

Phys525:
Quantum Condensed Matter Physics: Quantum Criticality
Basics, Dynamics and Topological criticality

Episode 14:

Application: $z=2$ free particle/Strong interacting fixed points of fermions

(A quick wrap up!)

$$\mathcal{L} = \mathcal{L}_f + \mathcal{L}_b + \mathcal{L}_{bf}$$

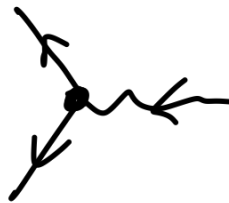
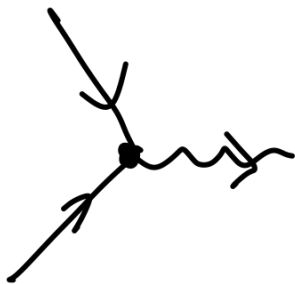
χ : Fermionic (e)
 ϕ : Bosonic (2e)

$$\mathcal{L}_f = \chi_\alpha^* \partial_\tau \chi_\alpha - \chi_\alpha^* \frac{\nabla^2}{2} \chi_\alpha - M_f \chi_\alpha^* \chi_\alpha \quad \alpha = \uparrow, \downarrow$$

$$\mathcal{L}_b = \phi^* \partial_\tau \phi - \phi^* \frac{\nabla^2}{4} \phi + (\epsilon_b - 2M_f) \phi^* \phi$$

$$\mathcal{L}_{bf} = g \chi_\alpha^* \sigma_{y,\alpha\beta} \chi_\beta^* \phi + c.c. \quad -M_b$$

↑
 Anti-Symmetric



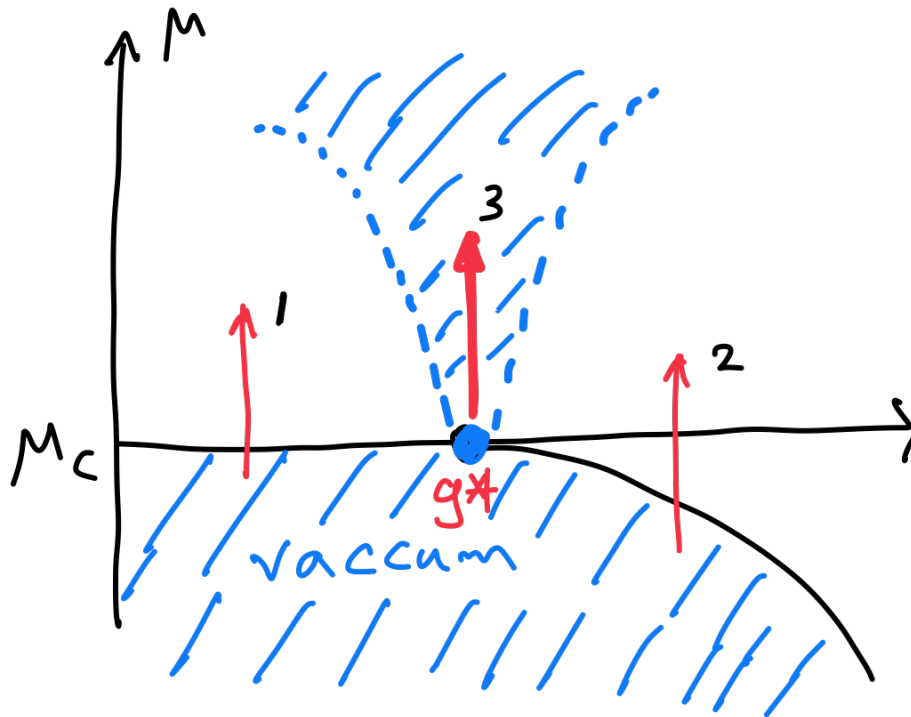
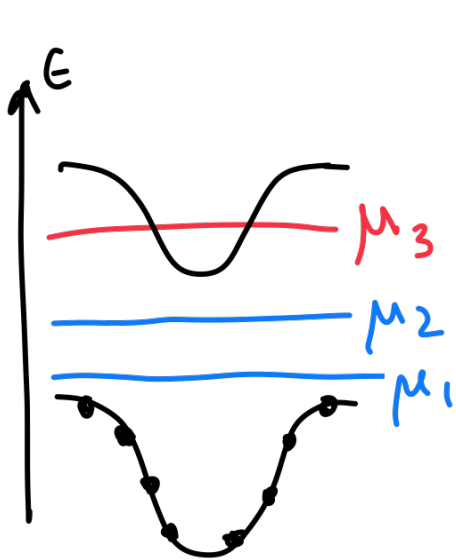
Grassmannian field

