

Phys525: Quantum Condensed Matter Physics:

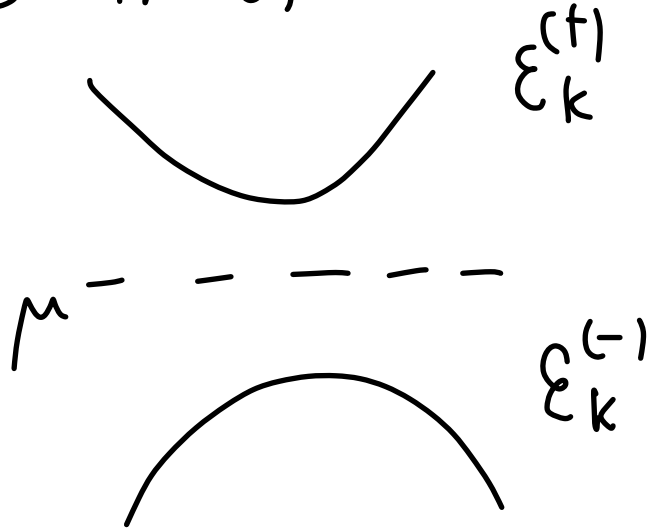
Episode 4: PHC as an internal emergent symmetries (ES):
slightly formal stuff and PHC in SSB states

- Emergent PHC/CC symmetries: case analysis
- A) due to interactions with background crystal structures, i.e. band structures with $U(1)$ charge symmetries. (Often need tuning!!)
- B) due to mutual interactions between particles via spontaneous symmetry breaking. (Very robust!!)

Today's plan

- Some more formal/general discussions on PH transformations and basic algebraic structures.
- More realistic Haldane lattice model
- PHC in SSB states (particle-particle interaction driven)

C-Transformation \rightarrow



$$b_{k,+}^{\dagger} \rightarrow b_{-k,-}$$

$$b_{k,+} \rightarrow b_{-k,-}^{\dagger}$$

$$H \rightarrow H \quad \text{if} \quad \epsilon_k^{(+)} = -\epsilon_{-k}^{(-)}$$

More formally, $C^{\dagger} H(p, m) C = -H(-p, m)$

$$H = \sum_k \epsilon_k^{(+)} b_{k,+}^{\dagger} b_{k,+} + \epsilon_k^{(-)} b_{k,-}^{\dagger} b_{k,-}$$

$$= \left[\sum_k \epsilon_k^{(+)} b_{k,+}^{\dagger} b_{k,+} - \epsilon_{-k}^{(-)} b_{-k,-}^{\dagger} b_{-k,-} \right]$$

$\rightarrow a_{k,+}^{\dagger} a_{k,+}$

- More technical discussions on PHC transformation and PHC symmetry as supplementary materials.

(A)

$$H = \sum_{i,j} \Psi_i^\dagger H_{ij} \Psi_j = \tilde{\Psi}^\dagger \mathcal{H} \tilde{\Psi}$$

$$\tilde{\Psi} \rightarrow [\tilde{\Psi}'^\dagger]^\dagger = [\tilde{\Psi}^\dagger U^\dagger]^\dagger, \quad \tilde{\Psi}^\dagger \rightarrow \tilde{\Psi}'^\dagger = \tilde{\Psi}^\dagger U^\dagger$$

$$H \Rightarrow \tilde{\Psi}'^\dagger U^\dagger \mathcal{H} U \tilde{\Psi}'$$

$$\boxed{\text{Tr } AB = \text{Tr } B^T A^T}$$

$$= \tilde{\Psi}'^\dagger \underbrace{U^\dagger \mathcal{H} U}_{\mathcal{H}'} \tilde{\Psi}',$$

thus

$$-U^\dagger \mathcal{H} U = \mathcal{H} = -U^\dagger K \mathcal{H} K U$$

then H remains invariant.

