

Phys525:  
Quantum Condensed Matter Physics: Quantum Criticality  
Basics, Dynamics and Topological states

Episode 12: Correlation function in “Z” Vs **Time Ordering** in quantum Ground states:

Slightly more formal stuff on Space-time temporal correlations

# imaginary time evolution (QFT: Vac-Vac Amplitude)

- Extract ground state properties (today)
- Extract space-time symmetries suitable for EFTs (today)
- Evaluate the dynamic correlations

# From classical correlation functions to causal Green's functions (retarded and advanced)

$$C(r, \tau; 0, 0) = \langle \psi(0, 0) \psi^*(r, \tau) \rangle_{Z_{qft}(D=d+1)}$$

$$\rightarrow G(\vec{r}, \tau; 0, 0) = \langle g.s. | -i\mathcal{T} b(\vec{r}, \tau) b^\dagger(0, 0) | g.s. \rangle$$

$$\rightarrow G^R(\vec{r}, t; 0, 0) = \langle g.s. | -i[b(\vec{r}, t), b^\dagger(0, 0)] | g.s. \rangle \theta(t)$$

$$\rightarrow G^A(\vec{r}, t; 0, 0) = \langle g.s. | -i[b(\vec{r}, t), b^\dagger(0, 0)] | g.s. \rangle \theta(-t)$$

$$C(\tau) = iG(\tau = it);$$

$$C(\Omega) = -G(i\Omega);$$

$$G(i\Omega \rightarrow \Omega + i\delta) = G^R(\Omega)$$

$$\rightarrow G^R(\Omega) \rightarrow G^R(\vec{r}, t)$$

$$Z = \langle 0 | e^{-H\tau} | 0 \rangle$$

$$= \int \mathcal{D}\phi e^{-\int d\vec{r} dz \mathcal{L}(\{\phi(\vec{r}, \tau)\})}$$

$S(\{\phi(\vec{r}, \tau)\})$

Eg.  $\tau \rightarrow \infty$

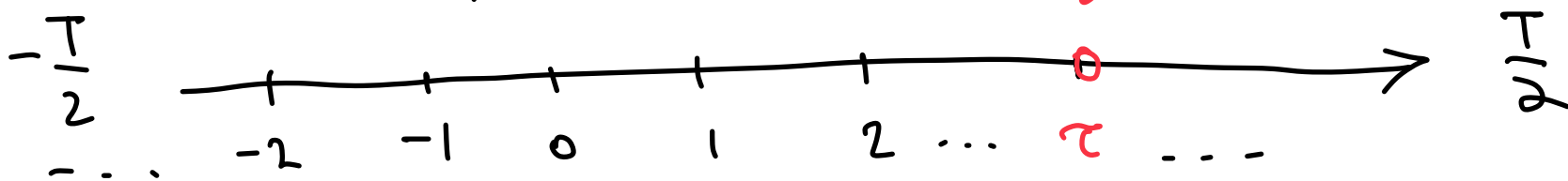
$$= -\frac{1}{\tau} \ln Z$$

$$\langle \phi \rangle_Z = \frac{\int \mathcal{D}\phi \phi(\vec{r}, \tau) e^{-S(\{\phi(\vec{r}, \tau)\})}}{\int \mathcal{D}\phi e^{-S(\{\phi(\vec{r}, \tau)\})}}$$

$$\langle \phi^*(\vec{r}, \tau) \phi(0,0) \rangle_Z = \frac{\int \mathcal{D}\phi \phi^*(\vec{r}, \tau) \phi(0,0) e^{-S}}{\int \mathcal{D}\phi e^{-S}}$$

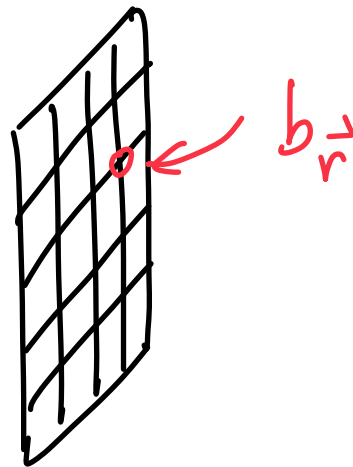
$$\langle \phi \rangle_Z = \frac{\int \Delta\phi \phi(\vec{r}, \tau) e^{-S(\{\phi(\vec{r}, \tau)\})}}{\int \Delta\phi e^{-S(\{\phi(\vec{r}, \tau)\})}}$$

What  $\langle \phi \rangle_Z$  stands for?