$$\{\chi_{\alpha R}, \chi_{\alpha' R}\} = \{\chi_{\alpha R}, \chi_{\alpha L}\} = 0$$

$$\chi_\alpha = \chi_{\alpha R} + i \chi_{\alpha I}, \quad \chi_{\alpha R}^2 = \chi_{\alpha I}^2 = \frac{1}{2}, \quad \{\chi_{\alpha L}, \chi_{\alpha' L}\} = \{\chi_{\alpha L}, \chi_{\alpha' R}\} = 0$$

Related
 $z=2$

General fermionic Model (3d)



attractive
interaction
↓
 g

1: free-fermion Universality class

2: free-boson Universality class

3: $SO(2,1)$ Conformal field Theory / CFT

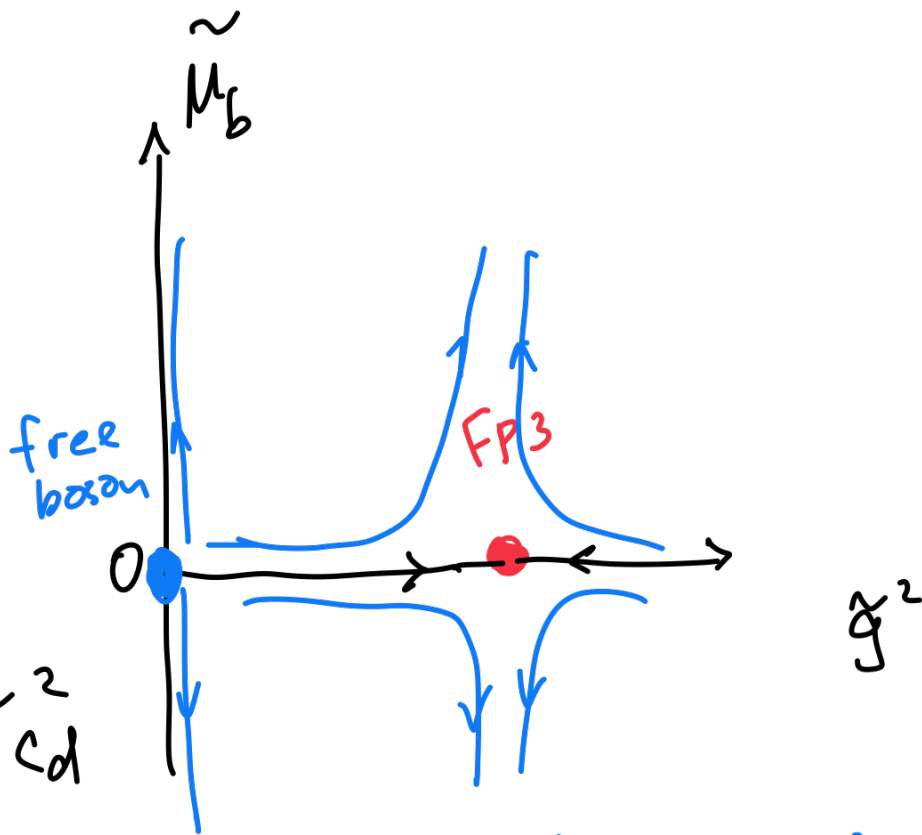
RGE

$$\frac{d\tilde{M}_f}{dt} = -2\tilde{M}_f$$

$$\frac{dg^2}{dt} = (d-4)g^2 + g^4$$

$$\frac{d\tilde{M}_b}{dt} = -2\tilde{M}_b + 4\tilde{M}_f + g^2 c_d$$

$$\frac{1}{\tilde{Z}_b} \frac{d\tilde{Z}_b}{dt} = g^2$$



fixed pt 2: $\tilde{M}_b^* = 0, g^{*2} = 4-d$

$\tilde{M}_f^* \rightarrow \infty$

$$\left(\beta_{\tilde{M}_b} = 0 \rightarrow 2\tilde{M}_f = \tilde{M}_b(n) - \frac{g^{*2} c_d}{2} \wedge \right)$$

fixed pt 2:

$$d < 4$$

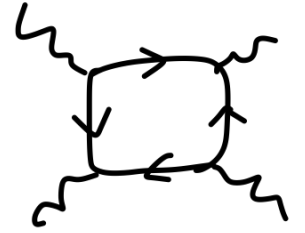
$$\tilde{M}_b = 0, \tilde{g}^{*2} = 4 \cdot d, \tilde{M}_f \rightarrow \infty$$

= Bosons interacting with massive fermions

$$\tilde{M}_b = 0,$$

$$d_{\text{eff}} = \varphi^* (\partial_t - \frac{\nabla^2}{2}) \varphi + \lambda |\varphi^* \varphi|^2 + \dots$$

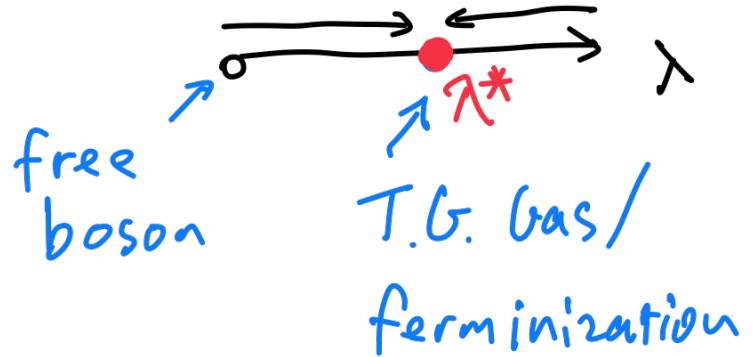
$$\tilde{\lambda} \propto \tilde{g}^{*4}$$