$$\textcircled{B} \quad \Psi_i \rightarrow \Psi_i^{\dagger} = U_{ji}^{\dagger} \bar{\Psi}_j^{\dagger}, \quad \Psi_i^{\dagger} \rightarrow \Psi_i' = U_{ij} \bar{\Psi}_j$$

$$H = \sum_{i,j} \bar{\Psi}_i^{\dagger} H_{ij} \Psi_j = \sum_{i,j} \Psi_{\alpha}^{\dagger} U_{i\alpha} H_{ij} U_{\beta j}^{\dagger} \bar{\Psi}_{\beta}^{\dagger}$$

$$\underbrace{U_{\beta j}^{\dagger} H_{ij} U_{i\alpha}}_{U_{\beta j}^{\dagger} H_{ji}^* U_{i\alpha}}$$

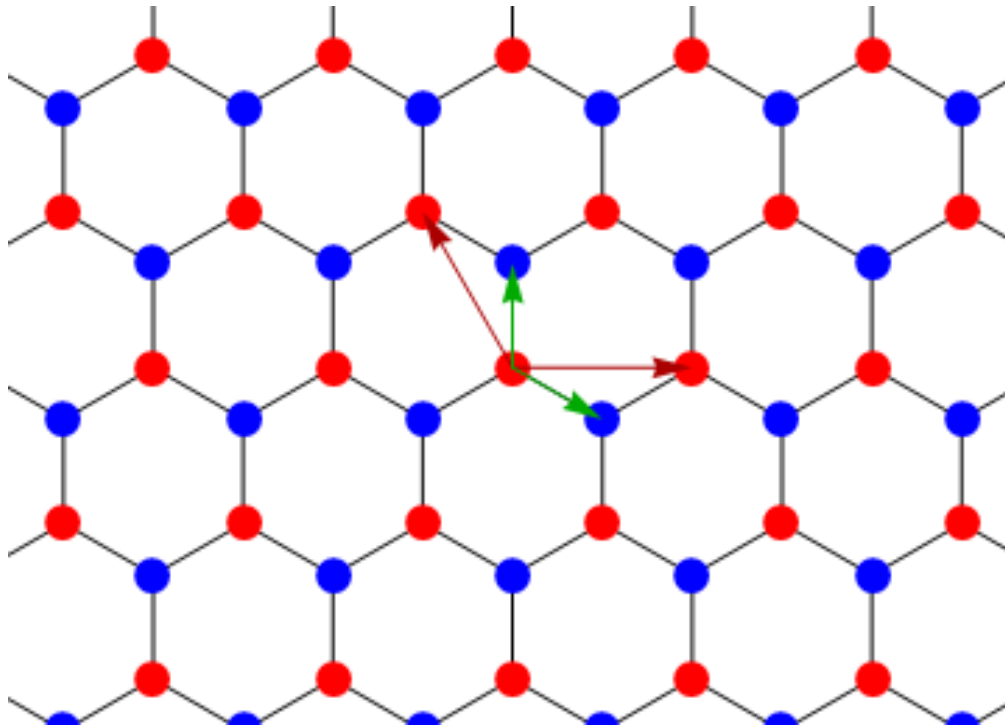
$$= \sum_{\alpha, \beta} \bar{\Psi}_{\beta}^{\dagger} H'_{\beta\alpha} \Psi_{\alpha}, \quad H'_{\beta\alpha} = -U^{\dagger} K H K U$$

$$H'_{\beta\alpha} = -U^{\dagger} K H K U$$

$$H = -C^{\dagger} H C \quad \text{or} \quad -H(-p, m) = C^{\dagger} H(p, m) C$$

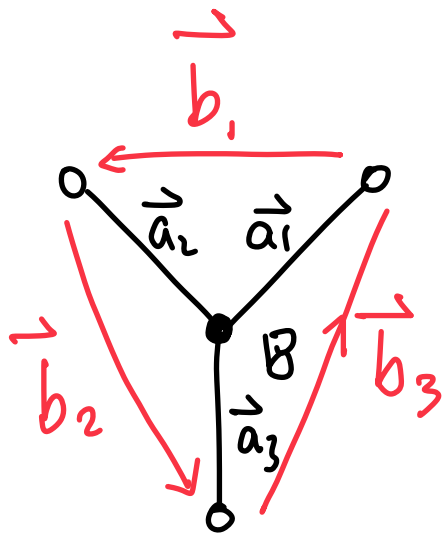
if H is invariant.

