

LECTURE 15

SUPERCONDUCTIVITY

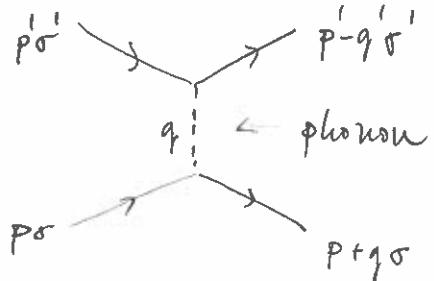
- Basic phenomenon: zero resistivity, Meissner effect
Cooper instability, BCS wavefunction

- Consider BCS "pairing Hamiltonian" $H = H_0 + H_1$,

$$H_0 = \sum_{k\sigma} \epsilon_k c_{k\sigma}^+ c_{k\sigma} \quad \epsilon_k = \frac{k^2}{2m} - \mu$$

$$H_1 = -\frac{1}{2} V \sum_{\substack{pp'q \\ \sigma\sigma'}} c_{p+q\sigma}^+ c_{p'-q\sigma'}^+ c_{p'+\sigma'} c_{p\sigma} \quad \leftarrow \begin{array}{l} \text{attractive interaction} \\ (\text{due to phonons}) \end{array}$$

- We will study a two-electron GF at finite T : (pair)



$$F(\vec{p}, \vec{p}', \vec{q}; \tau) = \langle T [c_{p+q\sigma}(\tau) c_{-p\sigma}(\tau) c_{p+q\sigma}(0) c_{-p\sigma}^+(0)] \rangle$$

$$\tau \in (-\beta, \beta) \quad \beta = \frac{1}{k_B T}$$

We calculate it using FO expansion in terms of unperturbed GF

$$G^0(\vec{p}, \tau) = \langle T [c_p(\tau) c_p^+(0)] \rangle$$

$$G^o(\vec{p}, \tau) = \sum_n e^{i\gamma_n \tau} G^o(\vec{p}, \nu_n)$$

$$G^o(\vec{p}, \nu_n) = \frac{1/\beta}{i\nu_n + \epsilon_p} \quad \nu_n = \frac{\pi}{\beta} (2n+1) \quad n=0, \pm 1, \dots$$

F is a two-electron GF, therefore its Matsubara frequencies are bosonic

$$F(\vec{p}, \vec{p}', \vec{q}; \tau) = \sum_m e^{i\omega_m \tau} F(\vec{p}, \vec{p}', \vec{q}; \omega_m)$$

$$\omega_m = (2\pi/\beta)m \quad m = 0, \pm 1, \dots$$

• Zero-order:

$$F^o(\vec{p}, \vec{p}', \vec{q}; \omega_m) =$$

$$= -\delta_{pp'} \sum_n G^o(p+q, \omega_m + \nu_n) G^o(-p, -\nu_n)$$

• In higher order we consider ladder diagrams that represent processes where electrons in a pair repeatedly scatter from one another:

$$F(\vec{p}, \vec{p}', \vec{q}; \omega_m) =$$

Because we assumed momentum-indep. interaction V in the BCS Hamiltonian the sums over $p_1, p_2 \dots$ separate and one can exactly sum all the diagrams as

$$F(\vec{q}, \omega_m) = \frac{F^0(\vec{q}, \omega_m)}{1 + \beta V F^0(\vec{q}, \omega_m)}$$

$$F^0(\vec{q}, \omega_m) = \sum_{\vec{p}} F(\vec{p}_1 \vec{p}_1, \vec{q}; \omega_m)$$

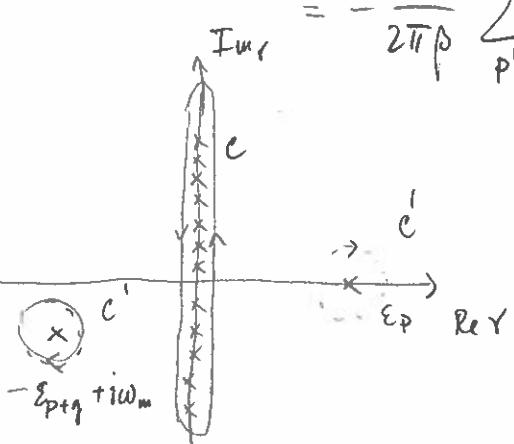
- SC instability is signaled by the divergence in pair propagator $F(\vec{q}, \omega_m)$,

$$1 + \beta V F^0(\vec{q}, \omega_m) = 0$$

$$F^0(\vec{q}, \omega_m) = -\frac{1}{\beta^2} \sum_{p, n} \frac{1}{(i\omega_m + i\nu_n + \epsilon_{p+q})(-i\nu_n + \epsilon_p)}$$

restriction on
 \vec{p} summation
to $|\epsilon_p| < k_B T_D$
 T_D - Debye temp.

$$I_{m, c} = -\frac{i}{2\pi\beta} \sum_{p'} \oint_C \frac{dy}{e^{py} + 1} \frac{1}{(\nu + \epsilon_{p+q} + i\omega_m)(-\nu + \epsilon_p)}$$



- The largest (negative) value of $F^0(\vec{q}, \omega_m)$ occurs when $\omega_m = 0 \Rightarrow$ we are interested in $F^0(\vec{q}, 0)$.

$$\begin{aligned}
 F(\vec{q}, 0) &= \frac{1}{\beta} \sum_p' \left[\frac{1}{(e^{\beta \epsilon_p} + 1)(\epsilon_p - \epsilon_{p+q})} - \frac{1}{(e^{-\beta \epsilon_p} + 1)(\epsilon_p + \epsilon_{p+q})} \right] \\
 &= -\frac{1}{\beta} \sum_p' \frac{1 - f_{p+q} - f_p}{\epsilon_{p+q} + \epsilon_p} \\
 &= \odot \frac{1}{2\beta} \sum_p' \frac{\tanh(\frac{1}{2}\beta \epsilon_{p+q}) + \tanh(\frac{1}{2}\beta \epsilon_p)}{\epsilon_{p+q} + \epsilon_p}
 \end{aligned}$$

The largest value occurs when $\vec{q} = 0$, therefore the denominator is

$$1 + \beta V F^*(0, 0) = 1 - V \sum_p' \frac{\tanh(\frac{1}{2}\beta \epsilon_p)}{2\epsilon_p}$$

$$\begin{aligned}
 \int_{-\infty}^{k_B T_0} d\epsilon N(\epsilon) \frac{\tanh(\frac{1}{2}\beta \epsilon)}{2\epsilon} &\approx N(0) \int_0^{k_B T_0} d\epsilon \frac{\tanh(\frac{1}{2}\beta \epsilon)}{\epsilon} \\
 &= N(0) \left[\ln\left(\frac{1}{2}\beta k_B T_0\right) - \underbrace{\int_0^{\frac{1}{2}\beta k_B T_0} dx \frac{\ln x}{\cosh^2 x}}_{\ln(\frac{\pi}{4\gamma})} \right] \\
 &\quad \xrightarrow{\frac{1}{2}\beta k_B T_0 \rightarrow \infty} \ln\left(\frac{\pi}{4\gamma}\right) \simeq 0.8187
 \end{aligned}$$

γ - Euler's const.

$$1 = V N(0) \ln \left(\frac{2\gamma \beta k_B T_0}{\pi} \right)$$

$$\Rightarrow T_c = 1.4 T_0 e^{-1/V N(0)}$$

BCS formula for superconducting critical temperature