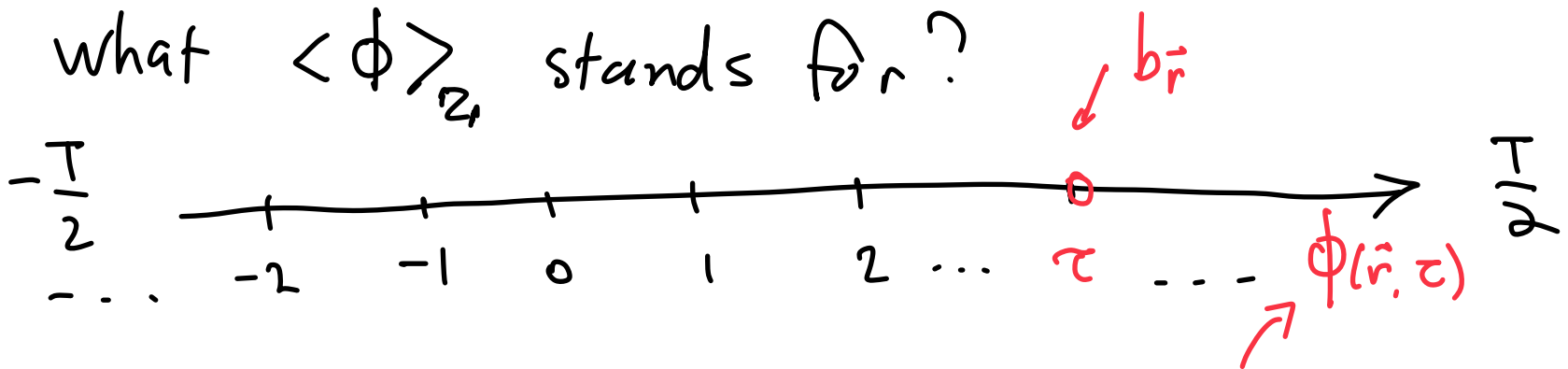
$$b_{\vec{r}} |\phi_{\vec{r}}\rangle = \phi(\vec{r}) |\phi_{\vec{r}}\rangle$$

↑  
"Coarse grained state"



$$\langle \phi \rangle_Z = \frac{\int \Delta\phi \phi(\vec{r}, \tau) e^{-S(\{\phi(\vec{r}, \tau)\})}}{\int \Delta\phi e^{-S(\{\phi(\vec{r}, \tau)\})}}$$

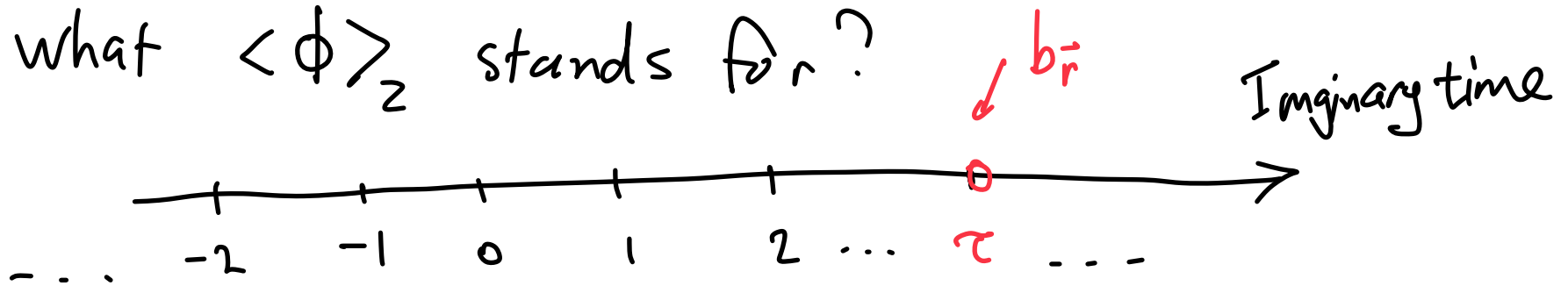
What  $\langle \phi \rangle_Z$  stands for?



