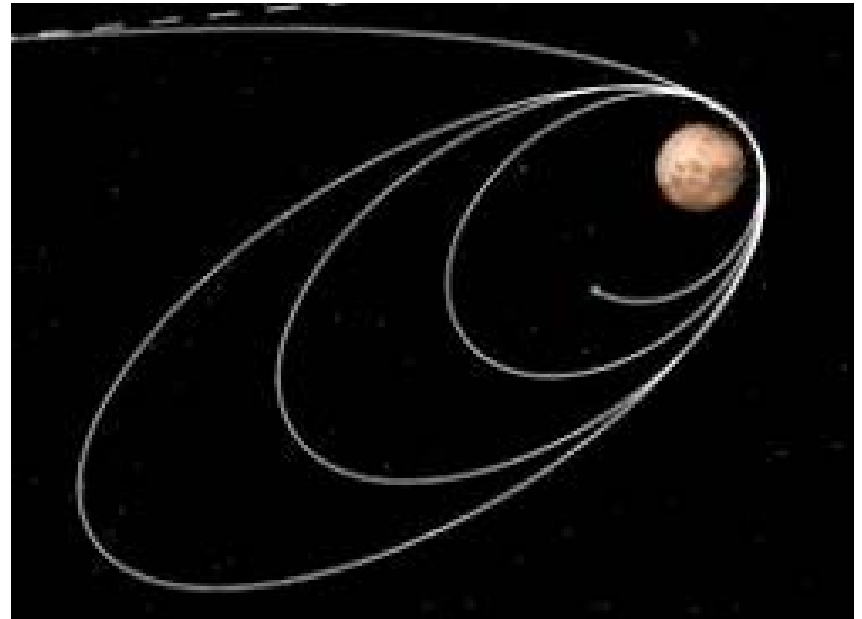
$$\frac{F}{m} = \frac{GM}{a^2} f \left(\frac{r}{a}, \frac{M_1}{M} \right)$$

where $f = \mathbf{UGLY}$.



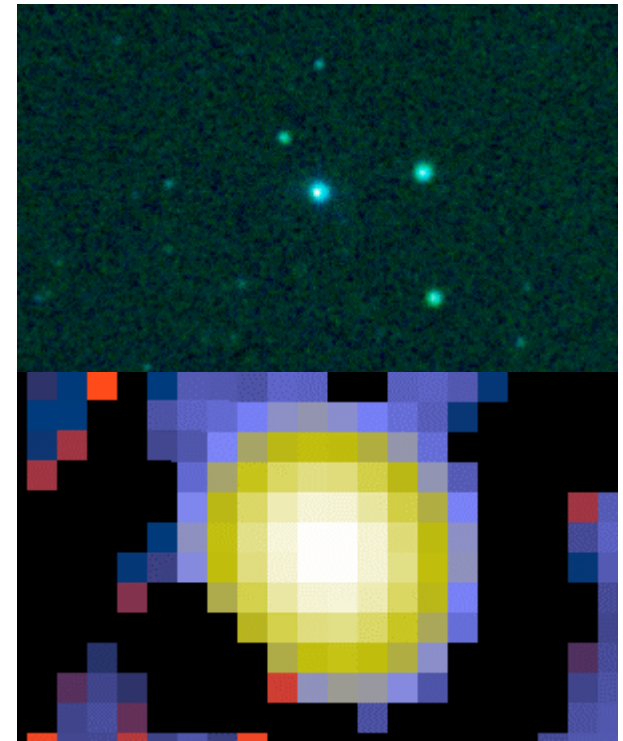
Tidal Circularization

- A topical analogy to understand how this works is aerobraking.
- Raising the tides dissipates orbital energy as friction in the star.
- Because the energy is removed at periastron, the apastron moves inward.



Quasars - Active Galaxies

- The nuclear activity of galaxies spans a wide spectrum ranging from normal galaxies through Seyfert galaxies and radio galaxies to quasars and QSOs.
- QSO - quasi-stellar object
- Quasar - quasi-stellar radio source
- The term QSO has fallen out of use.
- The key surprise about 3C 273 and 3C 48 was of course their distance.
- Luminosity of about 10^{47} erg/s from a small region -- you'll estimate their masses in the problem set.



3C 273 in B/R/J (DSS/2MASS)
and at 1.4 GHz (VLA NVSS)

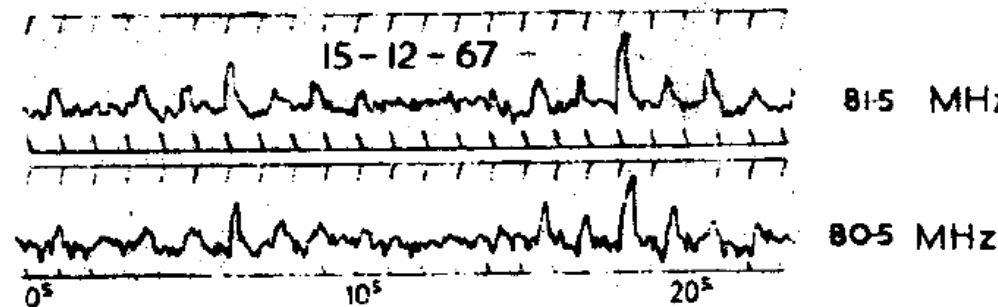
Hewish et al. (1968)

- Built a large multibeam antenna to look at scintillation of quasar radiation through the interplanetary plasma.
 - Sensitive
 - High time resolution
 - Low frequency



What did they observe?

