

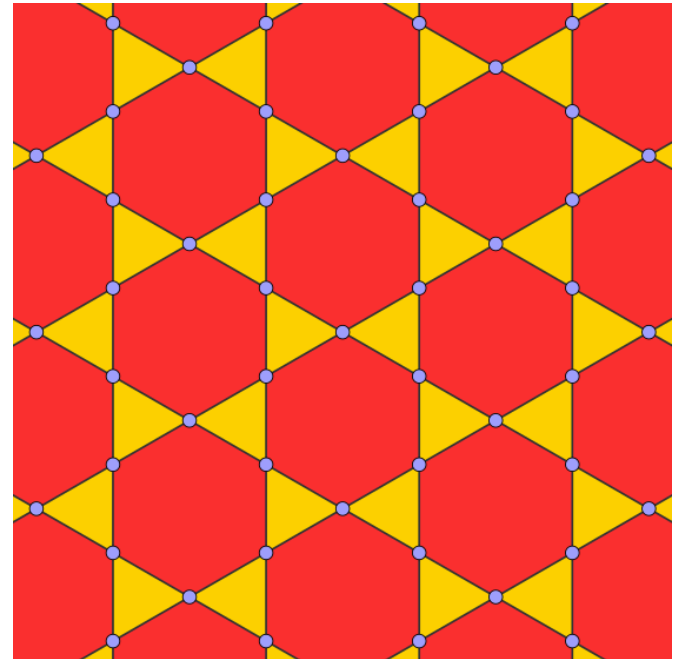
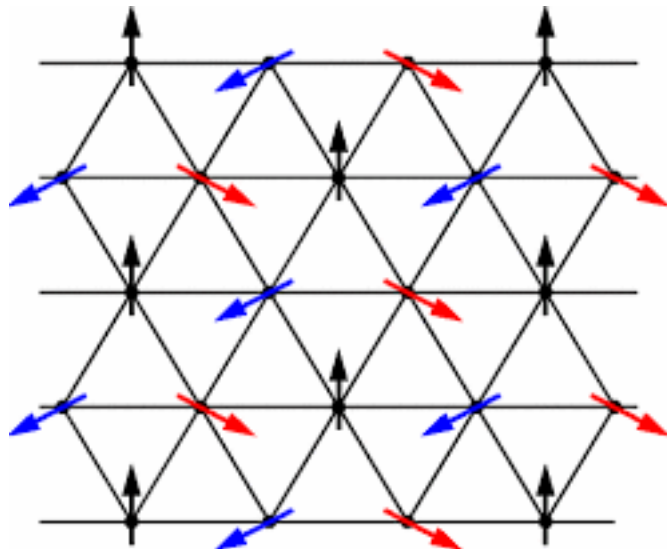
Phys525:
Quantum Condensed Matter Physics:
emergent symmetry and phenomena

Topological States, Topological ordered states and SPT: II



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Triangle and kagaome lattices (with $SU(2)$ exchange interactions; Z_2 spin liquids)