$$Z = \sum_{\{\phi(\vec{r}, \tau+1)\}} \sum_{\{\phi(\vec{r}, \tau)\}} \dots \langle \{\phi(\tau+1)\} | e^{-H\Delta\tau} b_{\vec{r}} | \{\phi(\tau)\} \rangle \dots$$

$$= \langle 0 | e^{-H(\frac{\tau}{2} - \tau)} \underbrace{\hspace{10em}}_{\text{After}} b_{\vec{r}} e^{-H(\tau + \frac{\tau}{2})} | 0 \rangle \underbrace{\hspace{10em}}_{\text{before}}$$

What  $\langle \phi \rangle_z$  stands for?



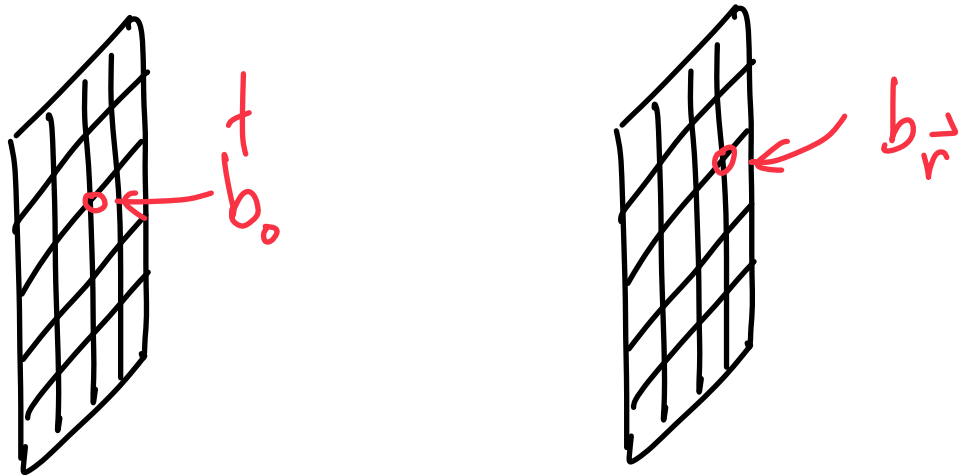
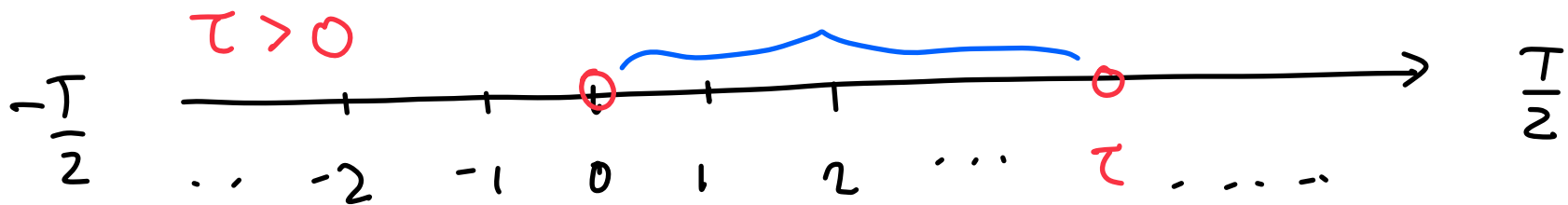
$$\langle \phi \rangle_z = \frac{Z(b_{\vec{r}}, \tau)}{Z} = \frac{\langle 0 | e^{-\left(\frac{T}{2} - \tau\right)H} b_{\vec{r}} e^{-\left(\tau + \frac{T}{2}\right)H} | 0 \rangle}{\langle 0 | e^{-TH} | 0 \rangle}$$

$$T \rightarrow \infty \Rightarrow \frac{|\langle 0 | g.s. \rangle|^2 e^{-E_{g.s.} T}}{|\langle 0 | g.s. \rangle|^2 e^{-E_{g.s.} T}} \cdot \langle g.s. | b_{\vec{r}} | g.s. \rangle$$

g.s. Quantum Average!