- “Bits of scruff that appeared at the same sidereal time each day”
- The pulses arrive slightly later at lower frequencies.
- Used the expected frequency drift due to the Earth’s motion to determine the position of the source on the sky.
- The position did not change over the year.



- The pulse frequency is constant except for the Earth’s Doppler shift.
- Suggested that it was pulsation of a white dwarf or neutron star.

Light through Plasmas

- A plasma is a fluid in which the charged particles are free to move (there are currents).

$$\nabla \times \mathbf{B} = \frac{4\pi}{c}\mathbf{J} + \frac{1}{c}\frac{\partial \mathbf{E}}{\partial t} \quad \nabla \times \mathbf{E} + \frac{1}{c}\frac{\partial \mathbf{B}}{\partial t} = 0$$

- Taking the time derivative of (1) and putting into the curl of (2),

$$\nabla \times (\nabla \times \mathbf{E}) + \frac{4\pi\partial \mathbf{J}}{c^2 \partial t} + \frac{1}{c^2} \frac{\partial^2 \mathbf{E}}{\partial t^2} = 0$$

$$\nabla \times (\nabla \times \mathbf{E}) = \nabla(\nabla \cdot \mathbf{E}) - \nabla^2 \mathbf{E}$$

- Let's assume that no net charge develops so the first term vanishes.

Now the currents

- A current is a charge times a velocity,

$$\mathbf{J} = n_e e \mathbf{v} \quad \text{and} \quad \dot{\mathbf{J}} = n_e e \dot{\mathbf{v}}, \quad m_e \dot{\mathbf{v}} = e \mathbf{E}$$

where we have assumed that $v \ll c$.

- Combining the last few equations gives,

$$-\nabla^2 \mathbf{E} + \frac{4\pi n_e e^2}{c^2 m_e} \mathbf{E} + \frac{1}{c^2} \frac{\partial^2 \mathbf{E}}{\partial t^2} = 0$$

- Let $\mathbf{E} = \mathbf{E}_0 e^{i(kx - \omega t)}$ and substitute

$$k^2 \mathbf{E} + \frac{\omega_p^2}{c^2} \mathbf{E} - \omega^2 \mathbf{E} = 0 \quad \text{with} \quad \omega_p^2 = \frac{4\pi n_e e^2}{m_e}$$

Dispersion Measure (1)

- EM radiation travelling through a plasma has the dispersion relation:

$$\omega^2 = \omega_p^2 + c^2 k^2 \quad \text{so} \quad v_g = \frac{d\omega}{dk} = c \frac{2k}{\omega} = c \sqrt{1 - \frac{\omega_p^2}{\omega^2}}$$

- So, the pulse travels faster at higher frequencies,

$$t = \frac{L}{v_g} \quad \text{so} \quad \frac{dt}{d\nu} = \frac{L}{c} \frac{\nu_p^2}{\nu^3 \left(1 - \frac{\omega_p^2}{\omega^2}\right)^{3/2}}$$

Dispersion Measure (2)

- We can also use an integrated quantity,

$$t_2 - t_1 = \frac{2\pi e^2}{mc} (\omega_2^{-2} - \omega_1^{-2}) \int_0^d n_e dl$$