$d = 2, 3$



$d = 1$



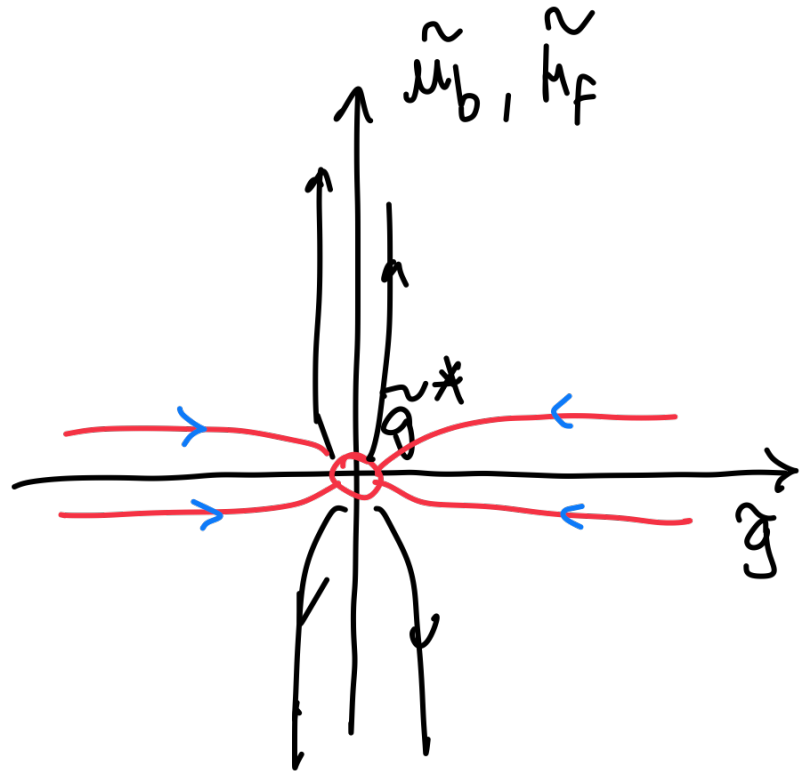
RGE

$$\frac{d\tilde{M}_f}{dt} = -2\tilde{M}_f$$

$$\frac{d\tilde{g}^2}{dt} = (d-4)\tilde{g}^2 + \tilde{g}^4$$

$$\frac{d\tilde{\mu}_b}{dt} = -2\tilde{\mu}_b + 4\tilde{M}_f^2 + \tilde{g}^2 c_d$$