Discussions on a more realistic Haldane model in honeycomb lattices



$$f = \begin{bmatrix} \psi_A^+ & \psi_B^+ \end{bmatrix} \begin{bmatrix} 0 & t_{ab} \\ t_{ba} & 0 \end{bmatrix} \begin{bmatrix} \psi_A(k) \\ \psi_B(k) \end{bmatrix} + \dots$$

$$t_{AB} = t_{BA}^* = \sum_{\alpha=1,2,3} e^{i\vec{k} \cdot \vec{a}_\alpha}$$

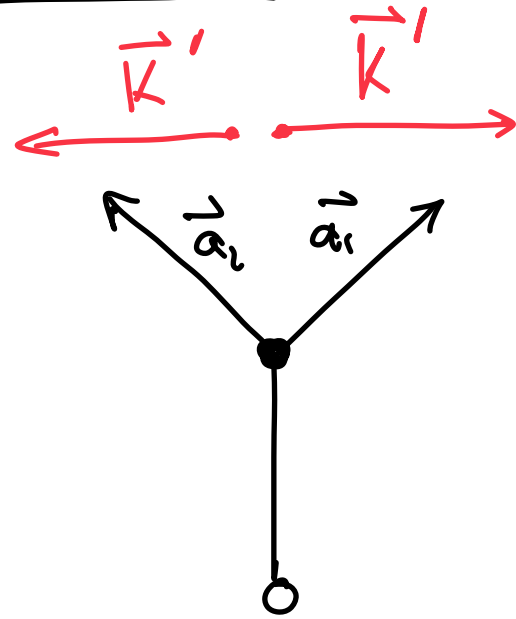
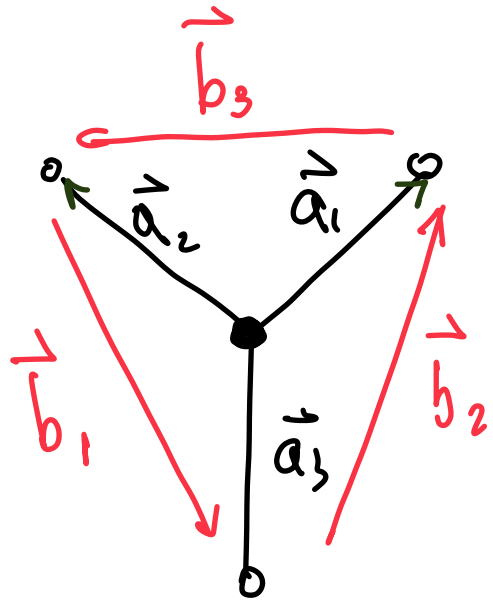


if $t_{AB}(\vec{k}) = 0$, then

$$(\vec{k} \cdot \vec{a}_1, \vec{k} \cdot \vec{a}_2, \vec{k} \cdot \vec{a}_3) = \left(0, \frac{2\pi}{3}, -\frac{2\pi}{3}\right) \text{ per}$$