- The dispersion constant is

$$D = (t_2 - t_1) / (\nu_2^{-2} - \nu_1^{-2})$$

and the dispersion measure is

$$DM \text{ (cm}^{-3}\text{pc)} = 2.410 \times 10^{-16} D \text{ (Hz)}$$

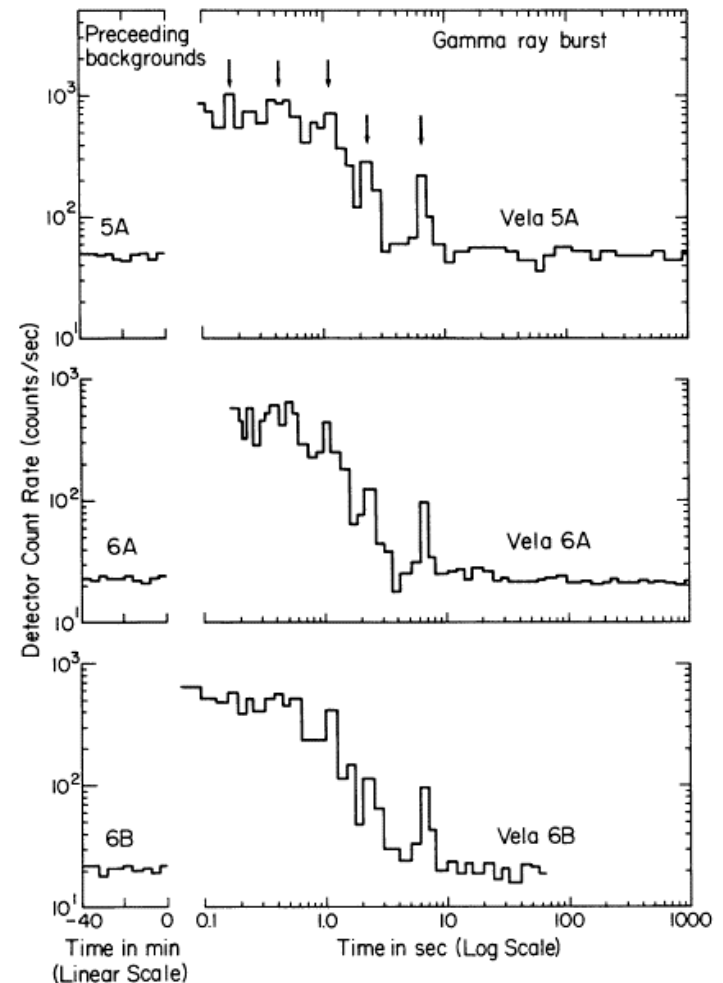
Why is it important?



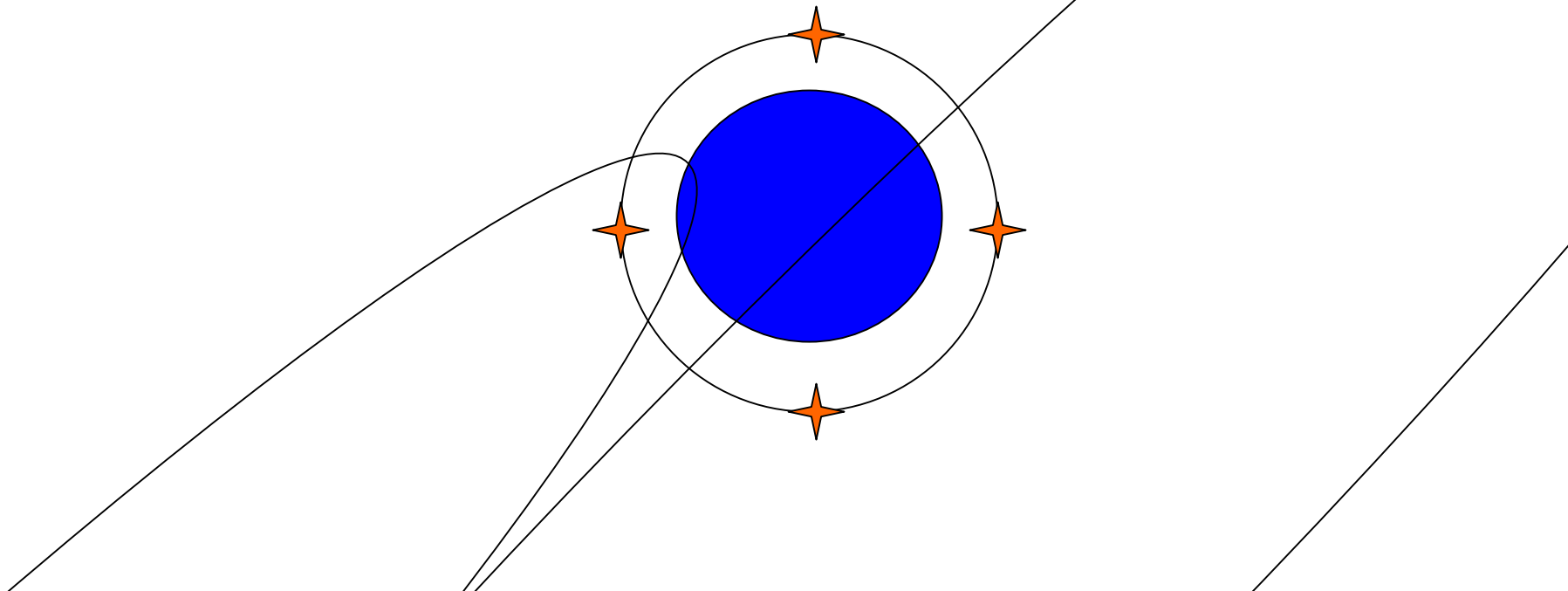
- The paper has outlined many of the techniques used in pulsar astronomy today:
 - DM distances
 - Pulse counting
 - Doppler positioning
 - The name “pulsar” for pulsating star (this ended up being the wrong model)
- How did they do it?
- The key was that they were looking for scintillation, so they
 - Looked at low frequencies where pulsars are brightest.
 - Used a fast time constant, multiple frequencies and lag cables to see the time shifts induced by scintillation.
 - Sensitivity - pulsars on average aren't that bright.

Klebesadel et al.

- Four Vela satellites in Earth orbit observed several bursts of gamma-rays.
- Could use timing of the burst arrival to determine the location on the sky and extra-terrestrial origin.



Locating GRBs



Fast Forward