$$\frac{1}{Z_b} \frac{dZ_b}{dt} = +\tilde{g}^2$$

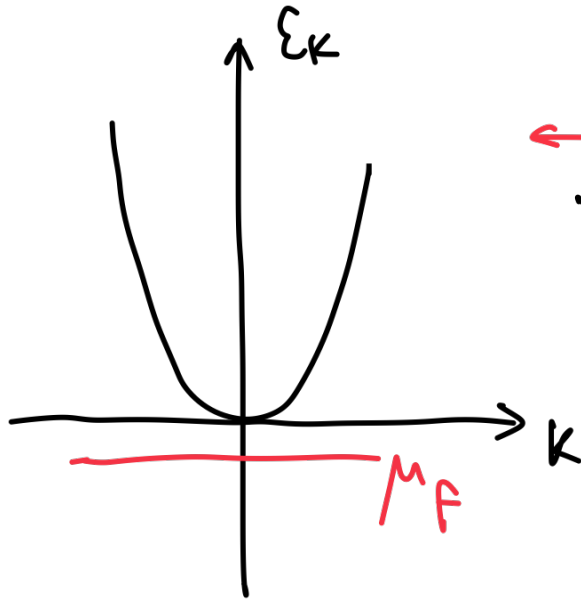


IR stable

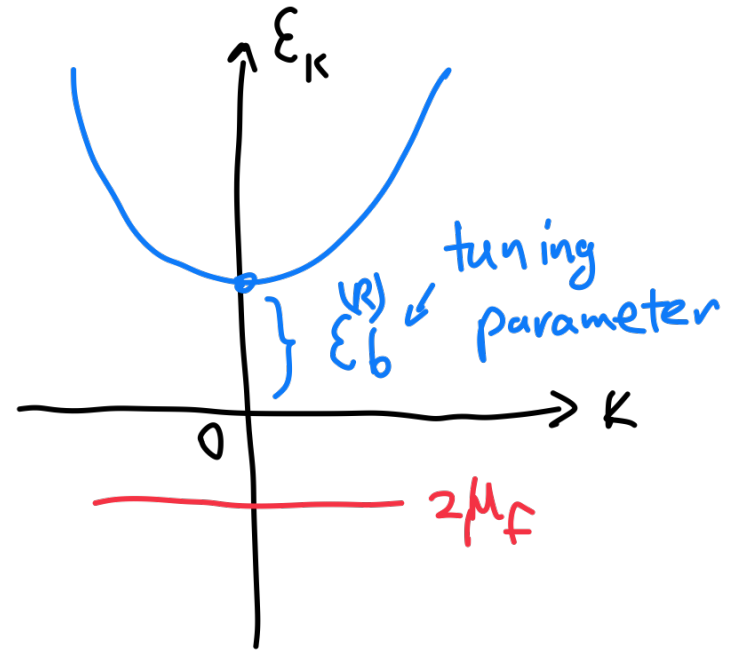
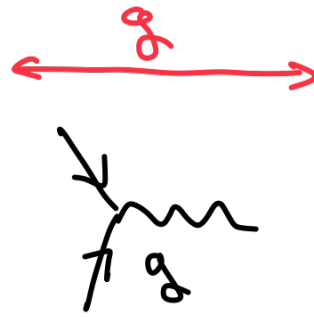
fixed pt 3: $\tilde{M}_f = \tilde{\mu}_b = Z_b = 0$
 $\tilde{g}^2 = 4-d = \tilde{g}^{2*}$