$3! = 6$ solutions

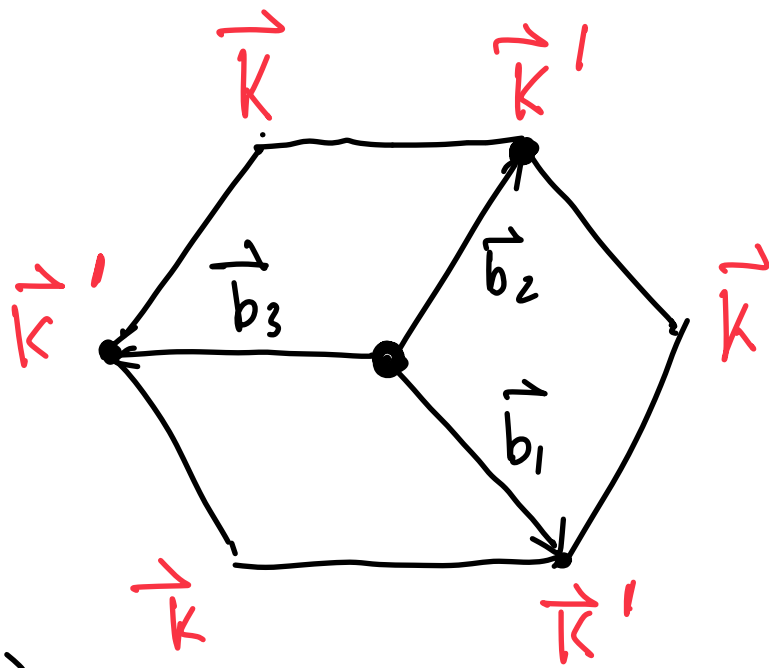
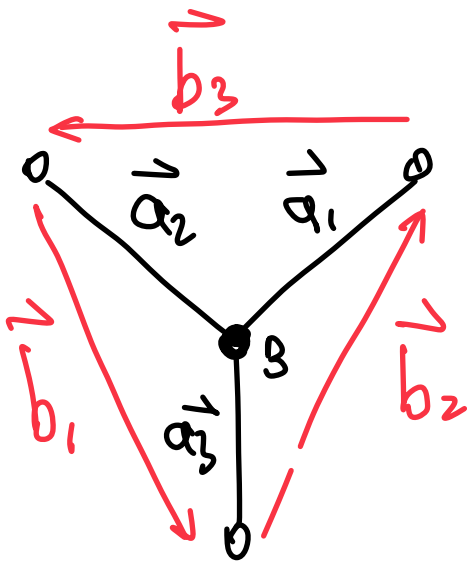
There three pairs of Solutions



$$\vec{k} \rightarrow \vec{k}' = -\vec{k}$$

$$(\vec{k}_1, -\vec{k}_1), (\vec{k}_2, -\vec{k}_2), (\vec{k}_3, -\vec{k}_3)$$

$$(0, \frac{2\pi}{3}, -\frac{2\pi}{3}) \rightarrow (0, -\frac{2\pi}{3}, \frac{2\pi}{3}). \quad \vec{k} \parallel \vec{b}_1, \vec{b}_2, \vec{b}_3$$



$$t_{ab}(\vec{k}) = \sum_{i=1,2,3} e^{i\vec{k} \cdot \vec{a}_i} = 0 + i\delta\vec{k} \cdot \sum e^{i\vec{k} \cdot \vec{a}_i} \vec{a}_i$$

$$\begin{bmatrix} 0 & t_{ab}(\vec{k}) \\ t_{ba}(\vec{k}) & 0 \end{bmatrix} = \begin{bmatrix} 0 & \pi_1 + i\pi_2 \\ \pi_1 - i\pi_2 & 0 \end{bmatrix}$$

$$t_{ab}(\vec{k}) \cong \Pi_1 + i\Pi_2 \quad \text{near } \vec{k}$$

$$t'_{ab}(\vec{k}) \cong -\Pi_1 + i\Pi_2 \quad \text{near } \vec{k}'$$

$$t_{ab}(\vec{k}) = 0 + i\delta\vec{k} \cdot \sum_i e^{+i\vec{k} \cdot \vec{a}_i} \vec{a}_i$$