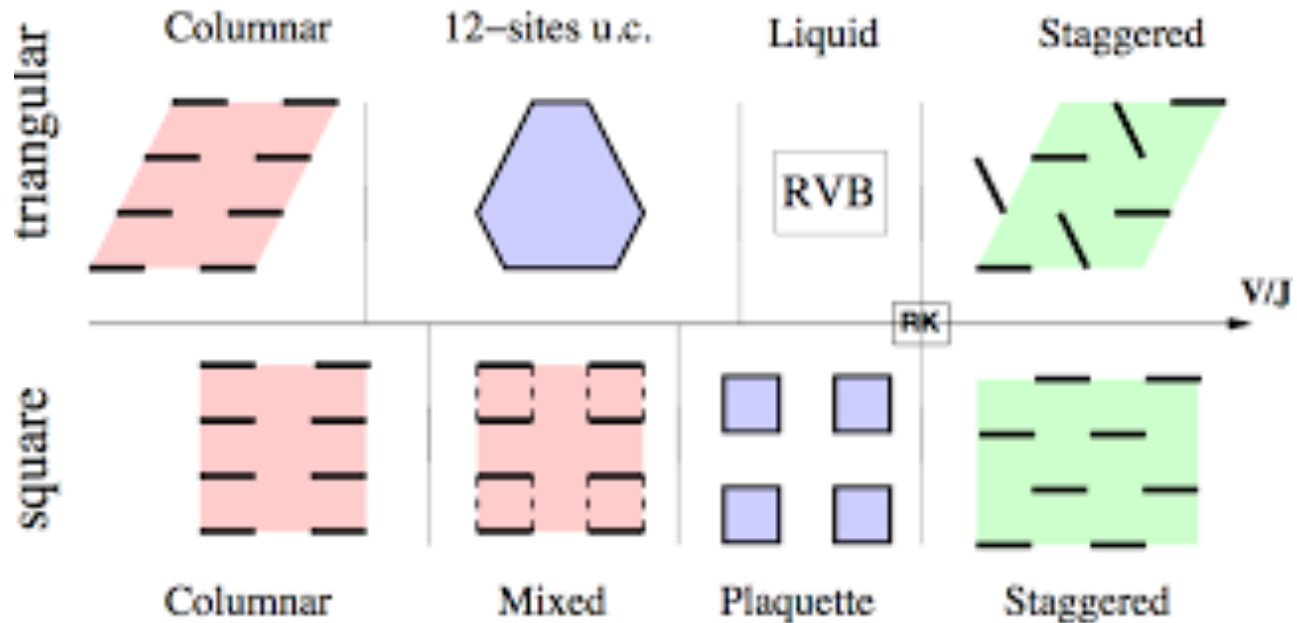
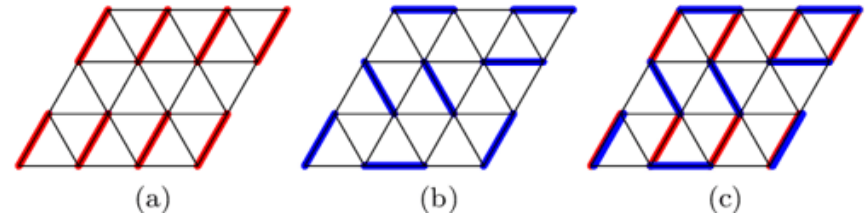


quantum dimer model (for spin liquids)

(Rokhsar, Kivelson, 89; Moessner, Sondhi, Fradkin, 2002-2003)

$$\begin{aligned}
 H = & -t \sum_r (|\text{dimer}\rangle \langle \text{dimer}| + \text{h.c.}) \\
 & +V \sum_r (|\text{dimer}\rangle \langle \text{dimer}| + |\text{dimer}\rangle \langle \text{dimer}|) \\
 & -h \sum_l (|\text{monomer}\rangle \langle \text{monomer}| + \text{h.c.}) \\
 & -\mu \sum_l (|\text{monomer}\rangle \langle \text{monomer}|),
 \end{aligned}$$

$\circ =$
 $\circ =$



Affleck-Kennedy-Lieb-Tasaki, 1987

1983

$$H_{AKLT} = J \sum_{\langle ij \rangle} (S_i \cdot S_j + \alpha (S_i \cdot S_j)^2), \quad d=0 \text{ Haldane}$$

$\alpha = 1/3$, AKLT

bilinear-biquadratic Model

[S=1, SU(2)] Model

Two site solutions

$$(S_i \cdot S_j) = \frac{s^2 - S_i^2 - S_j^2}{2} = \frac{s^2}{2} - 2$$

$$\alpha=0 \quad \begin{array}{cc} \frac{2}{1} & J \\ \frac{1}{S=0} & -J \\ & -2J \end{array}$$

$$d=1/3 \quad \begin{array}{cc} \frac{2}{0} & \frac{4J}{3} \\ & \boxed{-\frac{2J}{3}} \\ & \frac{1}{\boxed{-\frac{2J}{3}}} \end{array}$$