"No-boson" QFT

Connections to Feshbach Resonance



Atoms



Feshbach Molecules

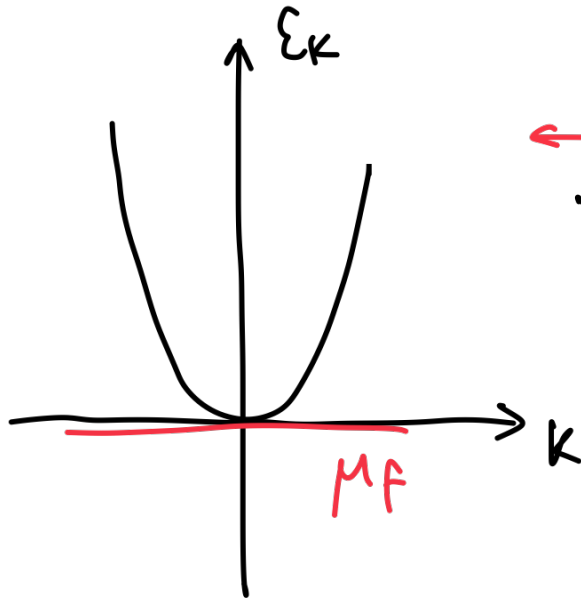
$$M_b = 2M_F - E_b^{(R)}$$

$$E_b^{(R)} = E_b - \frac{g^2}{2} \Lambda_w C_d$$



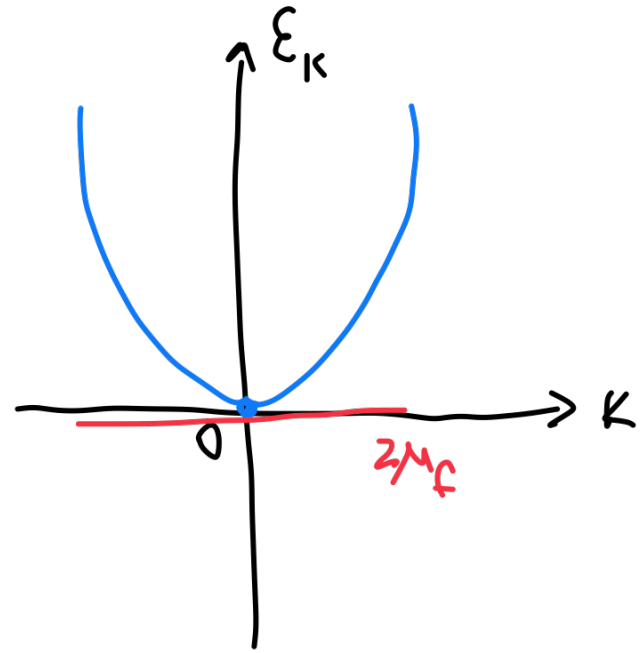
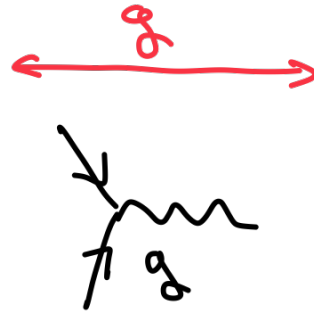
At Feshbach Resonance