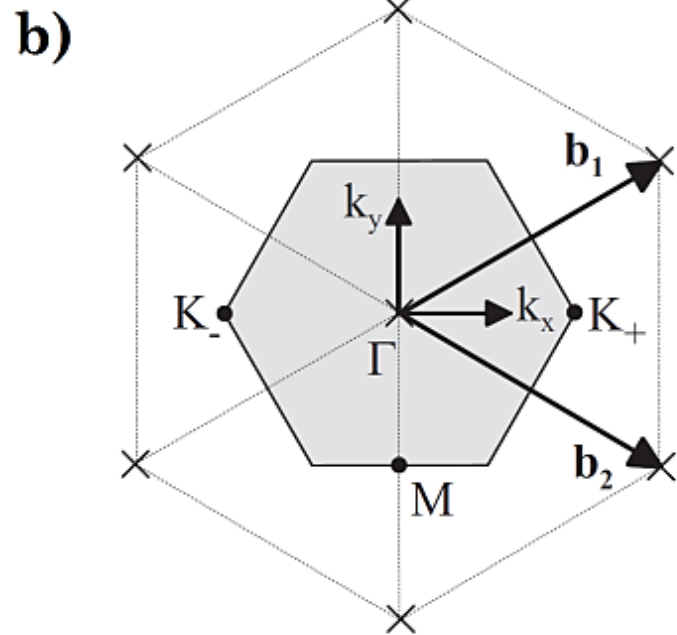
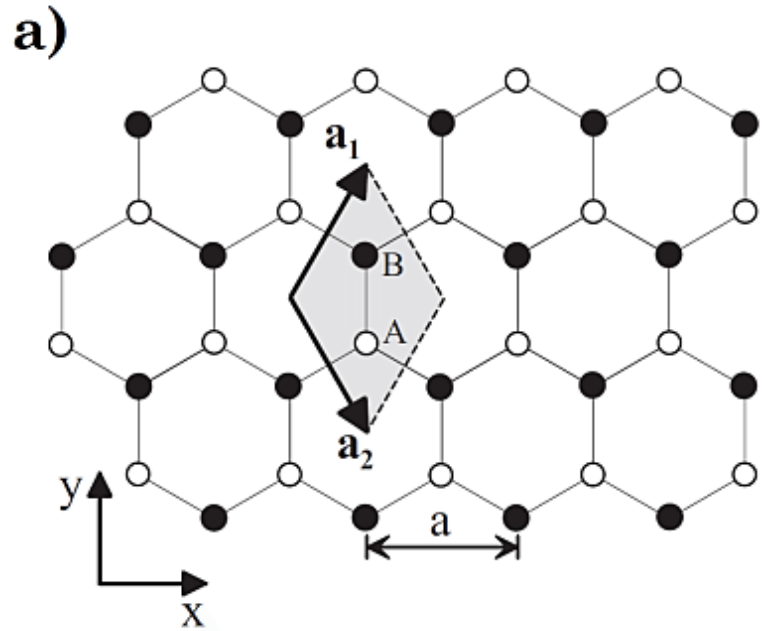
$$t_{ab}(\delta\vec{k}) = -t_{ab}^*(\delta\vec{k})$$

↑ Near \vec{k}

↑ Near \vec{k}'

$$H(p, m) = \tau_z \otimes \sigma_x p_x + \sigma_y p_y + m \sigma_z$$

τ_z : Valley Subspace ; $\sigma_{x,y,z}$: Sublattice Subspace



PHC Transformation:

$$C = K \tau_z \otimes \sigma_x$$

↘ Complex Conjugate

$$C^\dagger H(p, m) C = -H(-p, m)$$

Time Reversal Transformation:

$$T = K \tau_x, \quad T^\dagger H(p, m) T = +H(-p, m)$$

Chiral Transformation

$$S = i \tau_y \otimes \sigma_x, \quad S^\dagger H(p, m) T = -H(p, m)$$

more interesting case:
PHC symmetry without $U(1)$ symmetry
via SSB