Average over  $\sum$  vs Quantum State average





$$\langle \phi(\vec{r}, \tau) \phi(0, 0) \rangle_{\mathcal{Z}}^* = \langle \text{g.s.} | \underbrace{e^{H\tau} b_{\vec{r}} e^{-H\tau}}_{b(\vec{r}, \tau)} b_0^{\dagger} | \text{g.s.} \rangle$$

(Heisenberg Rep.)

$\tau > 0$

$$\langle \phi(\vec{r}, \tau) \phi^*(0, 0) \rangle_{\mathcal{Z}} \stackrel{T \rightarrow \infty}{=} \langle \text{g.s.} | \underbrace{e^{H\tau} b_{\vec{r}} e^{-H\tau}}_{b^{\dagger}(\vec{r}, \tau)} b_0^{\dagger} | \text{g.s.} \rangle$$

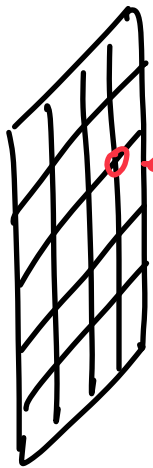
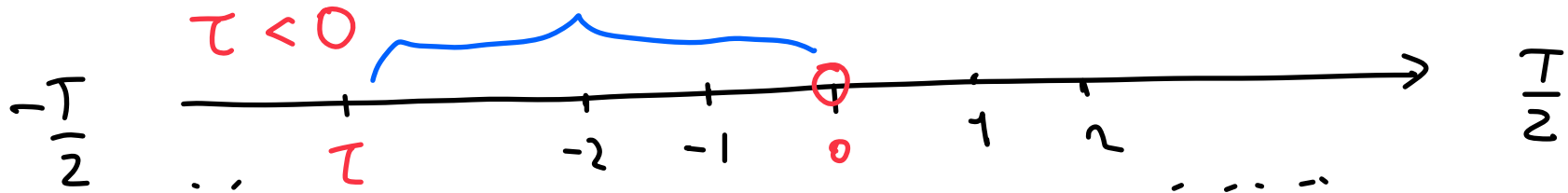
(Heisenberg Rep.)

$$\langle \text{g.s.} | b(\vec{r}, \tau) b_0^{\dagger} | \text{g.s.} \rangle$$

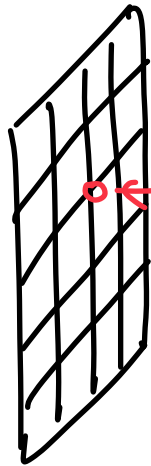
All excitations

$$= \sum_n \langle \text{g.s.} | b(\vec{r}, \tau) | n \rangle \langle n | b_0^{\dagger} | \text{g.s.} \rangle$$

Dynamic Correlation function  $\leftrightarrow$  excitations



$b_{\vec{r}}$



$b_0$

$$\langle \phi(\vec{r}, \tau) \phi(0, 0) \rangle_{\mathcal{Z}}^* = \langle \text{g.s.} | \overbrace{b_0 e^{-H(0-\tau)}}^{\tau} \underbrace{b_{\vec{r}} e^{-H\tau}}_{b(\vec{r}, \tau)} | \text{g.s.} \rangle$$

$\tau < 0$  (under  $\phi(\vec{r}, \tau)$ )  
 $\tau > 0$  (under  $b(\vec{r}, \tau)$ )

(Heisenberg Rep.)

# TIME ORDERING

$$\langle \phi(\vec{r}, \tau) \phi(0,0) \rangle_{\mathcal{Z}}^* \stackrel{\tau > 0}{=} \langle \text{g.s.} | \underbrace{e^{H\tau} b_{\vec{r}} e^{-H\tau}}_{b(\vec{r}, \tau)} b_0^\dagger | \text{g.s.} \rangle \stackrel{T \rightarrow \infty}{}$$

$$\langle \phi(\vec{r}, \tau) \phi(0,0) \rangle_{\mathcal{Z}}^* \stackrel{\tau < 0}{=} \langle \text{g.s.} | b_0^\dagger \underbrace{e^{H\tau} b_{\vec{r}} e^{-H\tau}}_{b(\vec{r}, \tau)} | \text{g.s.} \rangle \stackrel{T \rightarrow \infty}{}$$