$$\epsilon_b^{(R)} = 0$$



Atoms

$$\mu_F = 0$$



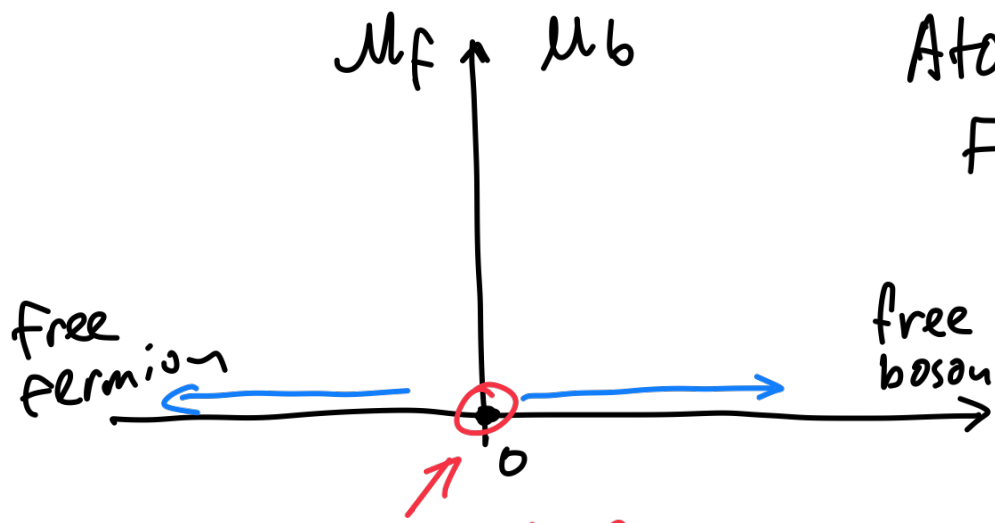
Feshbach Molecules

$$\mu_b = 0$$

$$\mu_b = 2\mu_F - \epsilon_b^{(R)} = 2\mu_F = 0$$

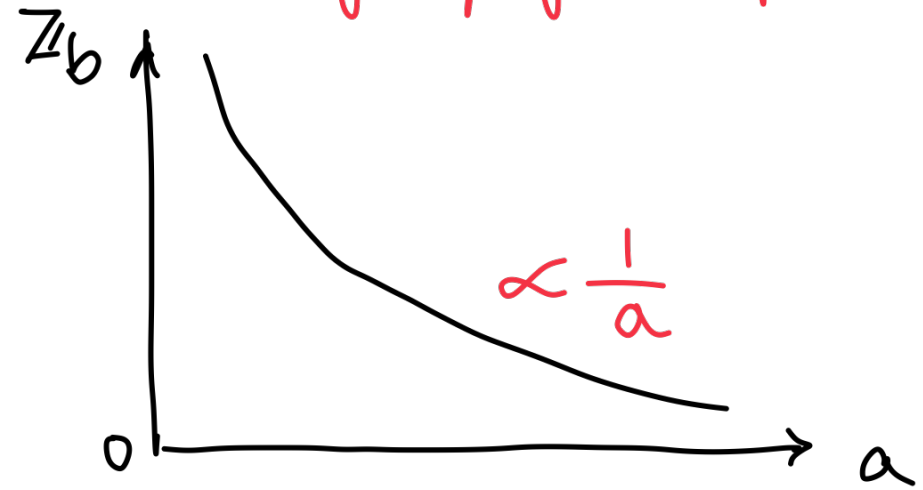
$$\epsilon_b^{(R)} = \epsilon_b - \frac{g^2}{2} \Lambda_{uv} C_d = 0$$

Atomic Gas Near Feshbach Resonance ($d=3$)



