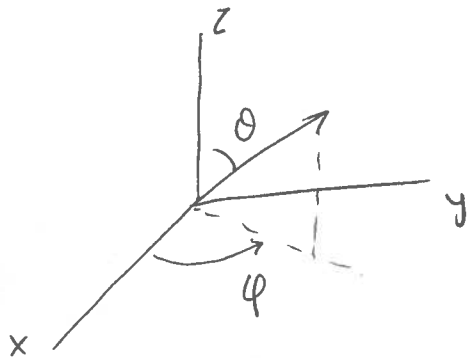


Home work III A

1) Bloch Sphere State $|\hat{\Omega}\rangle$ is defined as

$$\vec{S} \cdot \hat{\Omega} |\hat{\Omega}\rangle = \frac{\hbar}{2} |\hat{\Omega}\rangle$$

Following the previous NW set, $|\hat{\Omega}\rangle = \begin{pmatrix} \cos \frac{\theta}{2} \\ \sin \frac{\theta}{2} e^{i\varphi} \end{pmatrix}$



$$\langle S_x \rangle = \frac{\hbar}{2} \sin \theta \cos \varphi$$

$$\langle S_y \rangle = \frac{\hbar}{2} \sin \theta \sin \varphi$$

$$\langle S_z \rangle = \frac{\hbar}{2} \cos \theta$$

$\langle \vec{S} \rangle = \frac{\hbar}{2} \hat{\Omega}$ orientation along the direction of $\hat{\Omega}$

2) $|t=0\rangle = \begin{pmatrix} \cos \frac{\theta_0}{2} \\ \sin \frac{\theta_0}{2} e^{i\varphi_0} \end{pmatrix}$

$$\hat{n}_0 = \hat{n}_0(\theta_0, \varphi_0)$$

polar angle for \hat{n}_0 .

$$|t\rangle = \begin{pmatrix} \cos \frac{\theta_0}{2} e^{-i \frac{B\hbar t}{2}} \\ \sin \frac{\theta_0}{2} e^{i\varphi_0} \cdot e^{i \frac{B\hbar t}{2}} \end{pmatrix}$$

(Assume $H = \gamma \frac{\vec{B} \cdot \vec{S}}{2}$)

$$= \gamma \frac{B}{2} S_z$$

$$= e^{-i \frac{B\hbar t}{2}} \begin{pmatrix} \cos \frac{\theta_0}{2} \\ \sin \frac{\theta_0}{2} e^{i\varphi_0 + B\hbar t} \end{pmatrix}$$

$$\gamma = g \frac{e}{m} > 0$$

(differs from textbook def.)

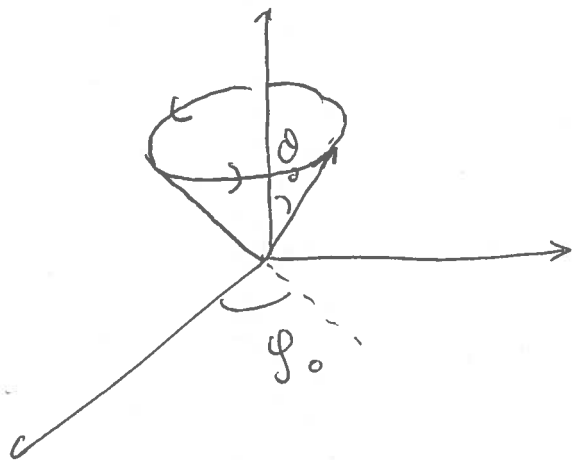
$$\hat{n}(t) = \hat{n}(\theta(t), \varphi(t)),$$

$$\theta(t) = \theta_0, \quad \varphi(t) = \varphi_0 + \omega t$$

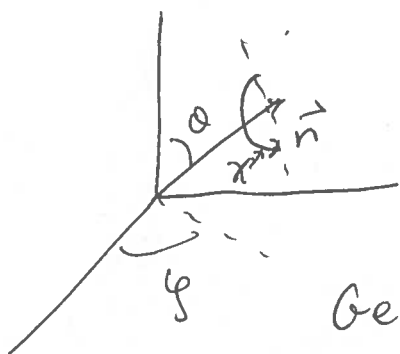
"Right hand Rule"

when it is plus.

"Larmor Precession"



$$3) U(\hat{n}, \chi) = \cos \frac{\chi}{2} - i \sin \frac{\chi}{2} \vec{\sigma} \cdot \hat{n} = e^{-i \frac{\vec{\sigma} \cdot \hat{n}}{2} \chi}$$

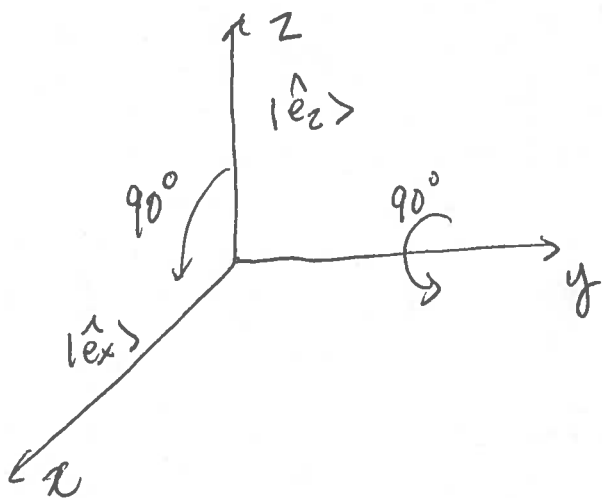


Generally,

$$|\hat{n}'\rangle = U(\hat{n}, \chi) |\hat{n}\rangle$$

After Rotation

Before Rotation



$$|\hat{\Omega}\rangle = |\hat{e}_z\rangle = \begin{pmatrix} 1 \\ 0 \end{pmatrix}$$

$$\hat{n} = \hat{e}_y, \quad \chi = \frac{\pi}{2}$$

$$U(\hat{e}_y, \frac{\pi}{2}) = \frac{1}{\sqrt{2}} (1 - i \sigma_y)$$

$$= \frac{1}{\sqrt{2}} \begin{pmatrix} 1 & -1 \\ 1 & 1 \end{pmatrix}$$

$$|\hat{\Omega}'\rangle = U(\hat{e}_y, \frac{\pi}{2}) |\hat{e}_z\rangle$$

$$= \frac{1}{\sqrt{2}} \begin{pmatrix} 1 \\ 1 \end{pmatrix} = |\hat{e}_x\rangle$$

desired state

4) in General, one can generate $SO(3)$ Rotations matrix using $SU(2)$ Rotations (First pointed out by Hamilton when studying classical motion of a top)

$$R_{\alpha\beta}(\hat{n}, \chi) = \frac{1}{2} \text{tr}(U^\dagger \sigma_\alpha U \sigma_\beta) \quad (*)$$

$SO(3)$ $\alpha, \beta = x, y, z$ $U = U(\hat{n}, \chi)$ — $SU(2)$ Rotation introduced before

$R_{\alpha\beta}(\hat{n}, \chi)$ is a 3×3 matrix

For $SO(3)$ Rotation, $r'_\alpha = R_{\alpha\beta} r_\beta$, $\alpha, \beta = x, y, z$
or

$$\vec{r}' = R \vec{r} \quad (**)$$

\downarrow After Rotation \downarrow vector before Rotation R — $SO(3)$ Rotation Matrix can be defined in (*)

if you can show R in (***) can be generated by R defined in (*), please come to my office

—— I'll take you out for a coffee!