Pull together,

$$\langle \phi(\vec{r}, \tau) \phi(0,0) \rangle_{\mathcal{Z}}^* = \langle \text{g.s.} | \mathcal{T} b_{\vec{r}}(\tau) b_0^\dagger | \text{g.s.} \rangle$$

$$\mathcal{T} b_{\vec{r}}(\tau) b_0^\dagger(0) = \Theta(\tau) b_{\vec{r}}(\tau) b_0^\dagger + \Theta(-\tau) b_0^\dagger b_{\vec{r}}(\tau)$$

Fluctuations near  $\langle \phi \rangle$ ,  $\phi \rightarrow \langle \phi \rangle + \phi$

$$\mathcal{L}(\{\phi(\vec{r}, \tau)\}) = |\partial_\tau \phi|^2 + |\nabla \phi|^2 + m^2 |\phi|^2, \quad m^2 = 4|\alpha|$$

$$Z = \int \prod_{\vec{r}, \tau} \Delta \phi e^{-\int d\vec{x} d\tau \mathcal{L}(\{\phi(\vec{r}, \tau)\})}$$

$$\approx \int \prod_{\omega, \vec{q}} \Delta \phi e^{-\sum_{\omega} \sum_{\vec{q}} \Phi^\dagger(\omega, \vec{q}) G^{-1}(\omega, \vec{q}) \Phi(\omega, \vec{q})}$$

$$= \prod_{\omega, \vec{q}} Z_{\omega, \vec{q}}$$

$$G^{-1}(\omega, \vec{q}) = \omega^2 + \vec{q}^2 + m^2$$

MF  $\rightarrow$

$$\langle \Phi^\dagger(\omega, \vec{q}) \Phi(\omega, \vec{q}) \rangle \approx \frac{1}{G(\omega, \vec{q})}$$

$\uparrow$   
 $Z_{\omega, \vec{q}}$

$\downarrow$   
HW Set II

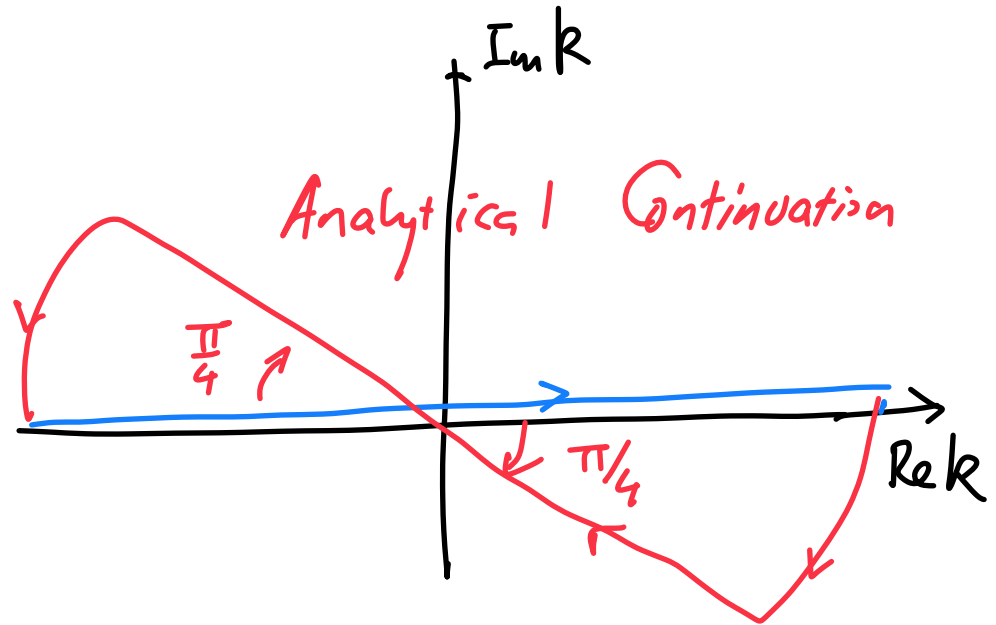
Time ordered in Non-Relativistic Theory (free particles)  
 $|0\rangle = |vac\rangle$

$$\langle 0 | T b_{\vec{r}}^\dagger(t) b_0(0) | 0 \rangle = \sum_{\vec{k}} \langle 0 | b_{\vec{r}}^\dagger(t) | \vec{k} \rangle \langle \vec{k} | b_0(0) | 0 \rangle$$