$1/a \leftarrow$ Scattering length

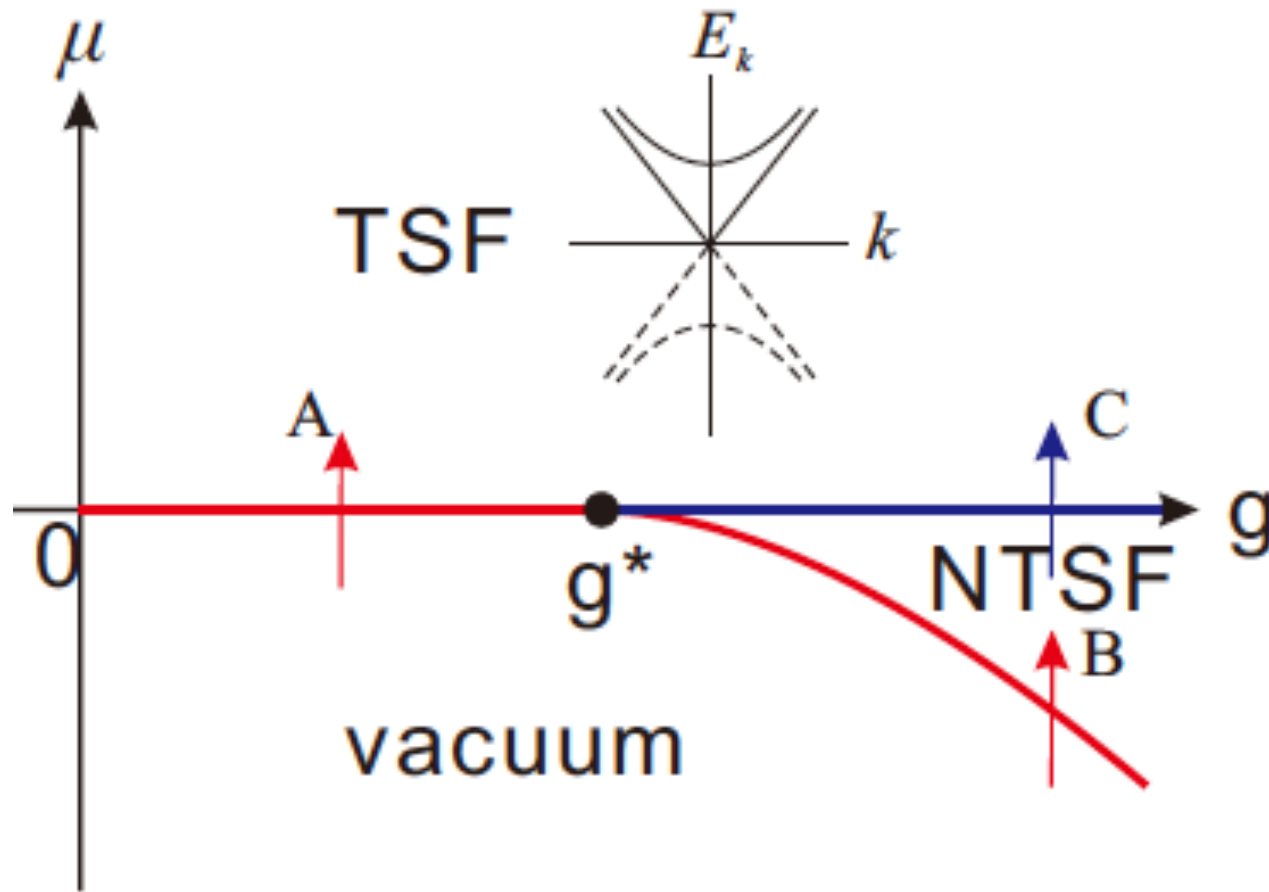
Feshbach Resonance (Strong Coupling Fixed pt)



Feshbach Resonance
 = $SO(2,1)$ CFT theory
 with $Z_b = 0$ ($\tilde{g}^2 = 4-d$)

$$\frac{1}{Z_b} \frac{dZ_b}{dt} = \tilde{g}^2 = \tilde{g}^{*2} = 1 \rightarrow Z_b(L) \sim \frac{1}{L}$$