$$H_{AKLT} = 2J \sum_{\langle ij \rangle} P_{ij}^{(S=2)}, \quad J > 0$$

Representation

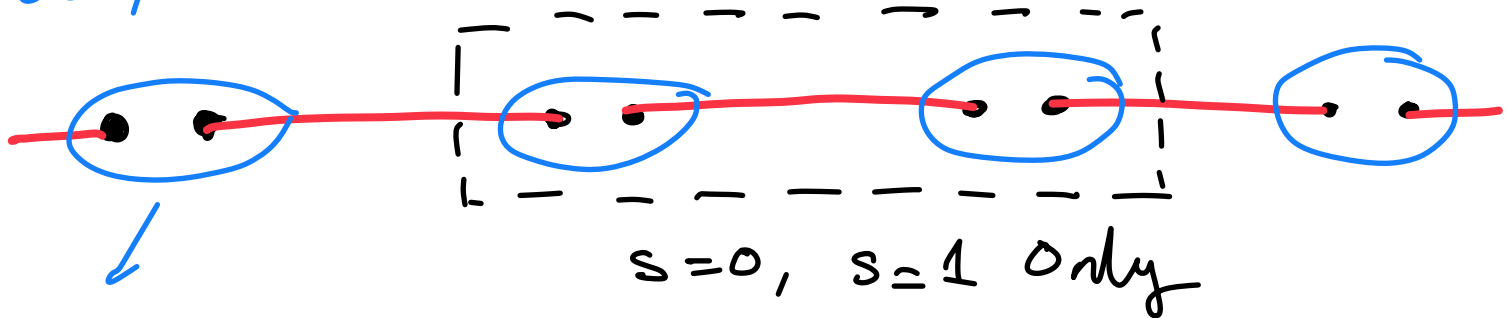
split $|S=1\rangle$
↓ ↓
 $|\uparrow\rangle, |\downarrow\rangle$

To impose $P_{ij}^{(2)}$

$$\left\{ \begin{array}{l} |\uparrow\uparrow\rangle \quad S_z = 1 \\ |\downarrow\downarrow\rangle \quad S_z = -1 \\ \frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle + |\downarrow\uparrow\rangle) \end{array} \right.$$

$$\frac{1}{\sqrt{2}} (|\uparrow\downarrow\rangle - |\downarrow\uparrow\rangle) \quad (\text{Needs to be removed})$$

VBS, Valence-bond solid as the ground state



"Project out $|s=1\rangle$ "

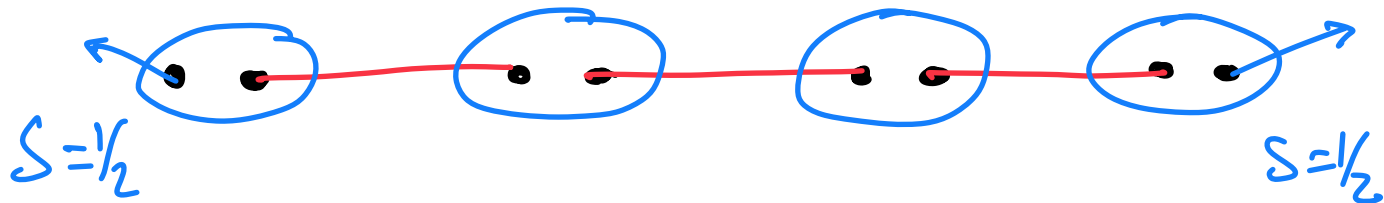
Elementary Properties known since 90s

1) String order by Denijs, Rommelse (1989)

$$\langle S_i^z S_j^z \rangle \rightarrow 0 \text{ as } i-j \rightarrow \infty,$$

but $\langle S_i^z \underbrace{e^{i\pi S_k^z}}_{i < k < j} S_j^z \rangle \neq 0 \text{ as } i-j \rightarrow \infty.$

2) Open boundary AKLT has 4-fold degeneracy.
(fractionalized)



(These features extendable to the Haldane phase.)