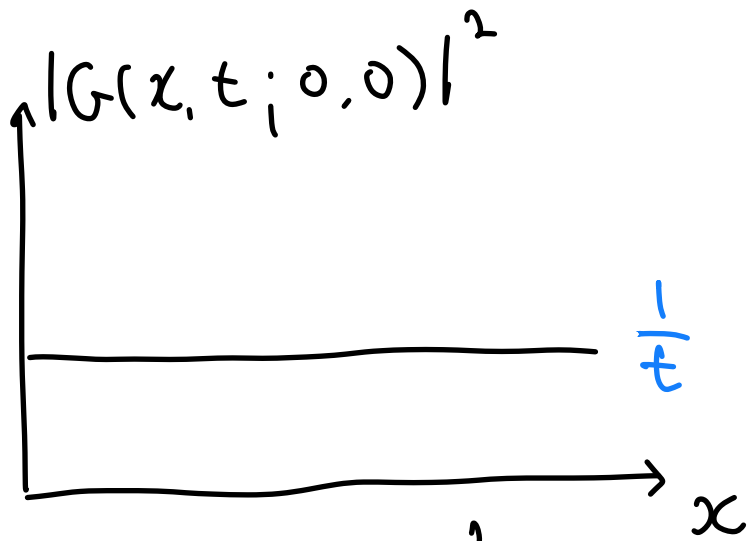
Single particle state  
 $(m=1)$

$$= \int \frac{d^3k}{(2\pi)^3} e^{i\vec{k} \cdot (\vec{r}-0)} e^{-i\epsilon_k t}$$

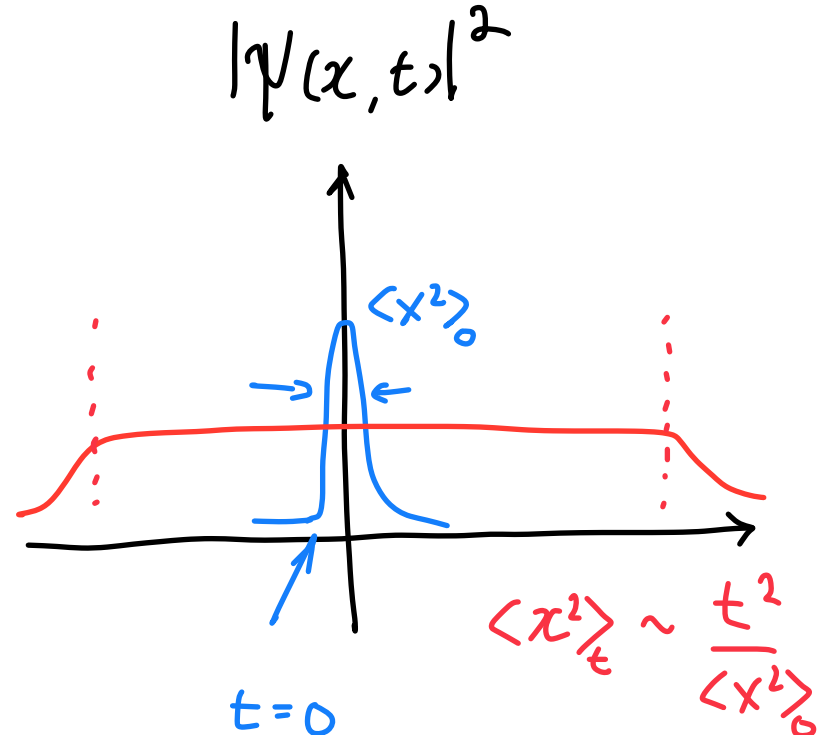
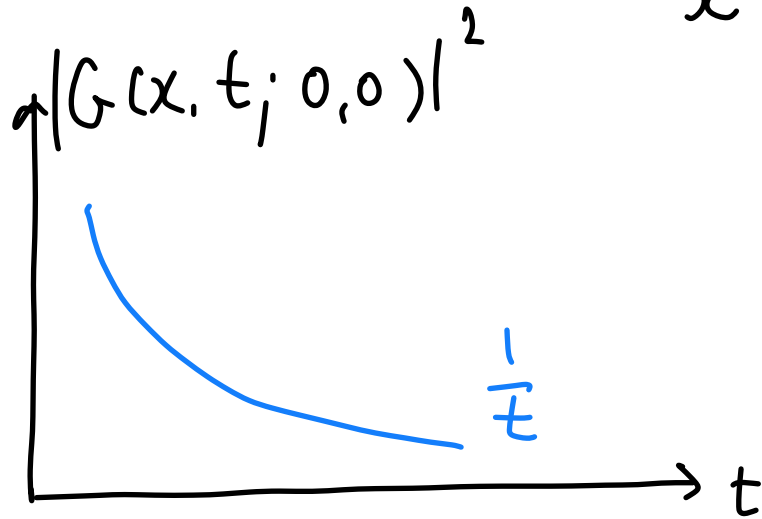
$\int \frac{d^3k}{(2\pi)^3}$  vol.  $\Omega$   $|\vec{k}|$   
 $\frac{4\pi^{-1/2}}{\sqrt{2\pi i}}$   $e^{i \frac{x^2}{2t} \cdot \frac{1}{\hbar^{1/2}}}$   
 $\rightarrow \vec{r}$