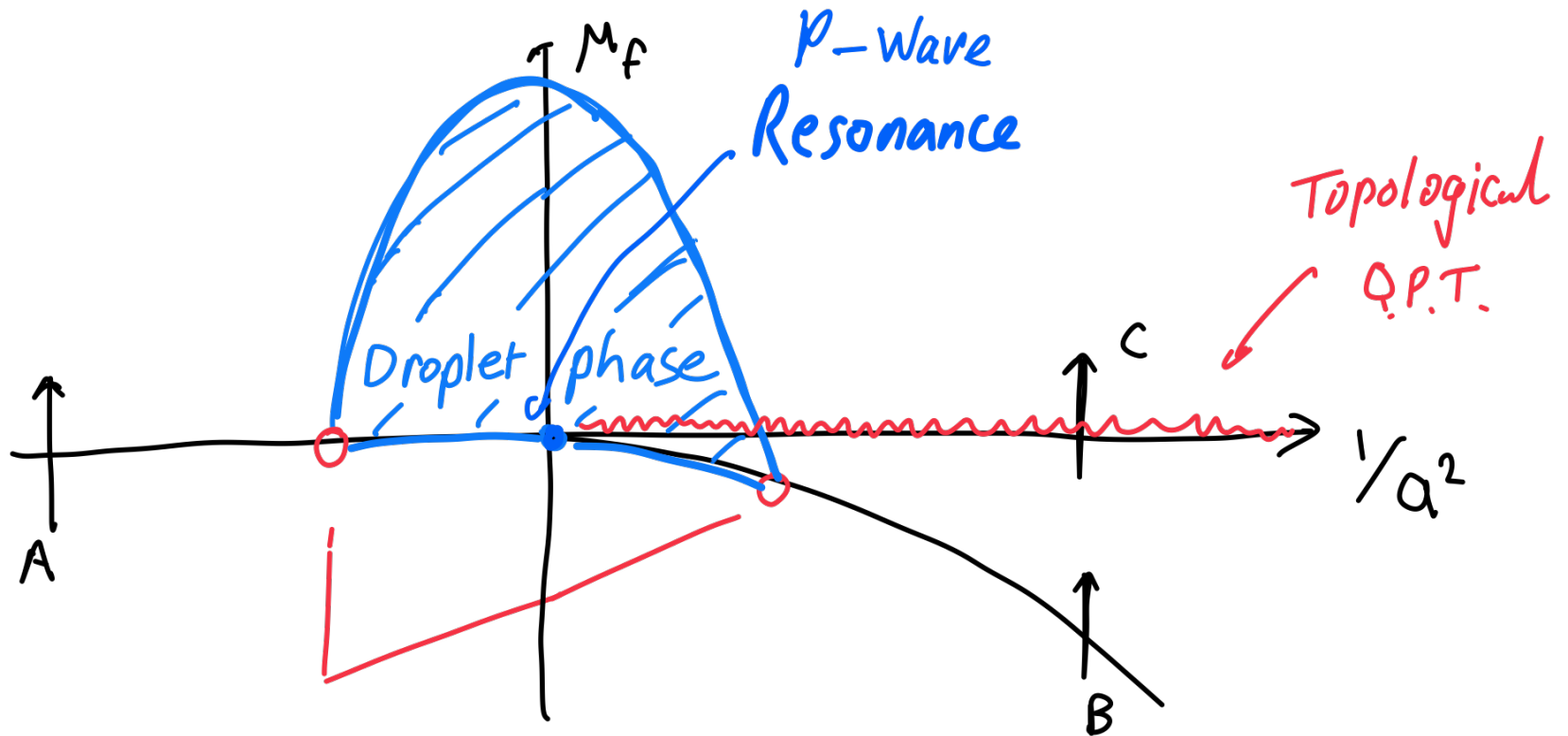
Example: Phase diagram of p+ip spinless SF ??



A: Free Fermion; B: Wilson-Fisher/Free boson; C: Majorana class
 g^* : QCP of $SO(2,1)$ CFT.

QFT/RGE Applications to p-wave interactions:

Tricritical Physics in Two-Dimensional p-Wave Superfluids, Fan Yang, Shao-Jian Jiang, and Fei Zhou, Phys. Rev. Lett. 124, 225701 (2020).



Tri-Critical