More Sophisticated Interpretation: Kennedy, Tasaki, 1992
on Haldane phase

$$\begin{cases} U^\dagger H U = \tilde{H} \\ U^\dagger S_i^z \prod_k e^{i\pi S_k^z} S_j^z U = S_i^z S_j^z \end{cases}$$

\tilde{H} has $\mathbb{Z}_2 \otimes \mathbb{Z}_2$ symmetry, or $\mathbb{Z}_2 \otimes \mathbb{Z}_2 = D_2$

(π -Rotations about x, y, z axes)

Ground state 4-fold degenerate; ferromagnetic Order.

$$U = \prod_{j < k} e^{i\pi S_j^z S_k^z}, \quad \tilde{H} = TBC.$$

Matrix-Product-State Rep (DMRG)

(~2011, Modern studies)

$$|g, s\rangle = \sum_{\{s_i\}} \alpha_{s_1, \dots, s_n} |s_1, \dots, s_n\rangle, \quad s_i = 0, \pm 1$$

$$\alpha_{s_1, \dots, s_n} = \text{Tr} A^{s_1} A^{s_2} \dots A^{s_n}$$

$$A^0 = +\sqrt{\frac{1}{3}} \sigma^z, \quad A^+ = -\sqrt{\frac{2}{3}} \sigma^+, \quad A^- = -\sqrt{\frac{2}{3}} \sigma^-$$

(ideas coming from QI, Verstraete, Cirac, Mung, 2008;
Perez-Garcia,
Applications in CMT start in 2009 on; Numerical applications
DMRG in 2011, Schollwock.)

Why SPT ?

(Symmetry protected topological states)

- are all spin disordered states (SU(2) spins) topologically ordered, i.e. is 1D AKLT (more discussions today) topologically ordered in the way how 2D Z_2 spin liquids are ordered (nonlocal order)?
- What happens to Z_2 TI or TSCs discussed in free fermion models with TRS when interactions are induced? (Next week).

Topological ordered states vs Symmetry protected topological states

a series of impressive attempts of using entanglement measure by
Chen, Gu, Wen et.al, 2009-2012 (in Group 2)

- Topologically ordered states (TOS) with long range entanglements, or topological entropy; or topological degeneracy on torus etc. (Levin, Wen, 2005; Kitaev, Preskill, 2005)
- Symmetry protected topological states (SPT) with short range entanglement entropy; no topological degeneracy on torus etc (discussions on AKLT).
- SPT states belong to a well defined phase only if the symmetry is respected. Otherwise not. (Can be a presentation topic.)

On surfaces of SPT vs TOS

- SPT always has **gapless surfaces or degenerate surfaces** disregarding how you cut. To gap it, **surfaces** are either SSB or intrinsic topologically ordered (no contradictions to what we had before!! see next page or come back next week).
- TOS not always have gapless boundaries. ex: Z_2 Spin liquids. However, it can have, say in FQHE.

What do we mean by saying “topological order” ?



Topics suggested for P3

- 1) Crossover from AKLT or other SPT into trivial phases.
- 2) Kennedy-Tasaki hidden $D_2 = Z_2 \times Z_2$ symmetry in the 1D Heisenberg spin model—extending AKLT feature to the Haldane phase.
- 3) Topological Entanglement entropy: how to choose subsystems to pick up the gamma- term/topological term —more formal one.

Topological entropy as a measure of topological order

(Levin and Wen, 05; Kitaev and Preskill, 05)

$$S_{VN} = \alpha L - \gamma, \gamma = \ln D$$

$$D = \sqrt{\sum_a d_a^2}$$

$$\rho_A = \text{Tr}_B \rho_{AB}$$

$$S_{VN} = -\text{Tr} \rho_A \ln \rho_A$$

- based on von Neumann entropy and quantum dimension D . $D=2$ for the toric code, \sqrt{q} for FQHE with filling factor $1/q$, $q=3,5,7..$