QM 100

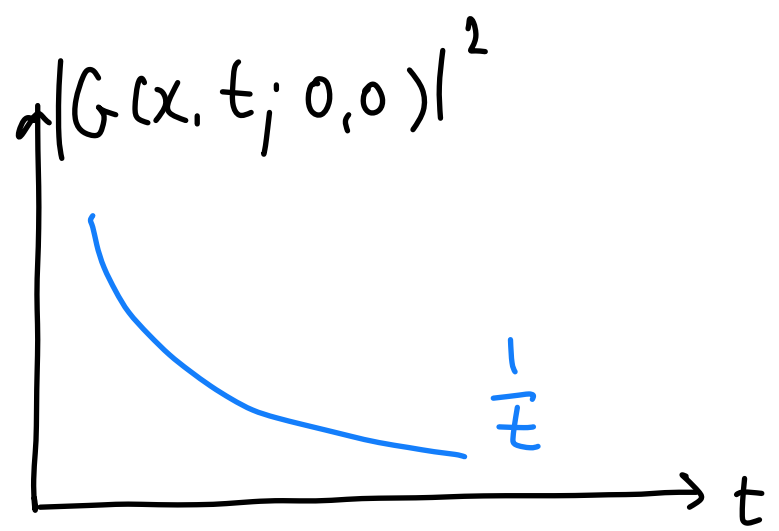
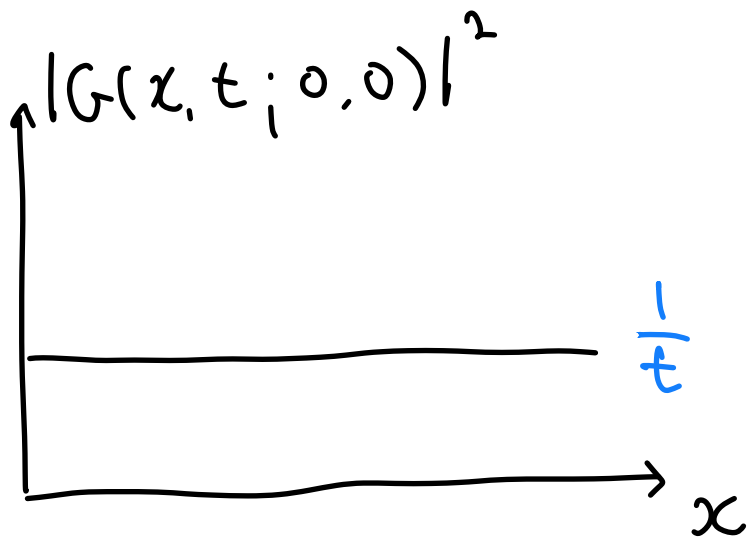


$0$   $\mathcal{D}$   
 $(0, 0)$   $(x, t)$

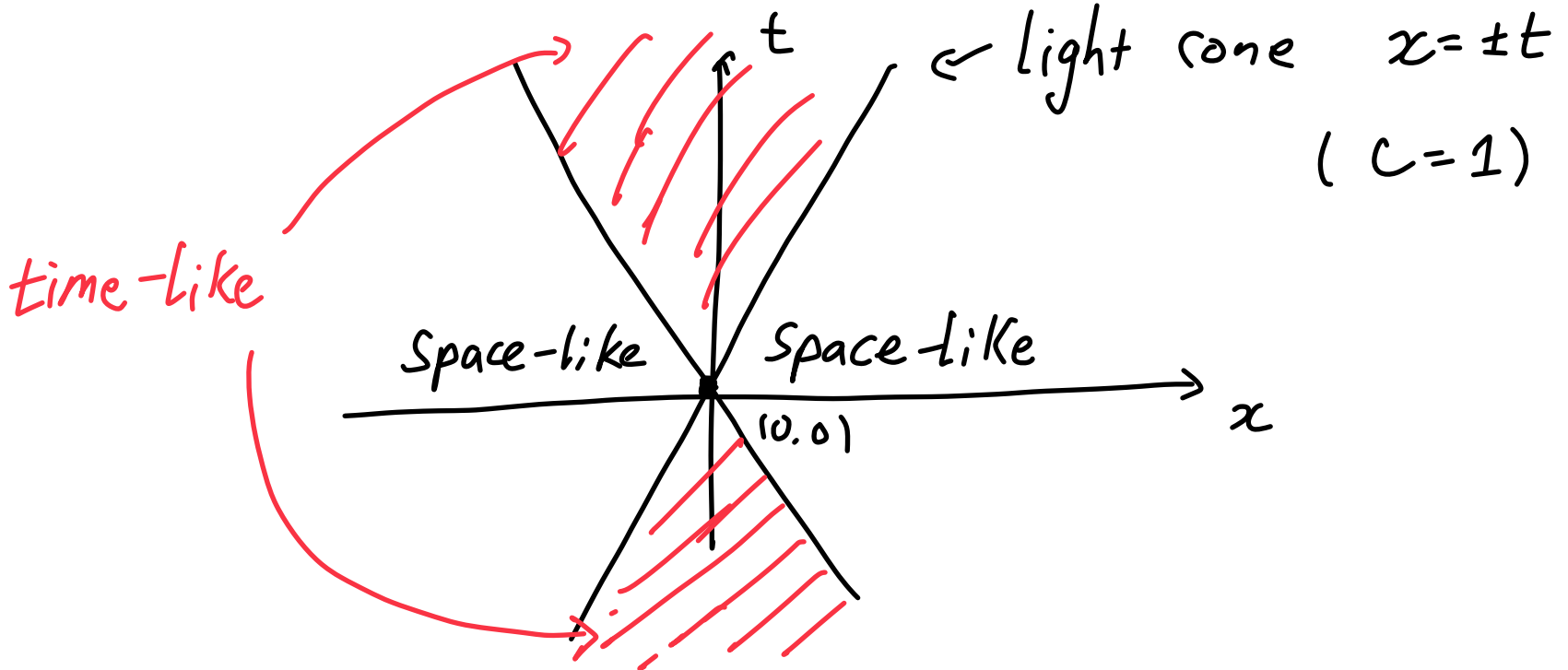


$\langle x^2 \rangle_0 \rightarrow 0$

$$|\psi(x, t)|^2 \longrightarrow |G(x, t; 0, 0)|^2$$



What about  $Z=1$  Relativistic QCP?





$$\begin{aligned}
 G_R(\vec{r}, t; 0, 0) &= \int \frac{1}{(\omega + i\delta)^2 - k^2} e^{i(kx - \omega t)} \frac{d\omega}{2\pi} \frac{dk}{2\pi} \\
 &= \Theta(t) \left[ -i \int \frac{e^{ik(x-t)}}{2k} \frac{dk}{2\pi} + i \int \frac{e^{ik(x+t)}}{2k} \frac{dk}{2\pi} \right] \\
 &= \frac{1}{2} \Theta(x-t) - \frac{1}{2} \Theta(x+t) \\
 &= \frac{1}{2} \Theta(t) \Theta(t - |x|)
 \end{aligned}$$

